

A Study of Translational Diffusion
Constants of Paramagnetic Species
In Solution

by

Hasan Bakor Abdullah Balkhoyor

A Thesis Presented to the

FACULTY OF THE COLLEGE OF GRADUATE STUDIES

KING FAHD UNIVERSITY OF PETROLEUM & MINERALS

DHAHRAN, SAUDI ARABIA

In Partial Fulfillment of the
Requirements for the Degree of

MASTER OF SCIENCE

In

CHEMISTRY

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Balkhoyor, Hasan Bakor Abdullah, M.S.

King Fahd University of Petroleum and Minerals (Saudi Arabia), 1993

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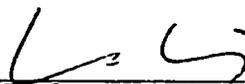
COLLEGE OF GRADUATE STUDIES

This thesis, written by **HASAN BAKOR BALKHOYOR** under the direction of his Thesis Advisor and approved by his Thesis Committee, has been presented to and accepted by the Dean of the College of Graduate Studies, in partial fulfillment of the requirements for the degree of **MASTER OF SCIENCE in Chemistry**.

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**THIS THESIS IS DEDICATED TO MY
PARENTS,
WIFE,
and
CHILDREN.**

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خلاصة الرسالة

اسم مقدم الرسالة : حسن بكر عبدالله بالخوير

عنوان الرسالة : دراسة معامل الانتشار الانتقالي للمركبات ذات الخواص المقاطيسية في المحاليل

التخصص : للكيمياء الفيزيائية

تاريخ التقديم : يوم الثلاثاء ٢ محرم ١٤١٤ هجرية

في هذا البحث تمت دراسة معامل الانتشار الانتقالي لمركبين لهما خواص مقاطيسية ، الاول : ثنائي كلورو - ثينين (٧- بنتادايثيل حقيقي) فثاديوم ، ويختصر (Cp_2VCl_2) في الكلوروفورم كمنيب، والثاني : S - ميثيل - N - O - مئوكسي - ساليسايدين - هايدرازين - كلويدايشوتو - فنتشولين - لكسي الفثاديوم ، ويختصر (VO[5-MeO Sal SB] (Phen))، في ثنائي كلورو ميثان كمنيب. وذلك عند درجة حرارة الغرفة. وقد استخدمت طريقة الانتشار خلال الانبوب الشعرية لأيجاد هذه المعاملات عن طريق تتبع شدة الطيف المقاطيسي مع مرور الزمن بواسطة جهاز طيف الطنين المقاطيسي الألكتروني . وبرسم علاقة للتغير بين الزمن وشدة الطيف ثم ايجاد افضل منحنى نظري يحاكي المنحنى التجريبي من خلال حل القانون الانتشار الثاني ل-Fick'S وتعويض في هذا القانون لأيجاد المعاملات بواسطة برنامج كمبيوتر كتب لهذا الغرض. وبأستخدام علاقة Stokes - Einstein التي تربط بين معامل الانتشار الانتقالي والزوجة وجد نصف القطر الحركي للجزيئات ومنه حسبت حجوم هذه الجزيئات. فكان معامل الانتشار الانتقالي للمركب الأول 0.520×10^{-5} سم²/ث ، وحجمه 250×10^{-8} سم³. ومعامل الانتشار للمركب الثاني 655×10^{-5} سم²/ث وحجمه 275×10^{-8} سم³.

في الجزء الثاني من الرسالة استخدمت اربع مجموعات مختلفة لمعاملات مقاطيسية للمركب الأول نكرت في الدوريات العالمية لتحقيقها على قيم عرض الخطوط للطيف المقاطيسي عند درجات حرارة مختلفة في مادة التولوين ، لأيجاد معامل التباين الدورتي (N) وكذلك محور الدوران لهذا المركب. وقد وجد ان مجموعة واحدة من المعاملات المقاطيسية حققت قيم عرض خطوط الطيف المقاطيسي بمعامل تباين دورتي (N) = 2.7 ± 0.5 وذلك على المحور $Z' = Y$ ، وتطابقت هذه القيمة مع نموذج Stokes-Einstein ومعادلة القيمة المانونة (Allowed - Values Equation) ومن معامل التباين الدورتي حسب زمن الدوران الذي يرتبط بعلاقة رياضية مع الزوجة والحرارة والحجم الجزيئي ومعامل التجانب بين جزيئات المذاب والمنيب (k) ، الذي حسبته قيمته 0.15 . وهذه القيمة تدل على ان المركب لكثرت ازلقي من المركب التمونجي فنادايل اسيتل اسيتونات (VOAA) (والذي يبلغ مقدار التجانب له في التولوين 0.53 ، ويؤيد ذلك معامل الألتصاق (S) الذي يأخذ القيم ($0 < S < 1$) فعندما يكون S يساوي صفر فنلك يدل على الأزلقي الكامل ، والواحد الألتصاق الكامل ، وبحساب معامل الألتصاق للمركب الأول وجد ان قيمته 0.0875 . وبمراجعة اشكال هذه الجزيئات يمكن تقدير ذلك فكما كان الجزيئ منتظما قل معامل التجانب بين جزيئات المنيب والمذاب وكذلك قل معامل الألتصاق ، والعكس كلما كان الشكل الجزيئ غير منتظم ازداد معامل التجانب و كذلك معامل الألتصاق .

ABSTRACT

The translational diffusion constants of two paramagnetic species, dichlorobis(η -cyclopentadienyl)vanadium, abbreviated as (Cp_2VCl_2) in chloroform, and S-methyl-N-5-methoxy-salicylidenehydrazine-carbodithioatophenanthroline-oxovanadium(IV), abbreviated as VO[5-MeO-Sal SB](Phen) in methylene chloride were measured at room temperature (25 ± 1 °C) by a capillary diffusion method. The translational diffusion processes were monitored by electron spin resonance spectral intensities as a function of time. The total spectral intensities were plotted as a function of time and from the best least squares fit a single value of the translational diffusion coefficient was obtained. A computer program based on Fick's second law of diffusion for the capillary diffusion cell has been written to get a single value of the diffusion constant. Using Stokes-Einstein for translational diffusion, the hydrodynamic radii and the molecular hydrodynamic volumes of the two paramagnetic species in solution were found to be 0.520×10^{-3} and 0.655×10^{-3} cm³/s, and 250 and 275 Å³, respectively.

From an analysis of the temperature dependent ESR line widths of Cp_2VCl_2 in toluene in the motional narrowing region by using four sets of the magnetic parameters available in the literature, only the set by Peterson and Dahl could explain the line width data. It was found that Cp_2VCl_2 was undergoing anisotropic reorientation with $N = 2.7 \pm 0.5$ at an axis $z' = Y$, where N is the ratio of R_{\parallel}/R_{\perp} and R_{\parallel} is the rotational diffusion constant along the molecular Z axis and R_{\perp} is the rotational diffusion constant perpendicular to the molecular Z axis. This value is supported by the Stokes-Einstein model and the allowed-values equation.

In addition, the rotational correlation times as a function of temperature were determined from the line width data. The anisotropic interaction parameter κ , determined from the plot of τ_c vs η/T was found to be 0.150. The anisotropic interaction parameter is a measure of the anisotropy of intermolecular interactions. The values show that Cp_2VCl_2 in toluene is more slip than VOAA in toluene which is more stick and its κ is 0.53. The Stickiness factor (S) is calculated to be 0.0875 for Cp_2VCl_2 , which is close to the slip limit ($S=0$). Therefore, it seems that κ value decreases with decreasing solute corrugation.

CHAPTER ONE

INTRODUCTION

CHAPTER 1

INTRODUCTION

1.1 INTRODUCTION

The purpose of this research is to study the translational motion of two paramagnetic species in solution by the capillary diffusion method. The first complex is *S*-methyl-*N*-5-methoxy salicylidene hydrazine carbodithioatophenathroline oxovanadium(IV), in methylene chloride, written simply as, VO[5-MeO-Sal SB](Phen), which has a monomeric octahedral structure as shown in Fig.(1.1) [1]. The second complex is dichlorobis(η -cyclopentadienyl)vanadium in chloroform, written as Cp₂VCl₂, its structure shown in Fig.(1.2). The translational motion will be measured at room temperature by a capillary diffusion method [2,3,4,5]. It has been previously shown that the translational diffusion processes can be monitored by electron spin resonance spectral intensities as a function of time [6]. The total spectral intensities are plotted as a function of time and from the best least-squares fit a single value of translational diffusion coefficient can be obtained [3]. Using the Stokes-Einstein equation the hydrodynamic radii of the two paramagnetic species in solution can be calculated.

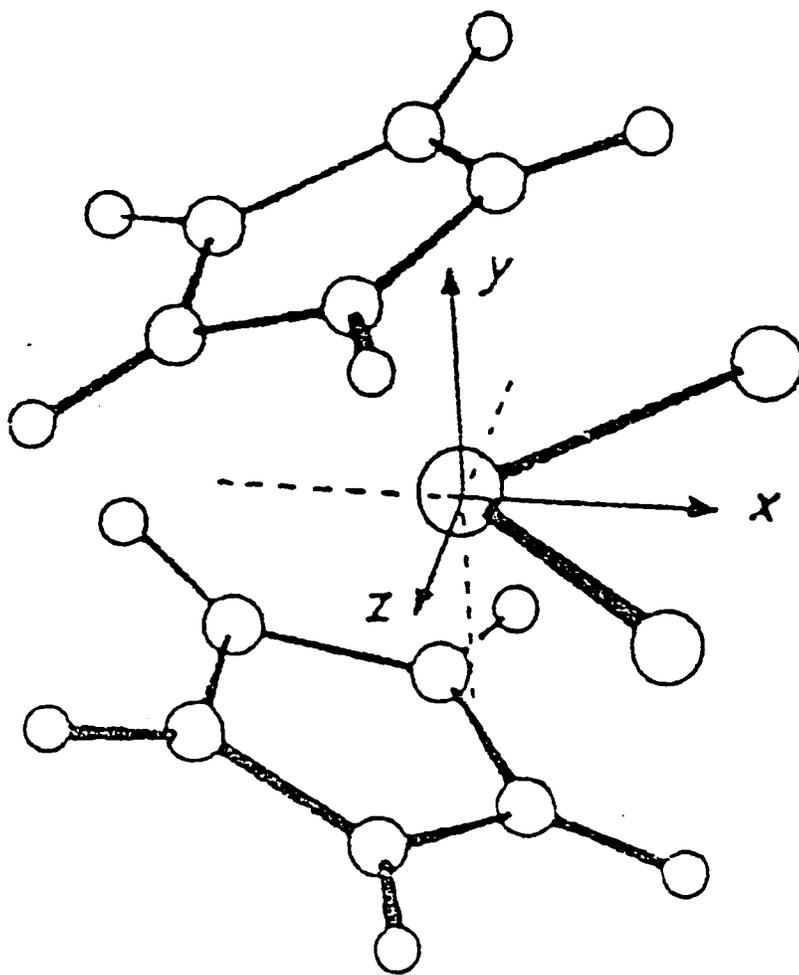


Fig. 1.1: Cp_2VCl_2

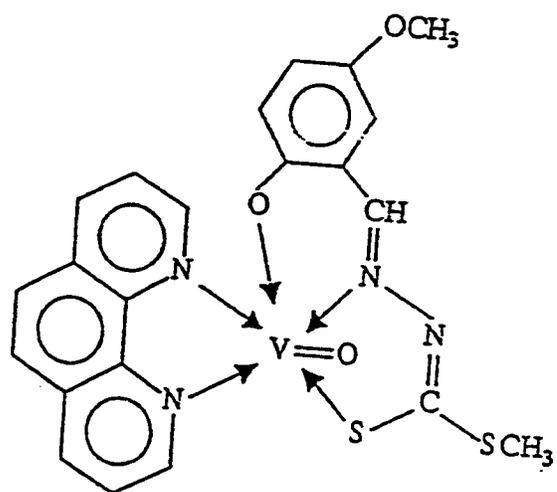


Fig. 1.2: VO[5-MeO-Sal SB](Phen)

The translational diffusion coefficient D is related to the viscosity by the Stokes-Einstein equation

$$D = \frac{kT}{6\pi\eta r_0} \quad (1.1)$$

where k is the Boltzmann constant, T is the absolute temperature, η is the shear viscosity, and r_0 is the hydrodynamic radius of an equivalent spherical volume [7], from which the molecular volume of the paramagnetic species can be determined. In the past the molecular volume has been determined using the Dreiding model [8], a computer program (Desktop Molecular Modeller) [1], and experimentally using such techniques as the porous disk method [9], and the electron spin resonance/spin exchange method [10].

The value of r_0 which determined experimentally by the capillary diffusion method will be used in recalculating the κ value from the correlation time, τ_1 , which is analyzed from molecular reorientation of paramagnetic species in solution as a function of temperature [1]. For spherical top or linear molecule this correlation time can be expressed as [11],

$$\tau_R = \left(\frac{4\pi r_0^3}{3}\right) \left(\frac{\kappa}{k_B}\right) \left(\frac{\eta}{T}\right) \quad (1.2)$$

Where η is the coefficient of shear viscosity of the solvent [12], and κ is an experimentally determined dimensionless coupling parameter called the anisotropic interaction parameter which measures the coupling of rotational motions of the probe to the translation modes of the fluid [11].

From the electron spin resonance relaxation studies [13], which have been analyzed in terms of the hydrodynamic free space model for molecular relaxation in liquids at various microwave frequencies [8,13,14], the analysis of the temperature dependent ESR line shapes in the motional narrowing and slow tumbling regions, the stickiness factors S were measured [15].

The stickiness factor is independent of molecular geometry and is zero in the slip and one in the stick limit, $0 < S < 1$. The stickiness factor can be calculated from the relationship

$$S = \frac{\kappa - C_{slip}^{hyd}}{1 - C_{slip}^{hyd}} \quad (1.3)$$

where C_{slip}^{hyd} is the empirical factor which depends on the non-hydrodynamic character of the friction and on the shape of the probe molecules [11], then from the recalculated value of κ we can obtain a more accurate value of S .

In the remainder of this chapter we give an overview of the transport process. The state of a system under diffusion is defined, then Fick's law of diffusion is set up mathematically and a solution of this law is derived. In Chapter II the apparatus used and its function in the experiment are described. In Chapter III we describe the experimental procedure used to set up the capillary diffusion cell using the ESR technique. In Chapter IV the analysis of the paramagnetic parameters are given. In Chapter V the conclusions of the experimental results and discussion are given.

1.2 TRANSPORT PHENOMENA

Transport phenomenon arises because any system in nonequilibrium will always try to approach equilibrium. A system is in equilibrium if the intensive state variables or parameters (temperature and density, for example) are constant at all times and have the same magnitude at all positions in the system.

From this definition follow other familiar properties of equilibrium state: at equilibrium, *gradients* of the parameters are absent, spontaneous processes do not occur, and there is no net exchange of matter or energy between the system and the surroundings, or between one part of the system and another.

If a system is in a state in which these conditions are not satisfied, the system is in nonequilibrium state. Suppose an equilibrium system is perturbed to a nonequilibrium state by some external or internal disturbance. At least one of the system parameters becomes a function of position. When the disturbance is removed, an irreversible decay process occurs spontaneously, and the system advances through a series of nonequilibrium state until equilibrium is reached.

The decay process is a transport process, the name arising because some quantity is transferred throughout the system as the system attempts to make all parameters independent of position.

A spatial dependence of the parameters is thus associated with a *flux* of some kind. It is logically assumed that the relation between a gradient X and its conjugate flux J is linear thus,

$$J = L X \quad (1.4)$$

If Eq. (1.4) is valid at all times, the system is said to be "close to equilibrium" at all times.

The rate at which the system approaches equilibrium is determined by the proportionality constant of Eq.(1.4) that is by the *transport coefficient* L .

It is a matter of experience that transport processes taking place in a system with n independent gradients can be represented by the relation,

$$J = \sum_{k=1}^n L_{ik} X_k \quad (i=1,2,3,\dots,n) \quad (1.5)$$

In the nomenclature of nonequilibrium thermodynamics, X is called a thermodynamic driving force, and L a phenomenological coefficient.

Equation (1.5) can be written in matrix notation,

$$\begin{bmatrix} J_1 \\ J_2 \\ \cdot \\ \cdot \\ \cdot \\ J_n \end{bmatrix} = \begin{bmatrix} L_{11} & L_{12} & \cdot & \cdot & \cdot & L_{1n} \\ L_{21} & L_{22} & \cdot & \cdot & \cdot & L_{2n} \\ \cdot & \cdot & \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot & \cdot & \cdot \\ L_{n1} & L_{n2} & \cdot & \cdot & \cdot & L_{nn} \end{bmatrix} \begin{bmatrix} X_1 \\ X_2 \\ \cdot \\ \cdot \\ \cdot \\ X_n \end{bmatrix}$$

which expresses the fact that if the off-diagonal elements of L exist, then the i th flux depends in a linear manner on the k th force as well as on its conjugate X_i . The relation between J_i and X_k is called thermodynamic coupling.

1.2.1 Uncoupled Transport Processes

When $L_{ik} = 0$ ($i \neq k$), Eq.(1.5) becomes $J_i = L_{ii}X_i$, a phenomenological representation of familiar empirical laws,

(1) Fourier's law of heat conduction,

$$\mathbf{q} = -\lambda \nabla T \quad (1.6)$$

The vector \mathbf{q} is the heat flux, the coefficient λ is the coefficient of thermal conductivity, and $\text{grad } T$ is the gradient of temperature.

(2) Fick's law of diffusion,

$$J_i = -D \nabla C_i \quad (1.7)$$

where J_i is the flux of chemical species i , C_i is the concentration of i , and D is the diffusion coefficient.

(3) Newton's law of viscosity,

$$\Pi_{yx} = -\eta \frac{\partial u_x}{\partial y} \quad (1.8)$$

Here Π_{yx} represents a momentum flux in the y direction caused by a fluid flow in the x direction with velocity component U_x , η is the coefficient of shear viscosity [16].

1.2.2 Coupled Transport Processes

The coupled phenomena arising if $L_{ik} \neq 0$, *thermal diffusion* of gases is an example of this processes in which the diffusion of gases depends not only on the concentration difference but also on the temperature difference. Also the *thermoelectricity* with coupling between flows of heat and of electrical current is another example [16].

1.2.3 Continuous and Discontinuous System

The fluid can be regarded as a continuum and the intensive state variables as continuous function of space and time. Such a system is called continuous. Also the phenomenological or transport coefficients, for both coupled and single processes, appertain to the fluid itself. At a microscopic level they could be described by the temperature

and pressure of the system, together with the microscopic properties of the components of the fluid, for example by the molecular dimensions and the intermolecular potential function, *thermal diffusion* is an example here for continuous system in which the vessel plays no part in the effect apart from containing the fluid. In a discontinuous system, the vessel plays an essential role, the fluid-vessel system is not a continuum, and the state variable are not continuous function of space because of discontinuities of the state variables. Here the phenomenological or transport coefficients are now function of the fluid and of the vessel [16].

1.2.4 Transport Processes At The Steady State

A process which involves time-dependent parameters is not easy to setup experimentally or to analyze theoretically. The process becomes far simpler to deal with if the state parameters are time independent. This can be achieved when the parameter gradients are externally constrained to constant values. In fact, the empirical transport laws, Fourier's law, etc., were conceived for this case.

Today, as well, most transport experiments are planned so that the state parameters are independent of time. The simplification applies especially to coupled phenomena. It is an experimental fact that, the case of a system with n gradients present, if one gradient is constrained by external

means, all gradients become constant in the course of time. The state that the system eventually reaches is called the *steady or stationary state*. The steady state has considerable theoretical significance. The steady state is the time-invariant ideal state for an open system. It is the counterpart of the time-invariant ideal state for a closed system—the *equilibrium state* [16].

1.3 Mass-Transfer

1.3.1 Modes of Mass Transfer

The movement of material from one location in solution to another, arises from differences in electrical or chemical potential at the two locations, or from movement of a volume element of solution. The modes of mass transfer are:

- (1) *Migration*. Movement of a charged body under the influence of an electric field (a gradient of electrical potential).

- (2) *Diffusion*. Movement of a species under the influence of a gradient of chemical potential (a concentration gradient).

(3) Convection. Stirring or hydrodynamic transport.

Generally fluid flow occurs because of natural convection (convection caused by density gradients) and forced convection, and may be characterized by *laminar flow* and *turbulent flow* [17].

1.3.2 Laminar And Turbulent Flow Through Tubes

The flow through tubes can be laminar or turbulent, depending on the magnitude of the dimensionless Reynolds number R_e , which is given by

$$R_e = \frac{\rho v d}{\eta} \quad (1.9)$$

where v is the fluid velocity (cm/sec), d is the tube diameter (cm), ρ is the fluid density (g/cm^3), and η is the fluid viscosity (poise). Physically the Reynolds number reflects the relative strength of the inertial and viscous forces in the moving fluid. Inertial forces increase with density and with square of velocity (ρv^2). Viscous forces increase with viscosity and velocity and decrease with the increasing tube diameter ($\eta v/d$). The Reynolds number is simply the ratio of these two terms.

At low Reynolds number ($R_e < 2100$) the flow called *laminar* or *viscous*, above this critical Reynolds number *turbulence flow* occurs [18].

1.4 The Diffusion Process

Mass transport on a molecular scale is called *diffusion* and on the scale of the apparatus is termed *bulk motion*, or more precisely, *convection* [18].

Diffusion is a process which leads to an equalization of concentrations within a single phase [4], in which matter is transported from one part of a system to another as a result of random molecular motions "*random walk*" [19] , or *Brownian motion* [7].

1.4.1 Types of Diffusion Coefficients

Fick studied the interdiffusion of two solutions of the same components but with different concentrations, and also diffusion from a solid into its saturated solution. His diffusion coefficients were therefore *interdiffusion coefficients* or *mutual diffusion coefficients*. In a pure liquid the individual particles move with random motion and have an equal opportunity of taking up any point in the total space occupied by the liquid. If it were possible to label a particle without otherwise changing its properties and to follow its motion through the unlabelled molecules, a *self diffusion coefficient* could be defined. Equally, in a multicomponent system, some labelled molecules of component 1 could be introduced so as to form a concentration gradient of labelled and unlabelled species 1 with an otherwise uniform environ-

ment. Application of Fick's law to the interdiffusion of labelled and unlabelled molecules of species 1 gives the *intradiffusion coefficient* of this species within the solution. A self diffusion coefficient is a special case of intradiffusion in a system which contains only the two distinguishable (by isotopic labelling, nuclear magnetic vector orientation, etc.) forms of the same chemical species. An other related experiment comprises adding at some point in space a trace of labelled species 1 to an otherwise homogeneous medium which may, or may not, be multicomponent, and may, or may not, contain unlabelled species 1. Attention is focused on the decay of the initially formed concentration gradient of the labelled species when it is at low concentration in an otherwise homogeneous environment. The resulting diffusion coefficient is termed a *tracer diffusion coefficient*. Also, the orientation of any molecule undergoes random fluctuations and so rotational motion also occurs. This motion ensures that, in the absence of any external field, the molecules do not take up any preferred orientation. When an external field is present the extent of rotational order produced depends upon a balance between the magnitude and direction of the applied field and the intensity of the rotational motion, which in turn depends upon the thermal energy of the system. Removal of the field causes the degree of rotational order which had been induced to decay back to zero at a characteristic rate, then a *rotational diffusion*

coefficient can be defined. All these diffusion coefficients can be related to the translational Brownian motion of the molecules which arises from random fluctuations in the positions of molecules in space as a result of the decay of a concentration gradient, and a *translational diffusion coefficient* can be termed for this motion [5].

1.4.2 Basic Hypothesis Of Mathematical Theory

Transfer of heat by conduction is also due to random molecular motion, and there is an obvious analogy between the two processes. This was recognized by Fick, who first put diffusion on a quantitative basis by adopting the mathematical equation of heat conduction derived some years earlier by Fourier. The mathematical theory of diffusion in isotropic substances is therefore based on the hypothesis that the rate of transfer of diffusing substance through unit area of a section is proportional to the concentration gradient measured normal to the section,

$$F = - D \frac{\partial C}{\partial X} \quad (1.10)$$

where F is the rate of transfer per unit area of section, C the concentration of diffusing substance, X the space coordinate measured normal to the section, and D is the diffusion coefficient.

In some cases, e.g. diffusion in dilute solutions, D can reasonably be taken as constant, while in other, e.g. diffusion in high polymers, it depends very markedly on concentration. If F , the amount of material diffusing, and C the concentration, are both expressed in terms of the same unit of quantity, e.g. grams or grams molecules, then it is clear from Eq. (1.10) that D is independent of this unit and has dimensions ($\text{length}^2 \text{ time}^{-1}$), e.g. $\text{cm}^2 \cdot \text{sec}^{-1}$. The negative sign in Eq.(1.10) arises because diffusion occurs in the direction opposite to that of increasing concentration.

It must be emphasized that the statement expressed mathematically by Eq. (1.10) is in general consistent only for an isotropic medium, whose structure and diffusion properties in the neighborhood of any point are the same relative to all directions. Because of this symmetry, the flow of diffusing substance at any point is along the normal to the surface of constant concentration through the point. This need not be true in an anisotropic medium for which the diffusion properties depend on the direction in which they are measured [19].

1.5 Differential Equation of Diffusion

The fundamental differential equation of diffusion in an isotropic medium is derived from Eq.(1.10) as follows. Consider an element of volume in the form of a rectangular parallelepiped whose sides are parallel to the axes of coordinates and are of lengths $2dx$, $2dy$, $2dz$. Let the center of the element be at $P(x,y,z)$, where the concentration of diffusion substance is c . Let ABCD and A'B'C'D' be the faces perpendicular to the axes of X as in Fig.(1.3).

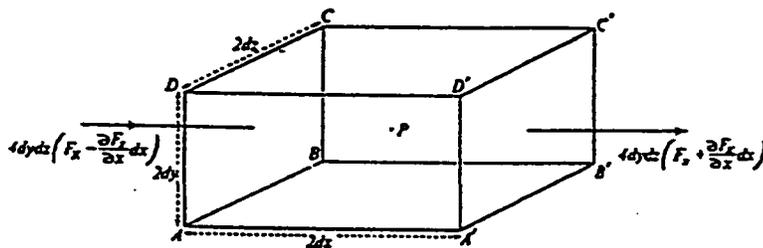


Fig.1.3: Element of Volume

Then the rate at which diffusing substance enters the element through the face ABCD in the plane $(x - dx)$ is given by

$$4 \, dydz \left(F_x - \frac{\partial F_x}{\partial x} dx \right)$$

where F_x is the rate of transfer through unit area of the corresponding plane of through P. Similarly the rate of loss of diffusion substance through the face A'B'C'D' is given by

$$4 \, dydz \left(F_x + \frac{\partial F_x}{\partial x} dx \right)$$

The contribution to the rate of increase of diffusion substance in the element from these two faces is thus equal to

$$- 8 \, dx dy dz \frac{\partial F_x}{\partial x}$$

Similarly from the other faces we obtain

$$- 8 \, dx dy dz \frac{\partial F_y}{\partial y} \quad , \quad - 8 \, dx dy dz \frac{\partial F_z}{\partial z}$$

But the rate at which the amount of diffusing substance in the element increases is also given by

$$\delta \, dx \, dy \, dz \, \frac{\partial C}{\partial t}$$

and hence we have immediately

$$\frac{\partial C}{\partial t} + \frac{\partial F_x}{\partial x} + \frac{\partial F_y}{\partial y} + \frac{\partial F_z}{\partial z} = 0 \quad (1.11)$$

If the diffusion coefficient is constant, F_x, F_y, F_z are given by Eq.(1.10), and Eq.(1.11) becomes

$$\frac{\partial C}{\partial t} = D \left(\frac{\partial^2 C}{\partial x^2} + \frac{\partial^2 C}{\partial y^2} + \frac{\partial^2 C}{\partial z^2} \right) \quad (1.12)$$

reducing simply to

$$\frac{\partial C}{\partial t} = D \frac{\partial^2 C}{\partial x^2} \quad (1.13)$$

if diffusion is one dimensional, i.e. if there is a gradient of concentration only along the X-axis, expressions of Eqs.(1.10 & 1.13) are usually referred to as Fick's first and second laws of diffusion [19,20].

1.6 Solution Methods of the Diffusion Equation when the Diffusion Coefficient is Constant

General solution of the diffusion equation can be obtained for a variety of initial and boundary conditions provided the diffusion coefficient is constant. Such a solution usually has one of two standard forms. Either it is comprised of a series of error function or related integrals, in which case it is most suitable for numerical evaluation at small times, i.e. in the early stages of diffusion, or it is the form of a trigonometrical series which converges most satisfactorily for large values of time. Another solution used the Laplace transform, is essentially an operator method by which both types of solution may be obtained. It is the most powerful of the three, particularly for more complicated problems [19].

1.6.1 Method of Separation of Variables [19,20,21,22,23]

A standard method of obtaining a solution of a partial differential equation is to assume that the variables are separable. Thus we may attempt to find a solution of Eq. (1.13) by putting

$$C(x, t) = F(x) \cdot f(t) \quad (1.14)$$

where $F(x)$ and $f(t)$ are function of x and t respectively .

Therefore

$$\begin{aligned}\frac{\partial C}{\partial t} &= F(x) \cdot f'(t) \\ \frac{\partial C}{\partial x} &= F'(x) \cdot f(t) \\ \frac{\partial^2 C}{\partial x^2} &= F''(x) \cdot f(t)\end{aligned}\tag{1.14 a}$$

substitution of Eq.(1.14 a) into Eq.(1.13) gives

$$F(x) f'(t) = D F''(x) \cdot f(t)\tag{1.14 b}$$

By separation of variables of Eq.(1.14 b) we get

$$\frac{f'(t)}{D f(t)} = \frac{F''(x)}{F(x)}\tag{1.14 c}$$

Since the left and right sides of Eq.(1.14 c) are respectively function of t only and x only, the equation can be satisfied if each side is separately equal to the same constant, let it be $(-k^2)$, and then we have two equations

$$F''(x) = -k^2 F(x)\tag{1.14 d}$$

and

$$f' = -k^2 D f(t) \quad (1.14 e)$$

In the capillary tube diffusion method, which will be discuss in the next chapter, the initial and boundary conditions for a tube close at ($x = 0$) and open at ($x = a$) are

$$\begin{aligned} \text{At } t=0, c=c_0 \text{ for } 0 < x < a, c=0 \text{ for } x > a \\ \text{At } t > 0, c=0 \text{ at } x=a \text{ and } \frac{\partial c}{\partial x}=0 \text{ at } x=0 \end{aligned}$$

To solve Eq.(1.14 d),

$$F''(x) + k^2 F(x) = 0 \quad (1.14 f)$$

the auxiliary equation for a second order linear differential equation is $m^2 + k^2 = 0$. Solving for m we get two complex roots $m = \pm ik$. Then the general solution for Eq. (1.14 f) has the form

$$F(x) = d_1 \cos kx + d_2 \sin kx \quad (1.14 g)$$

To find d_1 , d_2 , and k we use the initial and boundary conditions,

$$F' = (d_1 \sin kx)(-k) + kd_2 \cos kx \quad (1.14 h)$$

but we know that at $x = 0$, $F'(x)=0$ so that $F'(0) = kd_2 = 0$. Therefore $d_2 = 0$, and

$$F(x) = d_1 \cos kx \quad (1.14 i)$$

We also have at $x = a$, $F(a) = 0$, therefore $F(a) = d_1 \cos ka = 0$, but $d_1 \neq 0$ (if so $F(x) = 0$). Then $\cos ka = 0$, and solving for ka with respect to π at which the \cos function equal to zero we get

$$ka = \frac{2n+1}{2}\pi, \quad (1.14 j)$$

$$k = \frac{2n+1}{2a}\pi, \quad n=0,1,2,\dots$$

Then $F(x)$ becomes

$$F_n(x) = d_1 \sum_{n=0}^{\infty} \cos \frac{\pi(2n+1)}{2a}x \quad (1.14 k)$$

To solve the second equation, Eq. (1.14 e),

$$\frac{f'(t)}{f(t)} = -k^2 D$$

by integration

$$\begin{aligned} (\ln f(t))' &= (-k^2 D t + c_1)' \\ \ln f(t) &= -k^2 D t + c_1 \end{aligned}$$

we get for the function $f(t)$,

$$f(t) = c_1 \exp(-k^2 D t) \quad (1.14 l)$$

Substituting Eqs. (1.14 k & 1.14 l) into Eq. (1.14) we get

$$C_n(x, t) = \sum_{n=0}^{\infty} B_n \text{EXP}(-k^2 D t) \cos kx \quad (1.14 m)$$

where $B_n = c_1 d_1$. To find it's value we use the condition
 $c(x, 0) = 0$ at $x > a$ and $c(x, 0) = c_0$ at $0 < x < a$
 putting the value of k from Eq.(1.14 j), Eq.(1.14 m) will be

$$c(x, 0) = \sum_{n=0}^{\infty} B_n \cos \frac{\pi(2n+1)}{2a} x \quad (1.14 n)$$

which can be solved by Fourier analysis as

$$B_n = \frac{2}{a} \int_0^a c_0 \cos \frac{\pi(2n+1)}{2a} x \, dx \quad (1.14 \text{ o})$$

Upon integrating we obtain

$$\begin{aligned} B_n &= \frac{4c_0}{\pi(2n+1)} \sin\left[\frac{\pi(2n+1)}{2} x\right] \\ &= \frac{4c_0}{\pi(2n+1)} (-1)^n \end{aligned} \quad (1.14 \text{ p})$$

Substitution of B_n and k in Eq. (1.14 m) yields

$$c(x, t) = \sum_{n=0}^{\infty} \frac{4c_0(-1)^n}{\pi(2n+1)} \text{EXP}\left[\frac{-\pi^2(2n+1)^2Dt}{4a^2}\right] \cos \frac{\pi(2n+1)}{2a} x \quad (1.14 \text{ q})$$

By integration of Eq. (1.14 q) in the limit of $0 < x < a$ to get the total concentration at any time t we get the final result

$$c(x, t) = \frac{8ac_0}{x} \sum_{n=0}^{\infty} \frac{(-1)^n}{\pi^2(2n+1)^2} \text{EXP}\left[\frac{-\pi^2(2n+1)^2Dt}{4a^2}\right] \sin \frac{\pi(2n+1)x}{2a} \quad (1.15)$$

From Beer's law one can relate the concentration(c) to the absorption of electromagnetic radiation by a compound and to the intensity(I) of the spectrum, so $c(x,t)$ will be equal to $I(t)$ (the intensity of the spectrum at any time), and c_0 will be equal to the I_0 (the initial intensity of the spectrum or the reference one).

In Eq. (1.15), we can replace the ratio of x/a in which x is the total length of the capillary tube (l_t) or the maximum distance that the molecules can go, a is the small length of the tube (l_s) that is under study. Because of the uncertainty in the length of (l_s) which is under the electromagnetic radiation, a ratio R ($R = l_s / l_t$) can be substituted in the equation, we can rewrite Eq. (1.15) as

$$\frac{I(t)}{I(0)} = \frac{8}{\pi^2 R} \sum_{n=0}^{\infty} (-1)^n f_n(t) \sin \left[\frac{(2n+1)\pi R}{2} \right]$$

where (1.16)

$$f_n(t) = \frac{1}{(2n+1)^2} \exp \left\{ \frac{-(2n+1)^2 \pi^2 D t}{4 l_s^2} \right\}$$

A computer program for Eq. (1.16) has been written and it is listed in Appendix 1.

C H A P T E R T W O

A P P A R A T U S

CHAPTER # 2

APPARATUS

2.1 Introduction

Electron Spin Resonance spectroscopy (ESR), also called Electron Paramagnetic Resonance spectroscopy (EPR), discovered by Zavoisky in 1944, has become an essential tool for the study of the structure and dynamics of molecular systems containing one or more unpaired electrons [24].

2.1.1 Systems Studied by ESR

A large number of systems have been studied by ESR:

- (1) Biological system : hemoglobin, nucleic acid, enzymes.
- (2) Chemical system : polymers, catalysts, free radicals.
- (3) Conduction electrons : alkali and alkali earth metals.
- (4) Stable free radicals : in solid and solution, biradicals triplet states.
- (5) Gases : NO , NO₂ , SO₂.
- (6) Irradiated substances : ionic crystals solid and liquid.
- (7) Naturally occurring substance : petroleum, rubber, dolomite, coal.
- (8) Semiconductors : doped semiconductors.
- (9) Transition elements : single crystals, chelates, sandwich compound, alloys [25].

2.1.2 The Resonance Condition in ESR

Electron is a charged particle with angular momentum (orbital and spin) that possesses a magnetic moment, μ_e , given by

$$\mu_e = -g \beta J \quad (2.1)$$

Here, J is the total angular momentum vector in units of \hbar where \hbar is the Planck's constant, g is a dimensionless constant (the g -value, g factor, or spectroscopic splitting factor), and β is a constant called the Bohr magneton. The negative sign in the equation is a consequence of negative electronic charge.

Neglecting orbital angular momentum and considering only the total spin angular momentum, S , Eq.(2.1) can be written as

$$\mu_e = -g \beta S \quad (2.2)$$

The g -value for the free electron, g_e , is 2.0023 for an unpaired electron in an organic molecule or ion, g is generally within a few percent of this value.

The approximation made in Eq.(2.2) is valid for most discussions of the ESR spectra of the organic free radicals and transition metal complexes whose orbital angular momentum can be considered to be "quenched". Treating the g -value as an experimental quantity does not harm the present discussions, since deviations of the g -value from g_e can be accounted for by introduction of spin-orbital coupling.

Magnetic moments can be detected by their interactions with magnetic fields. In zero field, the magnetic moments of unpaired electrons in a sample are randomly oriented. In the presence of a magnetic field H , electron moments assume orientations with respect to the applied field, giving rise to $(2S+1)$ energy states (Zeeman splitting). The measurable components of μ_e are $(g\beta m_s)$, where m_s is the magnetic spin quantum number, which can take the values

$$+S, +(S-1), \dots, -(S-1), -S$$

The application of a magnetic field to an $S = 1/2$ (or larger) system is said to remove the spin degeneracy (i.e., the equal energy values of m_s in the absence of an applied magnetic field).

The energy of an electron moment in a magnetic field is given by

$$E = - \mu_e \cdot H \quad (2.3)$$

Combining eqs.(2.2 & 2.3), the expression

$$E = g \beta H m_s \quad (2.4)$$

results, assuming the direction of the applied field defines the Z-axis. When $S = \frac{1}{2}$ there are two energy levels, whose energy is linearly dependent on H . The separation between these energy levels (Fig. 2.1) at a particular value of the magnetic field, H_z , is

$$\Delta E = g \beta H_z \quad (2.5)$$

In an ESR experiment, an oscillating magnetic field perpendicular to H_0 induces transitions between the $m_s = 1/2$ and $m_s = -1/2$ levels, provided the frequency, ν , is such that the resonance condition

$$\Delta E = h \nu = g \beta H_0 \quad (2.6)$$

is satisfied. In ESR the frequency is held constant and the magnetic field is varied. At a particular value of the magnetic field, H_0 , resonance absorption of energy occurs, resulting in a peak in the spectrum (Fig. 2.1). The frequencies commonly employed in ESR experiments are in the microwave region, these frequencies and magnetic field strengths for g_e resonance absorption signals are given in Table 2.1.

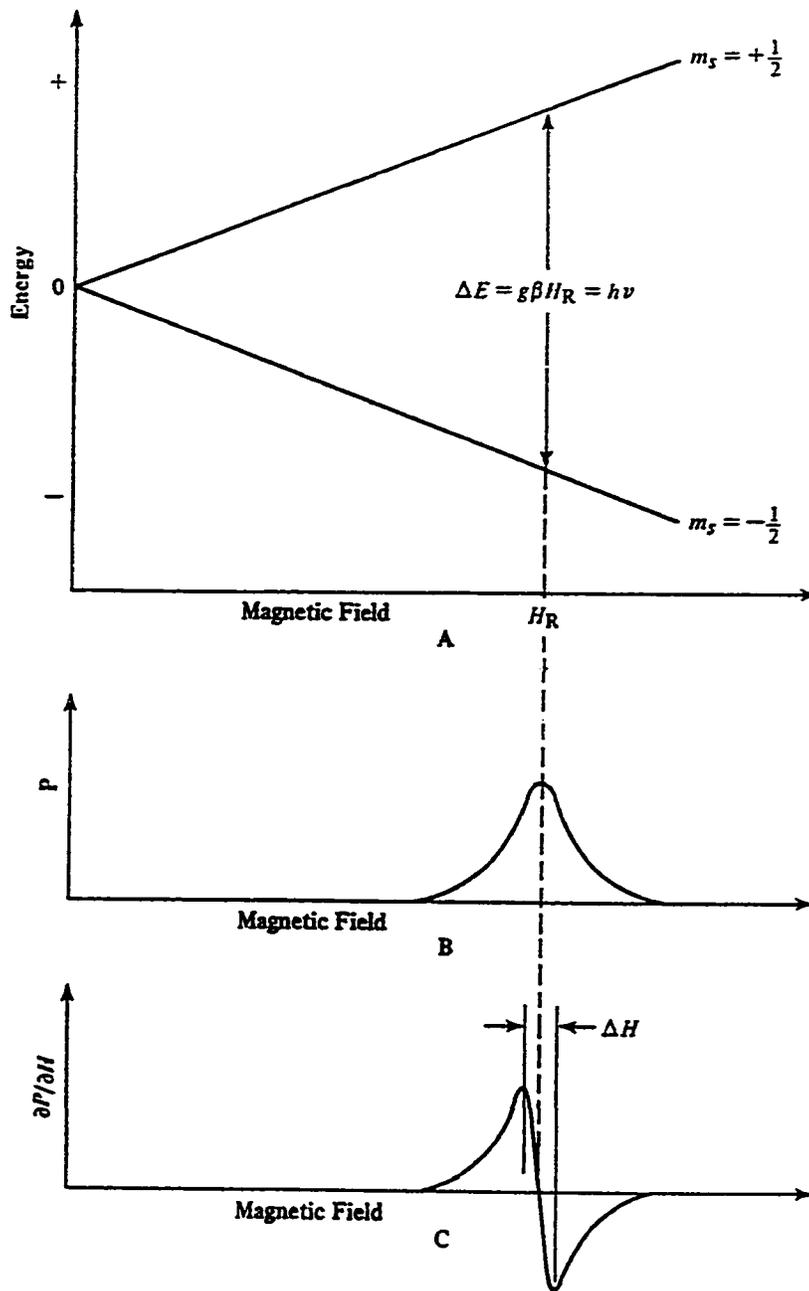


Fig. 2.1 Energy levels and spectra in ESR

Table 2.1 Spectrometer Frequencies and g_e Resonance Field

Designation	ν(Hz)	λ(cm)	ν(cm^{-1})	g_e(gauss)
X-band	$\sim 9.5 \times 10^9$	3.156	0.317	3,390
K-band	$\sim 23 \times 10^9$	1.303	0.767	8,207
Q-band	$\sim 35 \times 10^9$	0.856	1.168	12,489

2.1.3 ESR Instrumentation

The instrumentation used in ESR spectroscopy consists of a source of electromagnetic radiation, a sample holder, and appropriate detection equipment for monitoring the amount of radiation absorbed by the sample. In ESR spectroscopy, a magnetic field provided by an electromagnet is also required. Monochromatic radiation of the various frequencies used in ESR work (Table 2.1) is obtained from *klystrons*, which are electronic oscillators producing microwave energy. Spectrometers operating at X-band (3-cm wavelength) are the ones most commonly employed. The microwave radiation is transmitted along hollow rectangular metal pipes called *waveguides*.

Figure 2.2 gives a block diagram of a simple ESR spectrometer. The sample is placed at the center of the cavity, where the magnetic vector is at a maximum. Quartz tubes (approx. 3mm o.d.) are generally used to contain solid and solution samples. Unlike the NMR technique, the sample tubes are not rotated. The magnetic field is slowly and linearly increased until the resonance condition [Eq.(2.6)] is satisfied, at which point power is absorbed by the sample and a change in current in the detector crystal is monitored. A pair of Helmholtz coils is mounted around the cavity to increase sensitivity. Feeding the coils from an oscillator superimposes a variable amplitude sinusoidal modulation on the slowly varying magnetic field.

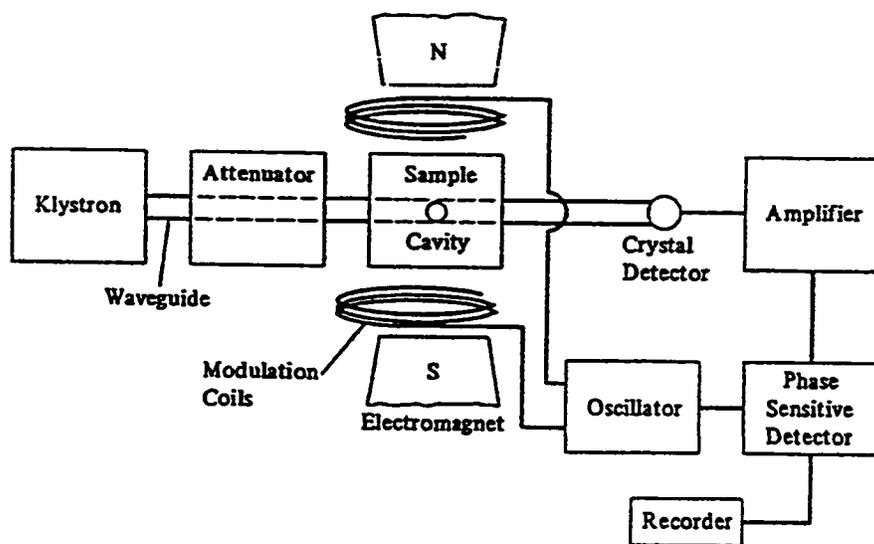


Fig. 2.2 Block diagram of a simple ESR spectrometer

The signal detected by the phase-sensitive detection system is proportional to the slope of the ESR absorption as the magnetic field passes through resonance. The recorder then presents the first-derivative spectrum. Many spectrometers are also equipped to present second-derivative spectra.

The spectral sensitivity of ESR depends on a variety of factors, but with a response time of 1 sec, as few as 10^{11} spins (approx. 10^{-12} moles) can be detected with currently available spectrometers. This sort of sensitivity suggests that ESR spectroscopy would be useful for trace analysis. A minimum detectable concentration is perhaps 10^{-9} M is more reasonable, where for qualitative measurement in aqueous solutions, 10^{-7} M is more reasonable, where for quantitative measurements the sample concentration should be greater than about 10^{-6} M. Concentrations greater than about 10^{-4} M should not be used in order to avoid spin-spin (exchange) broadening. For high-resolution studies, dissolved oxygen should be removed because it is paramagnetic and can contribute to line broadening.

In any ESR experiment, it is important to monitor the microwave frequency at which the spectrometer operates and the magnetic-field range swept during the experiment. Although the frequency is constant during an experiment, the frequency available from a given klystron will vary a little with tuning of the instrument. The frequency can be determined using a built-in frequency counters.

The magnetic field can be monitored using an NMR gaussmeter or by using samples of known g -value, for instance, the DPPH free radical, for which $g = 2.0036$. The magnetic sweep can also be checked using Fremy's radical or oxobis(2,4-pentanedionato)vanadium(IV) [24]. More theoretical and experimental details are given in [25,26,27,28,29].

2.2 VARIAN E-109 ESR SPECTROMETER

2.2.1 General Description

The individual units of the E-109 spectrometer are: An electronic (operator's) console, a microwave bridge, a sample cavity with connecting waveguide, a magnet, and a magnet power supply. The system is completely solid state in design with the exception of reflux klystron tube in the microwave bridge and the cathode ray tube in the oscilloscope. The basic electronic console contains the following modules and assemblies:

- (1) The E-203B Magnetic Field Controller provides direct control and regulation of the magnetic field in the magnetic air gap.
- (2) The E-207 High Frequency Module provides magnetic field modulation frequencies of 25 kHz 100 kHz.
- (3) The E-200 oscilloscope displays the ESR signals for visual observation and the klystron power mode for spectrometer tuning.

(4) The E-080B digital drive recorder graphically displays the ESR spectrum as a function of magnetic field or time. The E-102 microwave bridge provides controlled 9.5 GHz microwave frequency to the cavity and detection of ESR signals. The absorption reference arm design permits operation at microwave power levels to one μW . The E-231 series sample cavity is a multipurpose rectangular cavity with an 11 mm sample access. The cavity is designed to permit up to 40 G peak-to-peak field modulation at the sample.

The magnet used is 9-inch in size with the magnet power supply providing the d.c. power to the magnet. The magnet and microwave bridge are water cooled the chilled water supplied by a mechanically refrigerated closed circuit water chiller.

To measure the field value up to 6 significant figures a Varian E-500 NMR gauss meter is used. Similarly, to get the accurate value of the frequency, HP model 5342A microwave frequency counter is used.

2.2.2 Data Acquisition

2.2.2.1 Introduction

The E-935 Data Acquisition system consists of the following four units:

- (1) HP 9835 computer
- (2) E-935 Data Acquisition System
- (3) HP 1350 Graphics Translator
- (4) HP 1311A CRT Display

With the use of this computer, one can store the experimental data on a cartridge and later retrieve the data for various analyses, like plotting, manipulation of data, and calculations like g-factor, etc. The data collection becomes fast and efficient with the use of computer.

2.2.2.2 E-935 System Software

The software (provided by Varian Instrument Co.) is divided into four functional modules. Only one of these modules can be present in the memory of the HP 9835 computer at a time. The four functions are: calculation, manipulations, plotting, and scanning.

2.2.2.3 Magnet Scan

ESR data may be acquired while the larger magnet is swept. Magnet scan acquires ESR spectra containing 512, 1024, 2048 or 4096 data points, performs multiscan averaging, and continuously displays the current results on the display unit. Scan time may be in the range of 4-17,000 seconds.

2.2.2.4 Data File Capability

An ESR spectrum may be stored as a data file on a tape cartridge along with a file header containing descriptions of the sample. Experiment and the instrumental settings like field, frequency, power, modulation frequency modulation amplitude, number of scans, scan time, time constant, microwave power, etc. are stored in the file.

Such a data file may be retrieved from a tape cartridge, loaded into computer memory, and subsequently displayed on the display unit. Additional operations can be done on the display with the manipulation module. Also a printed catalogue of the files on a given tape cartridge may be obtained. The displayed spectrum can be immediately plotted.

2.2.2.5 Data Manipulation

With the help of the manipulation module , the following manipulations of ESR spectra may be performed using the interactive capabilities of the display unit and key board:

- (1) Base line adjustment of a spectrum.
- (2) Combination of two spectra.
- (3) Detection of erroneous feature from a spectrum.
- (4) Differentiation of a spectrum.
- (5) Digital filtration of a spectrum.
- (6) Double integration of a spectrum.
- (7) X , Y axes scale and shift modifications.
- (8) Editing of the file header and recording the spectrum after editing.

2.2.2.6 Tape Cartridge and File System

The tape cartridge serves two main purposes. First, the system tape cartridge contains all of the programs required for proper system operation. The second purpose is to provide for long or short term off-line storage of ESR data sets.

Before use, all tapes must be initialized as per given procedures. An initialized tape is divided into a total of

852 records of 256 bytes per record. There are two tracks per tape. Records 0-425 on the first track, while records 426-851 are on the second track. All information stored on a tape is saved in form of files. All data are contained in a particular file name. To retrieve any data later, a file name must be read to the computer.

2.3 VACUUM SYSTEM

For meaningful ESR line width studies, the sample should be evacuated in order to get rid of oxygen present in the air which is paramagnetic and can interfere with the line width measurements from the paramagnetic sample of interest.

A normal system manufactured by Pope Scientific Inc. was used with slight modification to suit our needs. One of the modifications is a Quick Fit arrangement to insert and remove the sample easily. Also, to measure the vacuum readily, we used a digital vacuum gauge manufactured by Granville-Phillips. Vacuum could be easily reached as low as 1 mTorr and most of the samples were sealed when the vacuum was of that order.

2.4 Capillary Diffusion Apparatus

This technique first introduced by Anderson and Saddington to study the diffusion of ions in solution by radioactive isotopes [30], the capillary cell is charged with the active solution and immersed in the inactive liquid for a known length of time so that normal diffusion can take place. The cell contents are then transferred to a suitable tray, evaporated to dryness, and the radioactive content measured in standard counting equipment. A reference tray is similarly prepared, a known volume of the active solution which has not undergone the diffusion process being used. The activity of this tray is measured immediately after that of the experimental tray, and from the values so obtained, that fraction of the original cell content which has passed out of it as a result of diffusion can be calculated. Knowing the length of the capillary cell and time of diffusion, the self diffusion coefficient of the ion may be calculated [30]. This method has been extensively used for determining self-and tracer diffusion coefficients of electrolytes [31,32,33]. Other studies on theoretical, concentration gradients effect, stirring rate effect, and temperatures effect have been done on this technique [34,35,36,37,38]. This method is shown to be reliable for ESR study, since it reproduces the values of diffusion coefficients that have been reported by other technique [3,9].

To use the capillary diffusion method with ESR spectrometer, a diffusion cell that can be used under solvent vapor pressure is shown in Fig. 2.3. The capillary diffusion cell was made of 1 mm i.d. heavy wall Pyrex tubing whose open end was ground flat. A small distortion of the capillary is unavoidable in joining the capillary to the 11 mm o.d. tubing (A). The total length of the capillary, l_t , was measured with a micrometer. The arm (D) was used to introduce a solid sample to the reservoir (C). Solvent was vacuum distilled into reservoirs (C) and (B) through the O-ring joint (E). The sample preparation and the experimental procedure will be described in the next chapter.

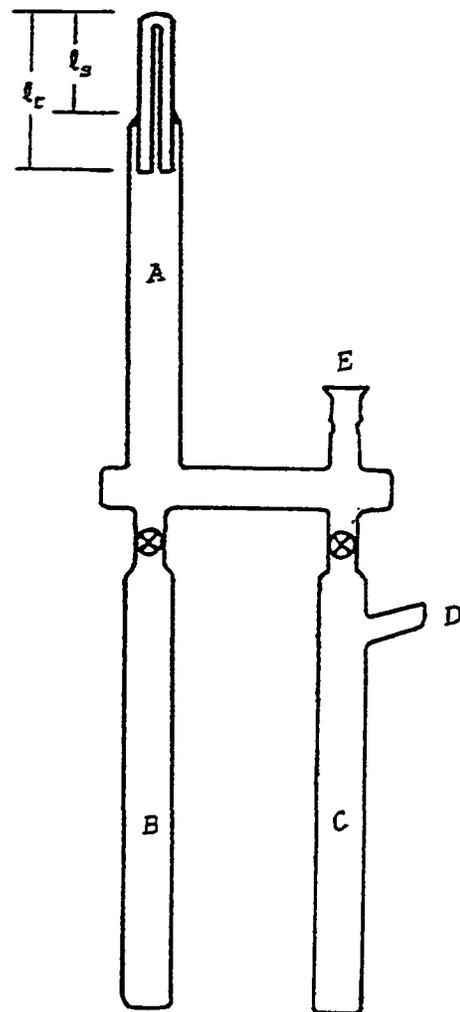


Fig. 2.3 A Capillary Diffusion Apparatus

CHAPTER THREE

PROCEDURE

CHAPTER 3

PROCEDURE

3.1 SAMPLE PREPARATION

3.1.1 VO[5-MeO-Sal SB] (Phen)

This compound was obtained through private communications [39]. The synthesis was carried out following the procedure outlined in reference [40]. The solvent used without further purification is methylene chloride, CH_2Cl_2 , (spectroscopic grade purchased from Fluka-AG). A weighted amount is introduced to the cell, and under vacuum the solvent is added the volume of the solvent measure from the reservoir (C) in the cell which is a 10 ml graduate cylinder, the concentration is about $(9.7 \times 10^{-3} \text{ M})$.

3.1.2 Cp_2VCl_2

This compound is available commercially (PFALTZ & BAUER, INC), the solvent used is the chloroform, CHCl_3 , (spectroscopic grade purchased from ARISTAR BDH). The preparation is similar to the above compound, the concentration is about $(8.7 \times 10^{-3} \text{ M})$.

3.2 Diffusion Cell Preparation

The following steps done to prepare the cell which contains the sample to be used in the ESR spectrometer:

1. Cleaning and drying the cell carefully.
2. Introduction of the weighted sample through the arm (D), then seal (D).
3. Check that all ports stop cock are close in the vacuum line, connect the cell from the O-ring joint to the vacuum line in one port.
4. Prepare the solvent needed (10-15 Ml) in a clean container and connect it to another port in the vacuum line.
5. Switch on the vacuum pump and open the main stop cock of the vacuum line, to let the pressure goes down.
6. Freeze the sample in reservoir (C) and the solvent container by two liquid nitrogen container.
7. Open the stop cock port at the two ports used, and also the stop cock in the cell at (B) and (C).
8. Degassing the cell and the solvent container until a good, low, stable pressure is obtained.
9. Close the main vacuum line stop cock, and the stop cock at (B) in the cell.
10. Remove the liquid nitrogen from the solvent, to start the evaporation of the solvent from the solvent container to condense at the freeze (C) reservoir which contains the sample.

11. One can repeat steps 6,7,8,9,10 to get more solvent in (C) for certain concentration.
12. Close the stop cock at (C) and open the one on (B), then repeat steps 6,7,8,9,10 for (B) , until good amount is collected from the solvent.
13. Close the stop cock at the solvent port, freeze (B) and (c), then open the stop cock at (B), (C), and the main vacuum line stop cock to remove the vapor of the solvent in the system.
14. Seal below the O-ring joint carefully, then close the two stop cock at (B) and (C) .
15. Close the main stop cock of, switch off the vacuum pump, and release the vacuum system slowly from any unused port to equalize the pressure.
16. Leave the cell to come to room temperature, then check the cell slowly to dissolve the sample with the solvent in (C).
17. Handle the cell pointing the capillary downward, then open slowly the stop cock at (C) to transfer the solution into the capillary through (A).
18. If a bubble of the vapor block the solution to inter the capillary, remove it by rubbing the capillary with fingers or small cloth slowly and carefully.
19. The excess liquid is poured back into (c), and the cell gently shaken to disperse any paramagnetic solution coating the wall of (a).

20. Put the cell in the ESR cavity at the close end to set up the parameters for the ESR spectrometer, use collars available to set the length of the capillary in the cavity and fix it to the cell by tape.
21. Maximum lines height should be obtained from the ESR parameter setting because these lines will decrease by time, so by the end of the experiment it will be readable.
22. Remove the cell from the cavity, open the stop cock at (B) slowly to transfer the solvent above the capillary cell in (A), put back the cell in the cavity, record the time, and take the first spectrum run.
23. The cell should not disturb in the cavity, three to four runs should be taken every day, with the same setting of the ESR spectrum parameter. Each run record the time and find the difference of time of this run and the time when the experiment start in seconds.
24. Analyze the spectrum of each run by finding the intensity of each line by the computer manipulation program, the first run can be taken as a reference, find the total intensity of the reference spectrum, enter these values into a data file for the computer program (in the appendix the program and the data files are listed for the two compounds) called Diff.
25. The data are collected for three to four weeks.

26. Run the program at the end will give two output files one to calculate the diffusion coefficient, and the other you can plot the total intensity as a function of time from the output file plot generated.

3.3 ESR Spectrometer Operation

The following is a brief description of how to operate the ESR spectrometer:

- (i). Put on the chiller water for circulation .
- (ii). Turn the console power on and bring the field set slowly to the central field value ($\approx 3300 - 3400$ G), through the zero value.
- (iii). In the microwave bridge the STANDBY setting to TUNE position, while keeping the microwave power around 40 Db.
- (iv). Adjust the iris and turn on the frequency counter.

To record the spectra the capillary must be at the active region of the cavity, this could be adjusted with the collars provided. Proper tuning procedure should be followed on the microwave bridge.

Keep the field set at 3300 G, scan range of 1000 G for the room temperature spectra. For good line shape analysis, the spectra taken must be free from artifact broadening due to the wrong selection of spectrometer parameters like scan speed,

time constant, microwave power, modulation amplitude, etc. To get the right signal, first increase the modulation amplitude until no peak no peak distortion or line width increase is noticed. This value of modulation amplitude could be used for the experiments subsequently. Second, increase the microwave power until the saturation of the peak occurs. Once the right spectra are plotted, then they are stored to the tape for Varian Spectrometer, and plotted also by the plotter, the ESR parameters for the compounds are tabulated in Table 3.1.

Table 3.1 The ESR Parameters For Experiments

Parameters	Compound I	Compound II
Number of Point	2048	2048
Number of Scans	4	1
Field Set (G)	3440.73	3410.00
Scan Range (G)	1000	800
Scan Time (sec)	120	180
Modulation Frequency (kHz)	100	100
Modulation Amplitude(G)	2	2
Receiver gain	1.25×10^4	4.0×10^4
DC Gain	Normal	Normal
Time Constant	0.25 s	0.25 s
Microwave Power	2 mW	20 mW
Microwave Frequency (GHz)	9.489787	9.51050

Compound I : VO[5-MeO-Sal SB] (Phen)

Compound II : Cp_2VCl_2

CHAPTER FOUR

ESR MAGNETIC PARAMETERS

CHAPTER #4

ESR MAGNETIC PARAMETERS

4.1 Introduction

The analysis of the spectra to extract the magnetic parameters of the complex is obtained from the rigid limit and room temperature spectra according to Wilson and Kivelson's method [41]. The parameters are necessary for a proper interpretation of the observed ESR spectra at various temperatures in terms of correlation time τ_1 and anisotropic rotational diffusion (N). The ESR spectra of Cp_2VCl_2 in chloroform ($CHCl_3$) at various temperatures, in the motional narrowing region, is analyzed to study the anisotropy of molecular rotation (N) defined as the ratio R_{\parallel}/R_{\perp} where R_{\parallel} is the rotational diffusion constant along the molecular principal axis of rotation, and R_{\perp} is the rotational diffusion constant perpendicular to the molecular principal axis of rotation [42,43,44]. The results are compared with the anisotropy of rotation (N) calculated from Stokes-Einstein model [11] and that calculated from the Allowed-Values Equation (AVE) [45]. This comparison demonstrates the applicability of these narrowing motion region equations to experimentally observable system as well as giving some valuable information about the structure. From the analysis of the temperature dependent ESR

line shapes in the motional-narrowing region the stickiness factor (S) which is a measure of all non-hard-core interactions is determined [15,46,47]. Also we will show how the dependence of line widths on temperature is utilized to obtain the anisotropic factor (N) and the correlation time τ_c . These experimental N values are compared with theoretical values [11,45].

The vanadyl radical has an electron spin of half (unpaired electron, d^1 -system) which interacts with the vanadyl nucleus which has the spin $I=7/2$ therefore the ESR spectrum consist of $(2I+1)$, that is eight lines.

4.2 Spin Hamiltonian and Definition of the Parameters

The appropriate spin Hamiltonian, H_s , required to describe the vanadyl paramagnetic systems in general is [48,49],

$$H_s = B \beta_0 g \cdot S + S \cdot A \cdot I + (\text{electron-electron dipole term}) + (\text{electron-electron exchange term}) + (\text{nuclear Zeeman term})$$

The first term is the electron Zeeman term and the second term is the electron-nuclear hyperfine interaction term, where B is the laboratory magnetic field, β_0 is electron Bohr magneton, g the electron g -value tensor, S is the electron spin angular

momentum operator, A is the electron-nuclear hyperfine tensor, and I is the nuclear spin angular momentum operator. The nuclear Zeeman term is not important in most work [50] and will be omitted.

The electron-electron dipole and exchange terms are significant only when very high concentration [51] and single crystals studies [52] are considered, so our Hamiltonian will be,

$$H_s = B\beta_0 g \cdot S + S \cdot A \cdot I \quad (4.1)$$

4.2.1 ESR Spectra in Solution

The ESR spectra of the vanadyl complexes in liquid solution can be described by the isotropic electron Zeeman and electric-nuclear hyperfine interactions included in the spin Hamiltonian [41,53] as in eq.(4.1), the allowed transitions when microwave field at frequency ω_0 is applied, occur at the field B which satisfies the resonance condition [41,53]

$$\omega_0 = g_0 \beta_0 \frac{B}{h} + A_0 M + \frac{1}{2} A_0^2 \left[\frac{I(I+1) - M^2}{g_0 \beta_0 (B/h)} \right] \quad (4.2)$$

where ω_0 is the microwave frequency in radians per second, h is the Planck's constant divided by 2π , $I = 7/2$, M is the

nuclear spin magnetic moment, B is the resonant value of the magnetic field corresponding to M ,

A_0 is the isotropic hyperfine constant in radians per second,

$$A_0 = \frac{1}{3} (A_x + A_y + A_z) \quad (4.3)$$

and g_0 is the isotropic g-factor,

$$g_0 = \frac{1}{3} (g_x + g_y + g_z) \quad (4.4)$$

where X, Y, and Z refer to the molecule fixed coordinates, and it has been assumed that g and A can be simultaneously diagonalized.

The contact hyperfine constant A_0 can be obtained to second order by taking the difference of B_M for line M and B_{-M} for line $-M$

$$A_0 = \frac{g_0 \beta_0 (B_M - B_{-M})}{2Mh} \quad (4.5)$$

furthermore,

$$g_o - g_s = g_s \left[\frac{B_s - \frac{1}{2}(B_M + B_{-M})}{\frac{1}{2}(B_M + B_{-M})} \right] - \frac{2A_o^2 h^2 [I(I+1) - M^2]}{g_s \beta_o^2 (B_M + B_{-M})^2} \quad (4.6)$$

where g_s and B_s are the isotropic g -value and the resonant value of the magnetic field (in Gauss), for standard of known g -value. For our system, the standard used is diphenylpicrylhydrazyl (DPPH), with g -value equal to 2.0037.

The first term in eq.(4.1) the electron Zeeman term, represents the large interaction of the electron spin with the laboratory magnetic field. This term yields the relation,

$$h \nu = g_o \beta_o B \quad (4.7)$$

where h is Planck's constant and ν is the microwave frequency of the ESR spectrometer. This equation can be reduced to [26],

$$\nu \text{ (in GHz)} = 2.80247 (g_o / g_e) B \text{ (in kG)} \quad (4.8)$$

For vanadyl complexes eq.(4.8) can be further reduced to,

$$\nu \text{ (in GHz)} = 2.847 \text{ B (in kG)} \quad (4.9)$$

This relation serves as a useful conversion factor. Most ESR spectrometers operate at $\nu = 9.5$ GHz (X-band), and thus $B = 3.3$ kG. The unpaired electron responsible for the ESR spectrum is confined largely to the vanadyl center, which interacts with the nuclear spin of ^{51}V ($I = 7/2$), the result is $(2I + 1)$ or 8 lines separated by a coupling constant with different intensities. This nuclear interaction is represented by the second term in eq.(4.1). It is apparent from the variation in the intensities of the lines that the line widths of the different hyperfine components are not equal and depend on the ^{51}V nuclear spin quantum number [41,53].

This line width variation indicates that the anisotropic hyperfine and Zeeman interaction are not completely average out by the tumbling of the molecules in the liquid, and they contribute significantly to the lines width. The relation of the important A_0 and g_0 parameters are shown in (Fig. 4.1) along with the lines widths (ΔH) and the peak heights (h) in a first derivative representation of the absorption curve. The spectra can be recorded with the opposite phase from that of Fig. 4.1 producing an upside down effect. The presentation most consistent with current trends in spectroscopy involves positive phase and increasing B from left to right.

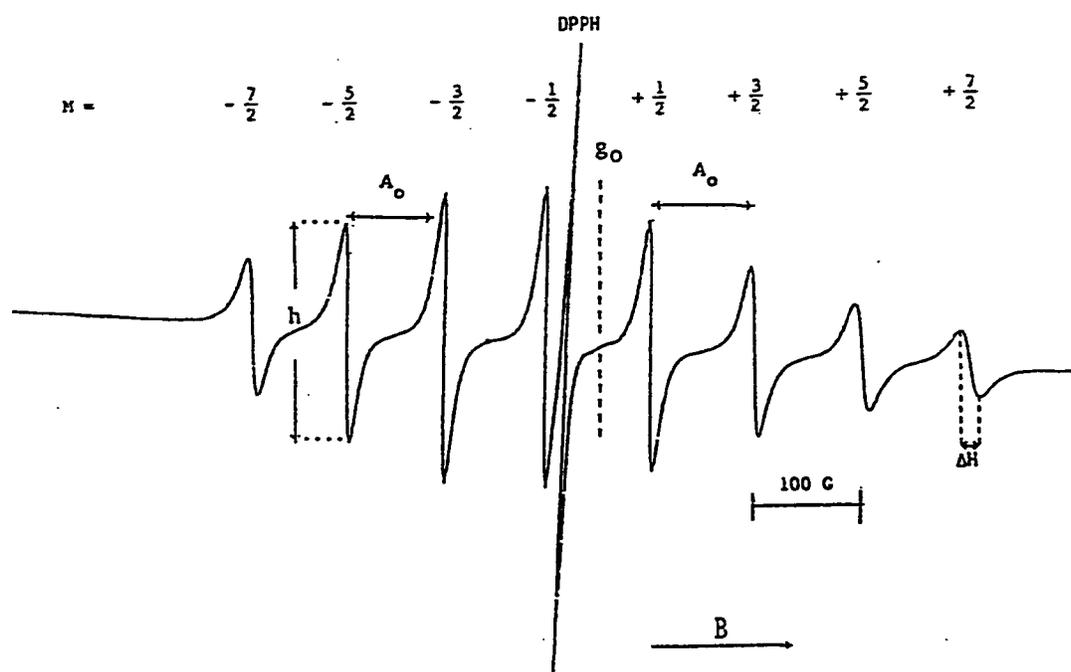


Fig. 4.1: A Typical X-band Spectrum Of Oxovanadium (IV) in Solution at Room Temperature.

4.2.2 ESR Spectra in the Absence of Molecular Motions

4.2.2.1 Anisotropic Effects

The glass spectrum of vanadyl probes can be described in terms of the Zeeman and nuclear-electronic hyperfine interactions by the spin Hamiltonian [41,49,54],

$$H_s = \beta_o (g_x B_x S_x + g_y B_y S_y + g_z B_z S_z) + A_x S_x I_x + A_y S_y I_y + A_z S_z I_z \quad (4.10)$$

a special case of eq.(4.1), both g and A depend on the orientation of the magnetic field in terms of molecular coordinates and the direction of B . The extremes in g and A lie along the principal X, Y, and Z axes. The glass ESR spectrum of a typical VO^{2+} at (77 k) is illustrated in Fig. 4.2 this glass spectrum of vanadyl complexes is governed by [41],

$$\omega_o = \frac{g\beta_o B_M}{h} + AM + \frac{(A_x^2 + A_y^2)(A_x^2 + A_z^2)}{8A^2} \frac{[I(I+1) - M^2]}{g\beta_o B_M/h} \quad (4.11)$$

where

$$g = (g_x^2 \cos^2 \theta + g_x^2 \sin^2 \theta \cos^2 \phi + g_y^2 \sin^2 \theta \sin^2 \phi)^{\frac{1}{2}}$$

and

$$gA = (A_x^2 g_x^2 \cos^2 \theta + A_x^2 g_x^2 \sin^2 \theta \cos^2 \phi + A_y^2 g_y^2 \sin^2 \theta \sin^2 \phi)^{\frac{1}{2}}$$

θ is the angle between B and unique molecular axes, and ϕ is the azimuthal angle. Note that the line for which $\theta = 0$ gives rise to the spectrum corresponding to g_x and A_x in (Fig. 4.2) while those for which $\theta = \phi/2$, $\phi = 0$ and $\pi/2$ give rise to the "doubled spectra" corresponding to (g_x, A_x) and (g_y, A_y) as in Fig. 4.2. The anisotropic tensors can be determined from ridged limit spectra, Fig. 4.2, as Wilson and Kivelson have done, for vanadyl acetylacetonate VOAA [41]. The spacings I, II, III, IV, and V in Fig. 4.2 correspond to,

$$A_x = \frac{II + III}{2} \quad (4.12)$$

$$A_x + A_y = IV + V \quad (4.13)$$

$$I = \frac{\beta_o B 2\delta_g}{h} + 4C \left(\frac{7}{2}\right) \quad (4.14)$$

where

$$\delta_g = \frac{1}{2} (g_x - g_y) \quad (4.15)$$

$$c = \frac{1}{4} (A_x - A_y) \quad (4.16)$$

although the values of g_x , g_y , g_z and A_x , A_y , A_z obtained from the glassy spectrum are compatible with the isotropic values obtained from the liquid spectrum, the most accurate determinations of the magnetic parameters are obtained by measuring g_0 and A_0 from the liquid spectrum, g_z , $g_x - g_y$, A_z and $A_x - A_y$ from the glass spectrum, and g_x , g_y , and A_x , A_y from this data and eqs. (4.14 and 4.15) [41].

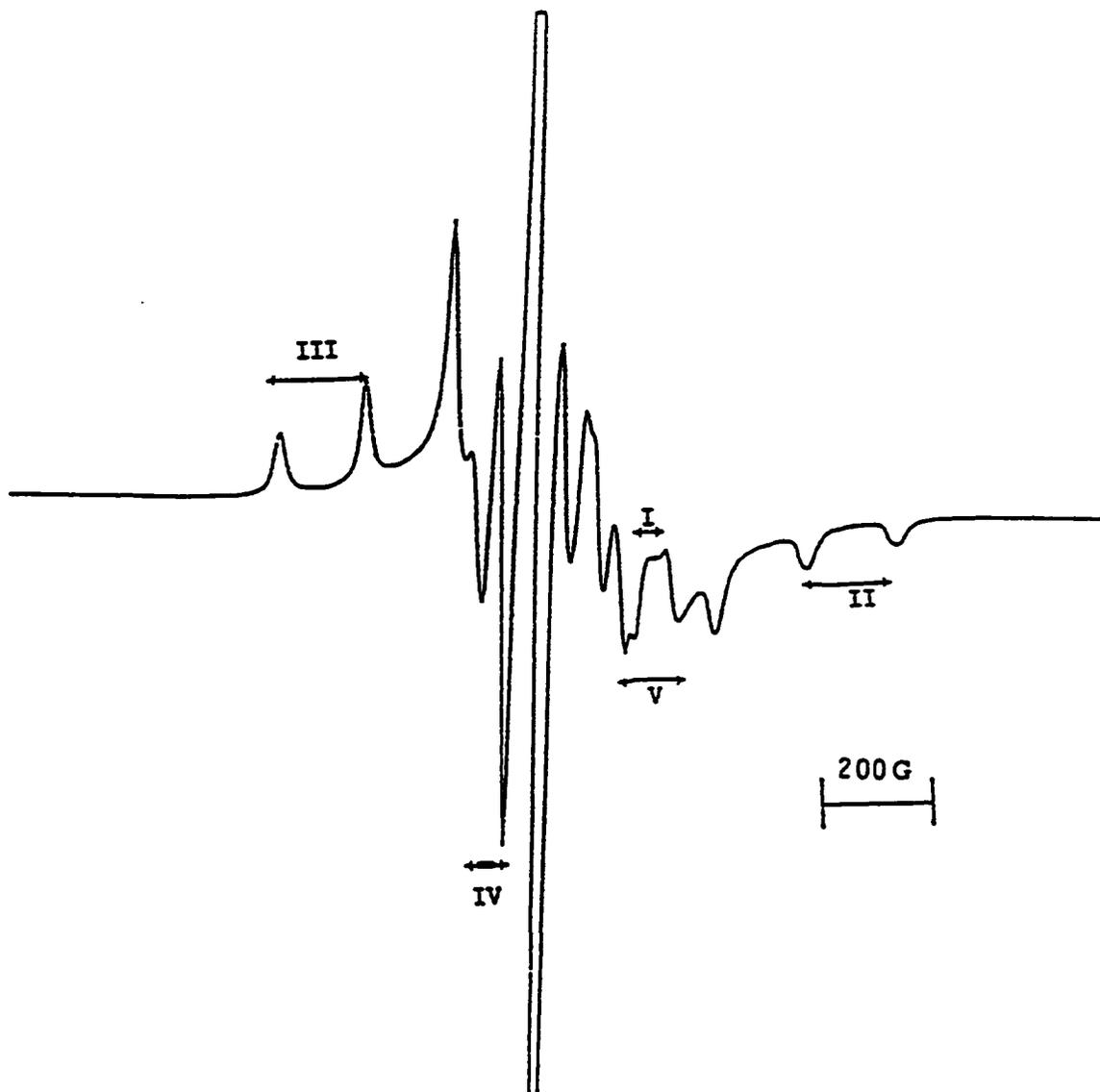


Fig. 4.2 Ridgid Limit spectra at 77 k

4.2.2.2 The Effects of Rotational Motion on ESR Spectra

The shape of an EPR spectrum is markedly dependent upon the rotational mobility of the paramagnetic spin probe. For example, in a nonviscous solvent nearly all vanadyl probes exhibit eight sharp lines. The rapid isotropic tumbling motion averages out all anisotropic effects discussed above, therefore sharp lines are obtained. If the rotational motion is slowed down by increasing the solvent viscosity, averaging is incomplete. The result is unequal broadening of the eight lines. Molecular motion is not necessarily isotropic. Preferential motion about one axis is likely to occur for any asymmetric spin probe, but the effects are particularly evident when surrounding medium is anisotropic. The ESR spectrum is very sensitive to the probe mobility. The best way to make use of this sensitivity is to describe the motion quantitatively in terms of the shape of the spectrum.

As discussed above, both the anisotropic hyperfine electron-nuclear interaction and the anisotropic spin-orbit interaction in a spin probe depend on the mutual orientation of the external magnetic field and of π -orbital of the unpaired electron. Rotations of the probe modulate these interactions, leading to fluctuations in local magnetic fields and changes of the ESR line widths. Rotational motion is characterized by correlation time (τ_R) which determines the effective frequency of rotation, f_{eff} (in Hz) = $(2\pi\tau_R)^{-1}$.

The principal g and A values are related to the isotropic g_0 and A_0 values through eqs (4.3 and 4.4).

The anisotropy of g and A tensors is applied in dynamic spin probe studies to get information about the motional state of the probe. Between the two limits of very rapid motion and a rigid glass, the spectra are quite complex. This complexity provides much of the information that can be obtained about molecular motion using spin probe techniques.

4.3 The Correlation Time Group

The correlation time data of vanadyl complexes usually divided into three groups,

Rapid rotations	10^{-11}	-	10^{-9}	sec.
Slow rotations	10^{-9}	-	10^{-6}	sec.
Very Slow rotations	10^{-6}	-	10^{-3}	sec.

The limits of these regions are essentially determined by the anisotropy of the magnetic interactions occurring in the probes.

4.4 Molecular Tumbling

The spin Hamiltonian can be divided into,

$$H = H_0 + H_1(\Omega) + e(t) \quad (4.17)$$

where

$$H_0 = \beta_0 g_0 B S_z + A_0 I \cdot S \quad (4.18)$$

where S_z is the electron spin along the Z-axis, S and I are the electron and the nuclear spin operator, B is the laboratory magnetic field, and β_0 is electron Bohr magneton, and H_1 is equal to,

$$H_1 = \beta_0 B \cdot g S + I \cdot A \cdot S \quad (4.19)$$

where g and A contain the an isotropic components of the electron g -tensor and electron-nuclear dipole hyperfine tensor after the isotropic components g_0 and A_0 have been subtracted. These tensors are defined in terms of molecular fixed axis whose orientation varies with time, with respect to the laboratory magnetic field due to molecular motion.

Thus the time dependence of H_1 is modulated by the reorientational motion of the molecule (Ω), and

$$h\epsilon(t) = \frac{1}{2} h \gamma_e B [S_+ \exp(-i\omega t) + S_- \exp(i\omega t)] \quad (4.20)$$

where

$$S_+ = S_x + iS_y$$

is the raising operator,

$$S_- = S_x - iS_y$$

is the lowering operator, γ_e is the magnetogyric ratio of a free electron and ω is the microwave frequency. $\epsilon(t)$ gives the interaction of the spin with the microwave radiation field which must be included in the Hamiltonian if the saturation effects are to be considered. Saturation effects are known to be negligible for experimental spectra discussed, only terms linear in $\epsilon(t)$ must be considered [54]. The term "fast" and "slow" refers to the randomly fluctuating part $H_1(t)$ spin Hamiltonian (where $H_1(t)$ has a time averaged value zero) in relation to a correlation time τ_f which characterizes the underlying stochastic process.

Fast motional theories are applicable if $|H_1(t)|\tau_p \ll 1$ where $H_1(T)$ is expressed in frequency units. On the other hand, the slow tumbling is attained when $|H_1(t)|\tau_p \gg 1$. Since the slow and the very slow rotations lie outside the scope of this work, only the fast tumbling rotations will be discussed.

4.5 Fast Tumbling

It is convenient at this point to define a line width parameter T_2 . If the peak to peak width is ΔH (in gauss), then $1/T_2 = \Delta H$. In general, the dependence of T_2 upon M , the component of the nuclear spin along the direction of the applied magnetic field is given by,

$$T_2^{-1}(M) = A + BM + CM^2 + DM^3 \quad (4.21)$$

where A , B , C , and D depend on the magnetic tensors. The line width contributions can be calculated with the Redfield relaxation matrix theory as applied by Freed and Fraenkel [55,56] or with the theory of Kubo and Tomita as applied by Kivelson [57].

In order to determine the anisotropy of rotation (N) and the axis of rotation, the motional narrowing analysis of Freed *et al.* [42,43,44] has been used and eq.(5) in reference [44] has been used.

In practice τ_p is obtained from B , C , and D coefficients because of the uncertainty in A [51]. The theoretical expressions for B and C coefficients will be discussed in the next chapter [42,43].

For isotropic Brownian $\tau_p = (6R)^{-1}$ where R is the rotational diffusion coefficient. For the eight lines spectrum for vanadyl probe, eq.(4.21) will give rise to four experimentally determined unknowns, A , B , C , and D . Since the line width coefficients B and C are function of τ_p values, in the case of two values of τ_p value not agreeing, the assumption of anisotropic rotation must be suspected. Fortunately, this restriction has been resolved by Freed and his co-worker [43,58]. The cubic dependence of $(1/T_2)$ on M given in eq.(4.21) is retained, but the motion is described by a rotational tensor R . This is often axially symmetric and is defined by two components R_{\parallel} and R_{\perp} , where R_{\parallel} is the rotational diffusion constant along the principal axis of rotation, and R_{\perp} is the rotation diffusion constant perpendicular to the principal axis of rotation. The anisotropy of rotation (N) is defined as the ratio of R_{\parallel}/R_{\perp} .

Two related correlation times $\tau(0)$ and $\tau(2)$ are given by,

$$\tau(0) = (6R_{\perp})^{-1} \quad (4.22)$$

$$\tau(2) = (2R_{\perp} + 4R_{\parallel})^{-1} \quad (4.23)$$

Reference should be made to the original literature [43,58] and more recent work [59] for details of this analysis.

CHAPTER FIVE

RESULTS AND DISCUSSION

CHAPTER # 5

RESULTS AND DISCUSSION

5.1 DETERMINATION OF THE MAGNETIC PARAMETERS

5.1.1 Isotropic Magnetic Parameters

The spectrum of Cp_2VCl_2 is shown in (Fig.5.1). The isotropic hyperfine constants, A_0 , and g_0 were obtained from B_M for line M and B_{-M} for line $-M$ by eq.(4.5) and eq.(4.6). A_0 and B_0 were determined for each pair of M and $-M$ lines and average over all pairs.

5.1.2 Anisotropic Magnetic Parameters

The rigid limit of Cp_2VCl_2 is shown in (Fig.5.2). The anisotropic tensors were determined from this rigid limit spectrum using second - order perturbation theory , and from the isotropic A_0 and g_0 determined from the liquid spectra, as Wilson and Kivelson have done for VOAA [41]. Spacings II and III , from the rigid limit spectrum are used in eq.(4.12) to determine A_z

$$A_z \text{ (G)} = (\text{ II} + \text{ III}) / 2$$

as the resonant lines for the magnetic field parallel to Z-axis ($H \parallel Z$) are concentric around g_z value. From the two terminal resonant lines in the glass spectrum, (Fig.5.2), which corresponds to $H \parallel Z$ ($B_{7/2}$ and $B_{7/2}$) value of g_z can be obtained as follows,

$$g_z = \frac{h\nu}{\beta_0} / \left[\frac{B_{7/2} - B_{7/2}}{2} \right] \quad (5.1)$$

where h is Planck's constant, ν is the microwave frequency in Hz, β_0 is the Bohr magneton. combining the values of A_x and g_z obtained from rigid limit spectrum with those of A_0 and g_0 obtained from the liquid spectra in eq.(4.3) and eq.(4.4), values of $A_x + A_y$ and also $g_x + g_y$ were obtained as follows,

$$A_x + A_y = 3A_0 - A_x \quad (5.2)$$

$$g_x + g_y = 3g_0 - g_x \quad (5.3)$$

the rigid limit resonant lines for $H \parallel Z$ are governed by,

$$\omega_0 = \frac{g_z \beta_0 B_M}{h} + A_x M + \frac{(A_x^2 + A_y^2) \left(\frac{63}{4} - M^2 \right)}{4 g_z \beta_0 B_M / h} \quad (5.4)$$

which is the case in eq.(4.11) where $\theta = 0$. Substituting A_y with respect to A_x from eq.(5.2) in eq.(5.4) will yields A_x value which can be used in eq.(5.2) again to get A_y . Using eq.(5.4) and resonant lines for $H // Z$, the tensors obtained above were further adjusted. Spacing I is

$$I \text{ (Hz)} = \left[\frac{\beta_o B (g_x - g_y)}{h} \right] + \frac{7}{2} (A_x - A_y) \quad (5.5)$$

where B (in gauss) is calculated from the microwave frequency (in Hz) and g_o by eq.(4.7), spacing I measured from the rigid limit (Fig.5.2) in gauss, should be used in eq.(5.5) in Hz. substituting A_x and A_y values obtained from eq.(4.5) and values of g_y from eq.(5.3) the following expression obtained,

$$I \text{ (Hz)} = \left[\frac{\beta_o B (2g_x + g_x - 3g_o)}{h} \right] + \frac{7}{2} (A_x - A_y) \quad (5.6)$$

which is solved to give the g_x value and then g_y is also obtained using eq.(5.3).

Again when $\theta = \pi/2$ and $\phi = 0$ and $\pi/2$ in eq.(4.11) yield the following two equations which govern the resonant lines for $H \parallel X$ -axis and $H \parallel Y$ -axis respectively,

$$\omega_o = \frac{g_x \beta_o B_H}{h} + A_x M + \frac{(A_x^2 + A_y^2) (A_x^2 + A_z^2) (\frac{63}{4} - M^2)}{8 A_x^2 g_x \beta_o B_H / h} \quad (5.7)$$

$$\omega_o = \frac{g_y \beta_o B_H}{h} + A_y M + \frac{(A_x^2 + A_y^2) (A_y^2 + A_z^2) (\frac{63}{4} - M^2)}{8 A_y^2 g_y \beta_o B_H / h} \quad (5.8)$$

The above two equations are used to verify the calculated values of the g and A tensors.

Several magnetic parameter values for Cp_2VCl_2 have been reported [62,65,66,67], these values are tabulated in table 5.1, the experimental result obtained from temperature variation and the line width measurements [13], will be analyzed by using the reported values to find the best reported value, from which we determine the anisotropy of rotation N .

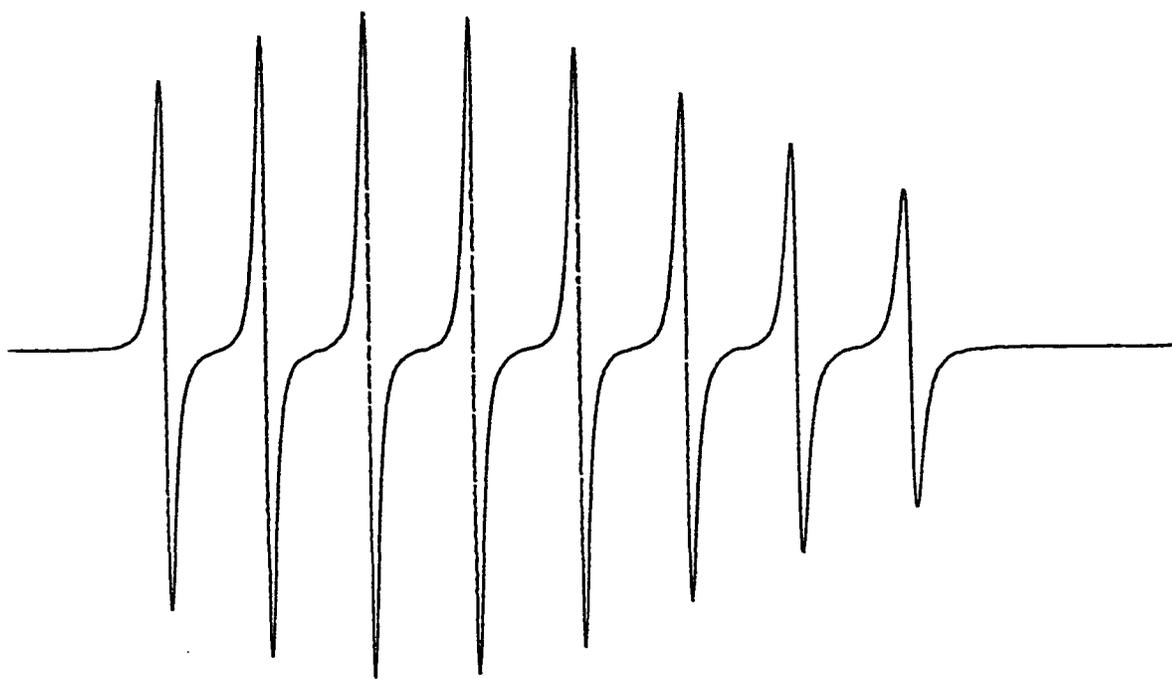


Fig. 5.1: X-band Spectrum of Cp_2VCl_2 at room temperature.

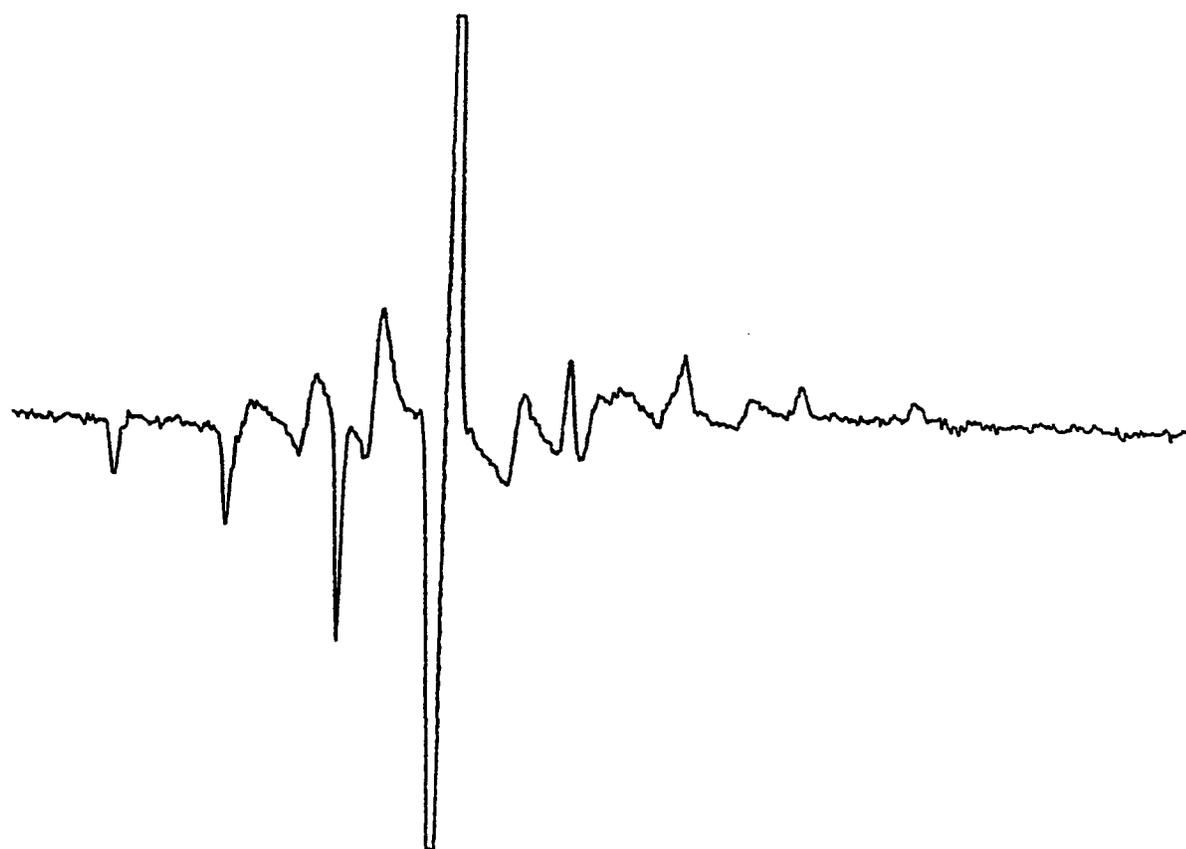


Fig 5.2: X-band Spectrum of Cp_2VCl_2 at $T = 77 \text{ K}$.

Table 5.1: Reported Magnetic Parameters for Cp_2VCl_2

Mag.Par in (G)	(a)	(b)	(c)	(d)
g_x	1.986	1.965	1.986	1.9802
g_y	1.971	1.946	1.971	1.9695
g_z	2.000	1.967	2.010	2.0013
g_0	1.98567	1.95933	1.98900	1.98367
A_x	79.808	81.974	76.444	80.600
A_y	126.183	130.065	123.818	125.500
A_z	16.177	29.511	20.457	20.600
A_0	74.05611	79.44862	75.13467	75.56670

(a) C. P. Stewart and A. L. Porte [66], Cp_2VCl_2 in CH_2Cl_2 + toluene.

(b) D. P. Bakalik and R. G. Hayes [67], Cp_2VCl_2 in CH_2Cl_2 + toluene.

(c) A. T. Casey and J. B. Raynor [65], Cp_2VCl_2 in toluene.

(d) J. L. Petersen and L. F. Dahl [62], Single-Crystal EPR Study of $\text{V}(\eta^5\text{-C}_5\text{H}_4\text{CH}_3)_2\text{Cl}_2$ Doped in $\text{Ti}(\eta^5\text{-C}_5\text{H}_4\text{CH}_3)_2\text{Cl}_2$.

5.2 MOTIONAL NARROWING ANALYSIS

The spectra in Fig. 5.3 are those of Cp_2VCl_2 in toluene taken over a temperature range (294 - 178 K). This free radical has an electron spin (1/2) which interacts with the vanadyl nucleus which has spin $I = 7/2$, therefore the ESR spectrum consists of eight lines each corresponding to a nuclear spin magnetic quantum number M :

$$M = I, I-1, \dots, -I$$

The line width parameter, T_2 (in gauss⁻¹), of each of these lorentzian hyperfine lines can be described by an expansion in M , the component of the nuclear spin along the direction of the applied magnetic field given by eq.(4.21).

The line widths and the peak to peak heights are obtained experimentally utilizing the computer graphic display which facilitates on screen line width measurements [13]. After measuring the line widths of each of these line, the widths are then entered into the computer program written for eq(4.21) for polynomial regression (called POLREG in Watfive computer language) where the coefficients of the best polynomial fit are output for the experimental values by using least square fitting, since our fitting is cubic four coefficients will be produce $A(0)=A$, $A(1)=B$, $A(2)=C$, and $A(3)=D$. The B and C values for different temperatures are shown in (Table 5.2) with the lines widths, also (Fig.5.4) shows the B Vs C plot of the experimental values.

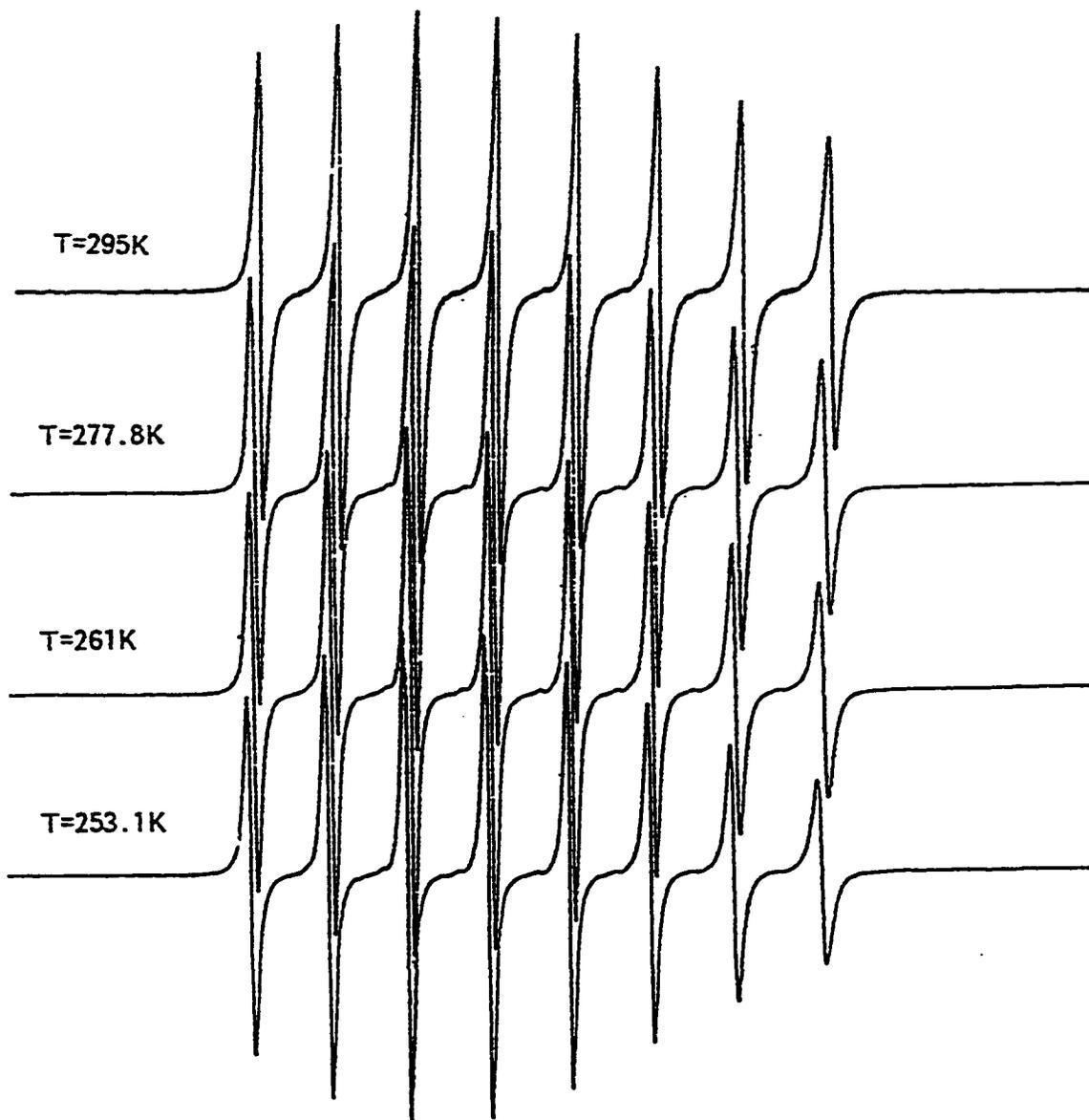


Fig.5.3 : Experimental spectra of Cp_2VCl at Different Temperature

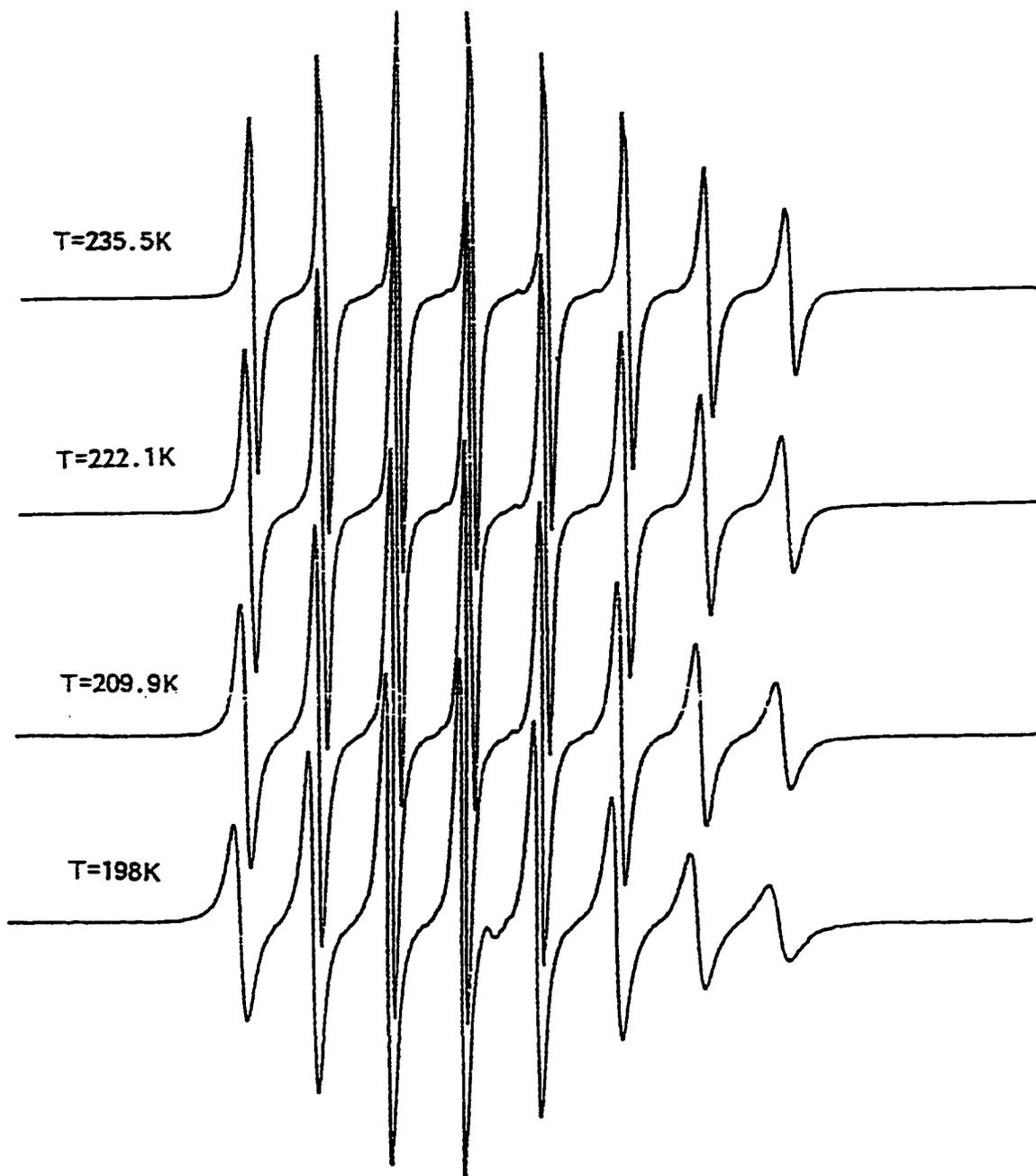


Fig.5.3(cont.) : Experimental Spectra of Cp_2VCl_2 at Different Temperature

Table 5.2 : Experimental Values of Temperature & Lines Widths

Measurement & B & C values

T.°C	21.7	4.6	-12.2	-20.1	-37.7	-51.1	-63.3	-75.2
ΔH(M)								
ΔH-7/2	6.523	6.475	6.679	6.902	7.598	8.771	10.647	13.757
ΔH-5/2	6.147	6.033	6.081	6.208	6.540	7.216	8.396	10.434
ΔH-3/2	5.982	5.846	5.824	5.888	6.026	6.463	7.223	8.736
ΔH-1/2	6.026	5.900	5.880	5.934	6.042	6.418	7.123	8.497
ΔH+1/2	6.258	6.147	6.223	6.328	6.549	7.066	7.992	9.692
ΔH+3/2	6.664	6.697	6.884	7.094	7.583	8.457	9.922	12.285
ΔH+5/2	7.274	7.448	7.813	8.154	9.088	10.538	12.847	16.433
ΔH+7/2	8.050	8.382	9.056	9.589	11.101	13.361	16.709	21.930
B	.2312	.2864	.3541	.4002	.5202	.6681	.9077	1.207
c	.09568	.1168	.1512	.1759	.2544	.3601	.5103	.7302
C/B	.41384	.4078	.4270	.4395	.4890	.5390	.5622	.6050

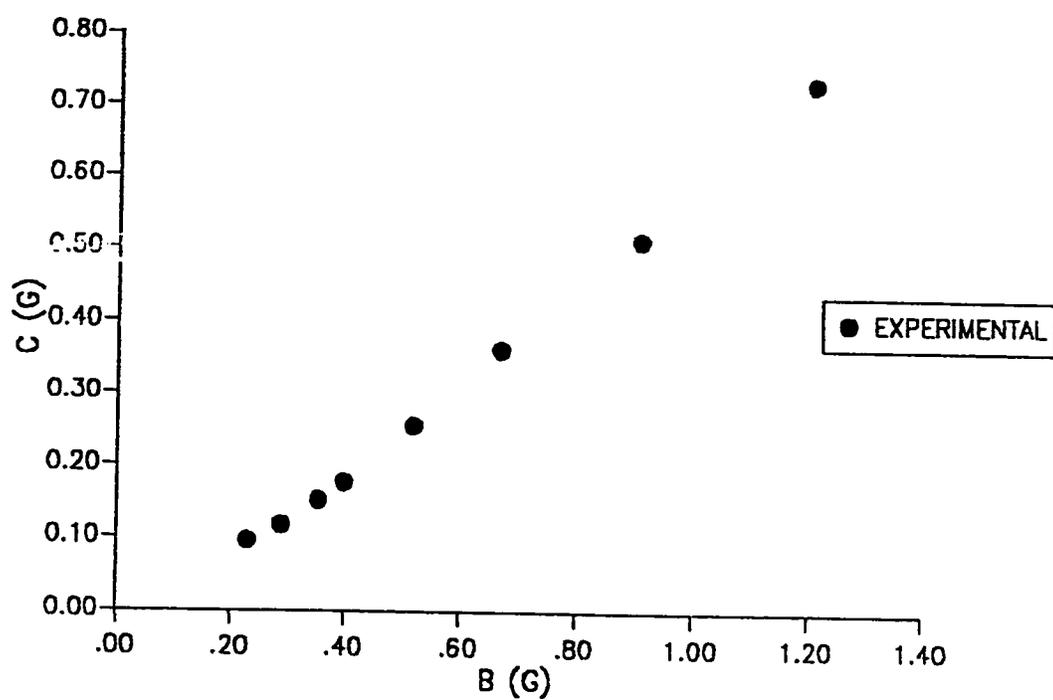


Fig. 5.4 : B Vs C Experimental Values for Cp_2VCl_2 at Different Temperatures.

In order to determine the anisotropy of rotation (N) and the axis of rotation, the theoretical calculations of B and C coefficients which are given by the following equations [59,60] were used. These equations were adopted from the motional-narrowing of Freed et al. [42,43,44].

$$C_o = \frac{8}{3} - [1 + (\omega_e \tau_o)^2]^{-1} - \frac{1}{\{3 [1 + (\omega_o \tau_o)^2]\}} \quad (5.9)$$

$$C_2 = \frac{8}{3} - [1 + (\omega_e \tau_2)^2]^{-1} - \frac{1}{\{3 [1 + (\omega_o \tau_2)^2]\}} \quad (5.10)$$

$$C = \frac{2}{\sqrt{30}} (0.8\pi)^2 (D_o \tau_o C_o + 2D_2^2 \tau_2 C_2) \quad (5.11)$$

$$B_o = \frac{16}{3} + \frac{4}{1 + (\omega_o \tau_o)^2} \quad (5.12)$$

$$B_2 = \frac{16}{3} + \frac{4}{1 + (\omega_2 \tau_2)^2} \quad (5.13)$$

$$B = \frac{-2}{\sqrt{30}} (0.1\pi \omega_o) (g_1 D_o \tau_o B_o + 2g_2 D_2 \tau_2 B_2) \quad (5.14)$$

where

$$\nu = g_o (8.7940147 \times 10^6) \text{ Sec}^{-1} \text{ G}^{-1}$$

$$\nu_p = \nu \frac{g_o}{g_e}$$

$$g_e = 2.00232$$

$$g_1 = (g_x - g_o) \left(\frac{3}{2}\right)^{\frac{1}{2}}$$

$$g_2 = \frac{(g_x - g_y)}{2}$$

$$\omega_o = 2 \pi \nu$$

$$\omega_a = \frac{A_o \nu_p}{2}$$

$$D_o = (A_x - A_y) \left(\frac{|\nu_p|}{2\pi}\right) \left(\frac{3}{8}\right)^{\frac{1}{2}}$$

$$D_2 = (A_x - A_y) \left(\frac{|\nu_p|}{8\pi}\right) \text{ MHz}$$

$$\tau_o = \tau_R N^{\frac{1}{2}}$$

$$\tau_2 = \frac{3 \tau_o}{(1+2N)}$$

and the magnetic parameters g 's and A 's are in table 5.1 .

The theoretical calculations of B and C coefficients are performed by varying τ_p in eqs(5.9, 5.14). The program written for eqs(5.9-5.14) is known as the ABC program (in Basic language). While performing theoretical simulation at the beginning we mainly look at the C/B ratio with the

anisotropic rotational diffusion and reorientation. Basically, the simulation of these curves using eqs(5.9-5.14) involves the fitting of three parameters the C/B ratio, the principal axis of rotation (Z'-axis) and N. Usually the approximate ratio of C/B gives some hint on the Z'-axis . Value of N is then varied to move the upper portion of the curve to approach the experimental points. Appendix 2 shows N, τ_1 values and the calculated B and C values for Z' = X, Y, Z, by using the reported magnetic parameters in table 5.1, at frequency of 9.511 GHz . plotting B Vs C for the different values of N for each Z'-axis with the B and C values of the experiment will show for which N value and on which Z'-axis the experimental values is fitted. In Fig. 5.5 we use the magnetic parameter set (a) in table 5.1, it is clear that the experimental points are below the calculated values even by increasing the N values for higher number. Similarly by using the values of set (b) in Fig. 5.6 the experimental points below the calculated one. In Fig. 5.7, by using set (c) values, the experimental points come above the calculated values. Fig. 5.8 by using set (d), shows that at Z'=X and Z the experimental points come above the calculated line, but at Z'=Y and for N=2.7 the experimental points are on the calculated line.

This suggest that Cp_2VCl_2 has a preferential rotation axis $Z'=Y$ at $N=2.7$ over the other two axes. In Fig. 5.9 we compare all sets (a,b,c,d) at $Z'=Y$ and $N=2.7$ to show clear fitting of the experimental values with the calculated one. Also in Fig. 5.10 we show the effect of changing the frequency on the fitting which is small effect.

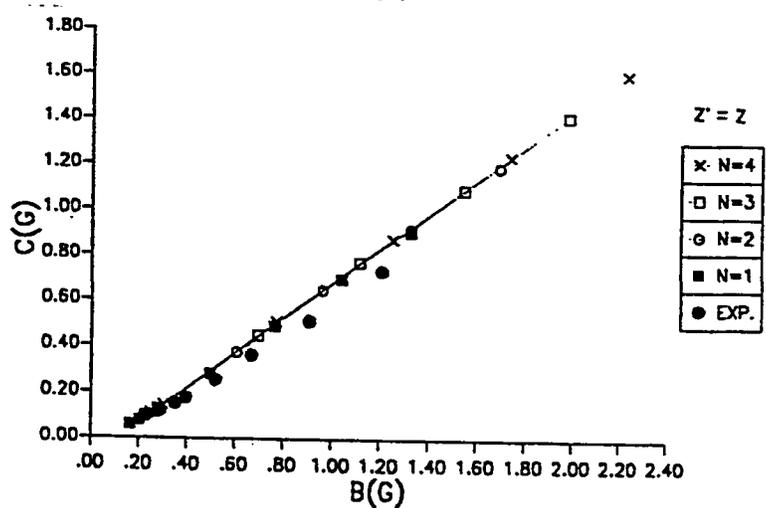
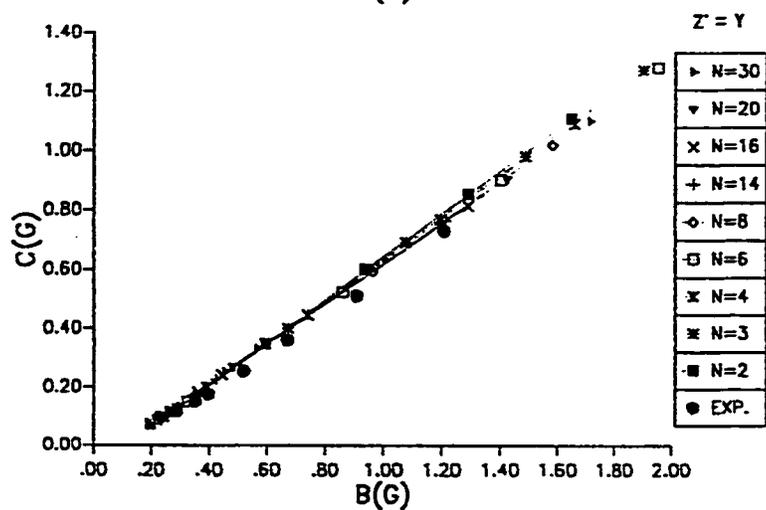
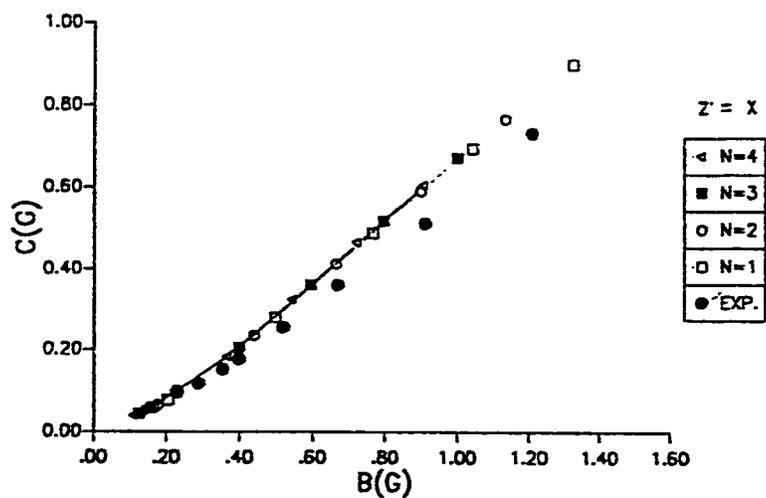


Fig. 5.5 : Experimental fit using the magnetic parameter set(a) for $Z'=X,Y,Z$

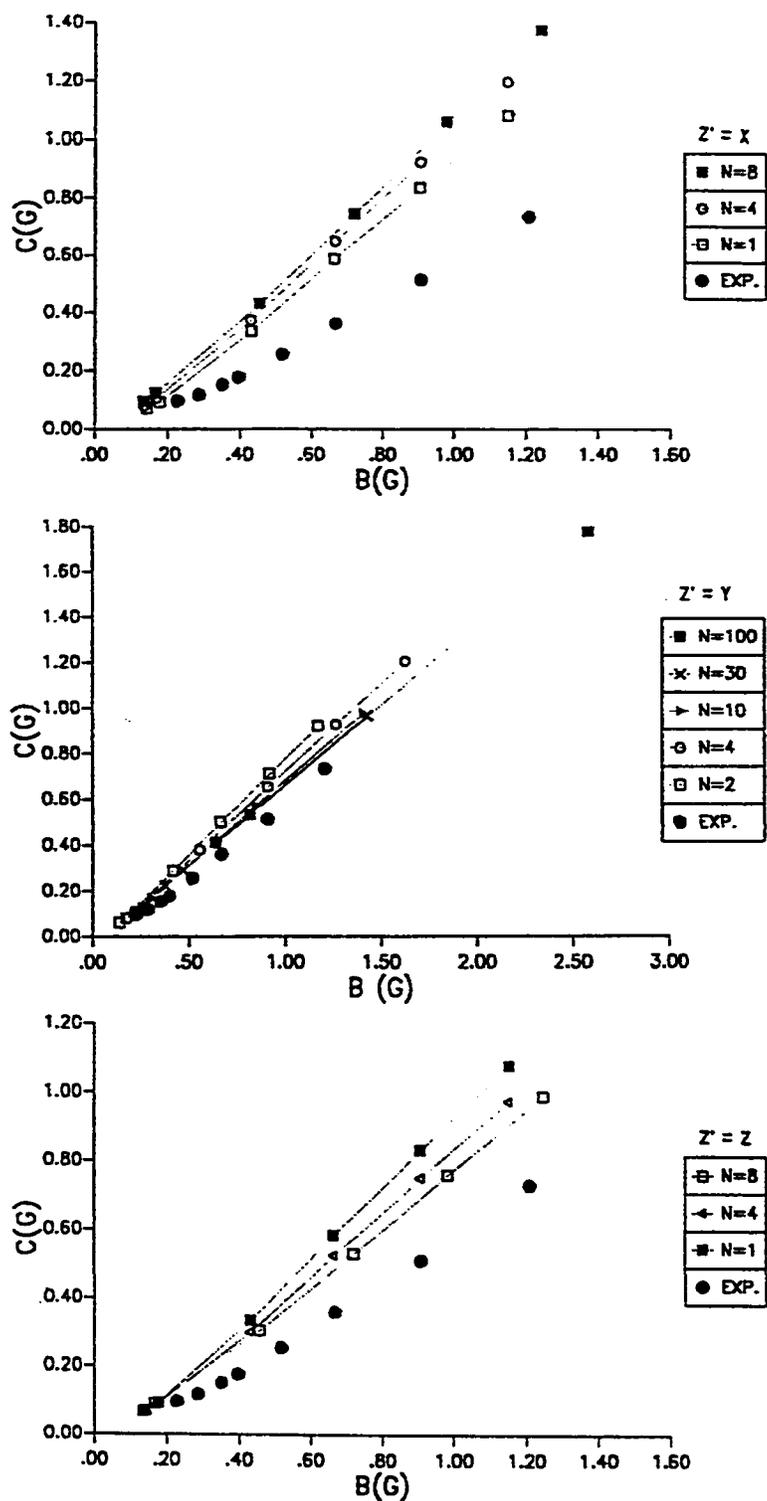


Fig. 5.6 : Experimental fit using the magnetic parameter $set(b)$ for $Z'=X,Y,Z$

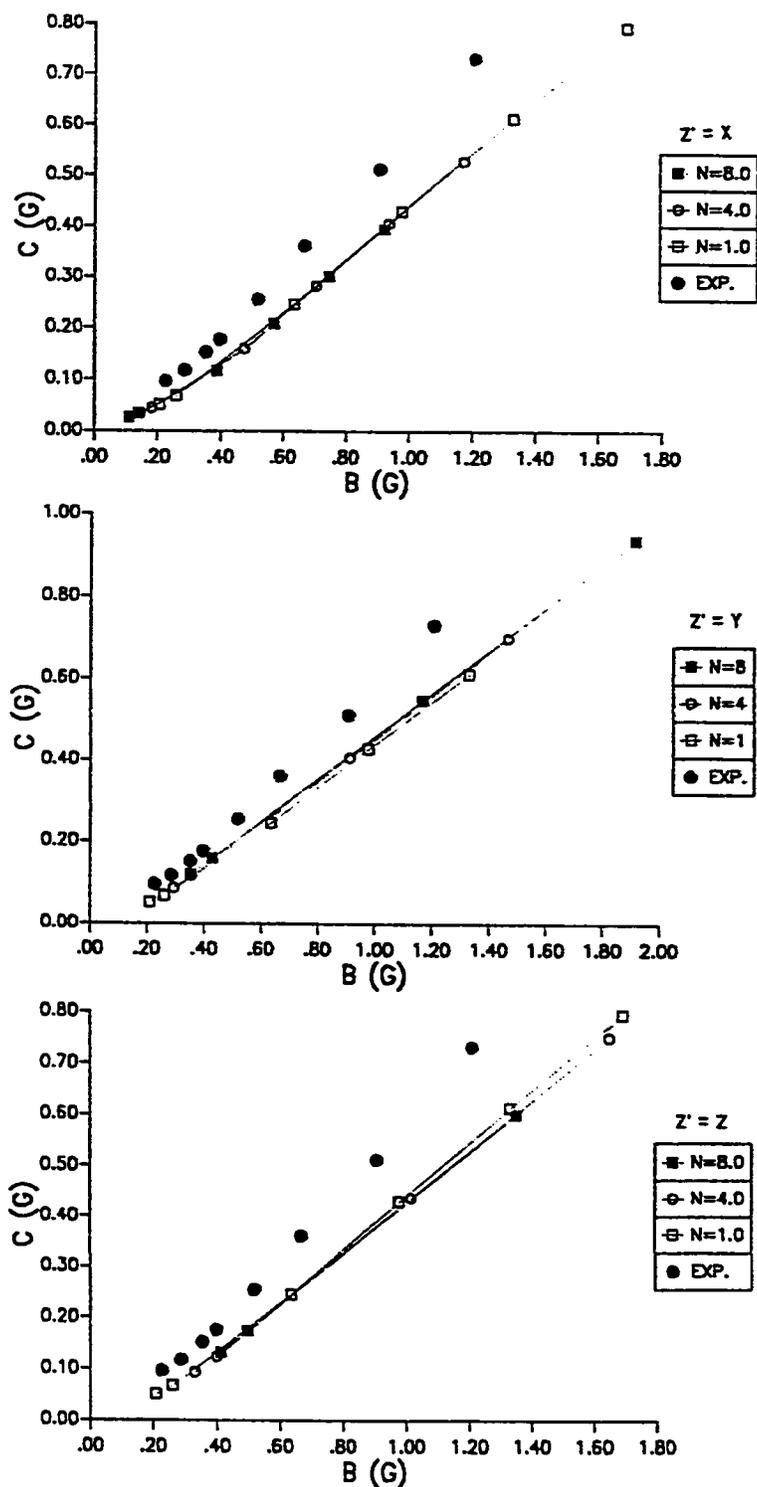


Fig. 5.7 : Experimental fit using the magnetic parameter $\text{set}(c)$ for $Z'=X,Y,Z$

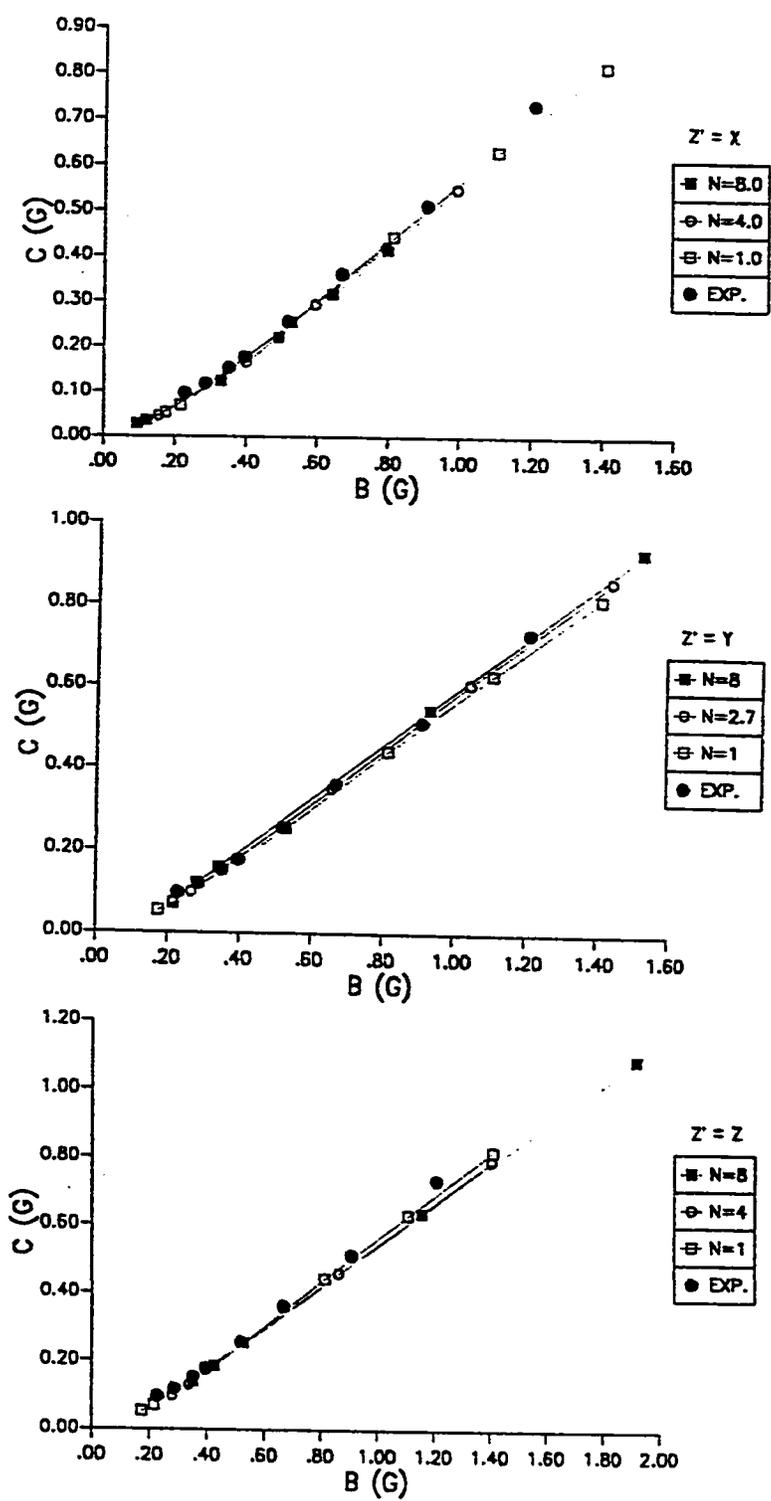


Fig. 5.8 : Experimental fit using the magnetic parameter set(d) for $Z'=X,Y,Z$ the best fit at $Z'=Y$ and $N=2.7$

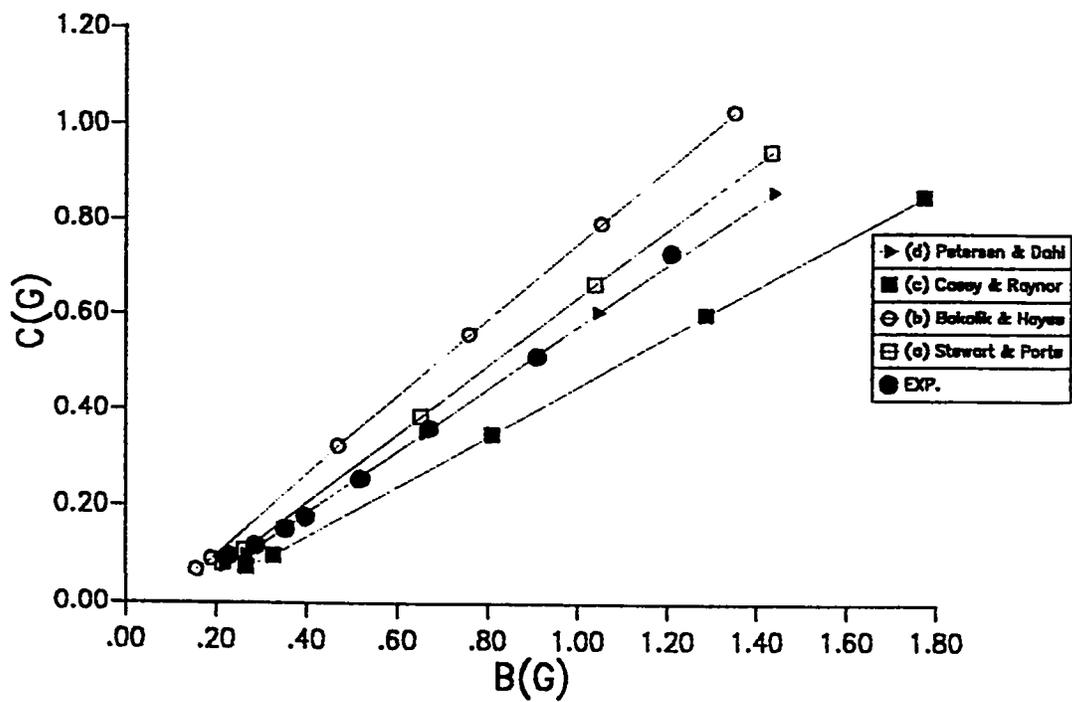


Fig. 5.9 : Experimental fit using the magnetic parameter sets(a,b,c,d) for $Z'=Y$ and $N=2.7$

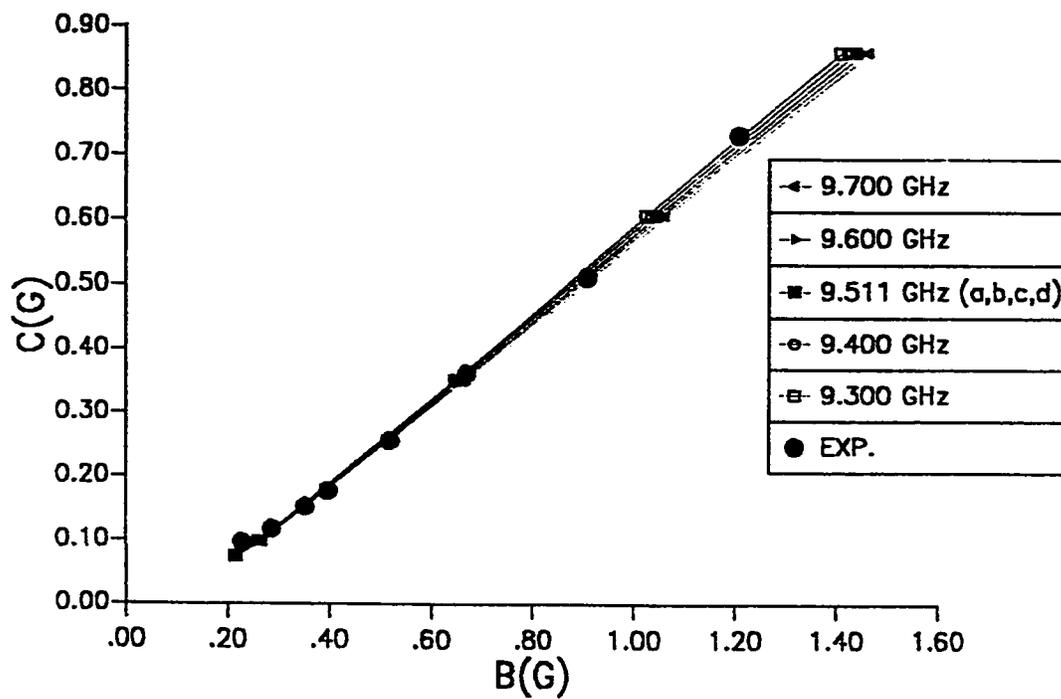


Fig. 5.10 : Experimental fit using the magnetic parameter set(d) for $Z'=Y$ and $N=2.7$ at different frequency.

Form the best fit plot in (Fig.5.8) one can generate an other plot between the C theoretical values Vs. τ_p (table 5.3), then we can extrapolate the experimental values C to fined τ_p experimental values from the straight line equation of C Vs. τ_p of the theoretical values, as in Fig.5.11 . By knowing the experimental τ_p at different temperature and the viscosity at different temperature we can use eq.(1.2) to find the relation (slope), if we plot the correlation time τ_p Vs. $\eta/T(K)$ for this experimental values as in Fig.5.12 . The viscosity values for toluene is corrected for different temperature using [61] the following equation,

$$\eta = \frac{a}{(b + t)^n} \quad (5.15)$$

where, $a = 18.954$, $b = 112.99$, $n = 1.6522$, t the temperature in($^{\circ}C$), the calculated value of η is in poise. Table 5.4 show the results. The slope we get is equal to

$$\text{The Slope} = \left(\frac{4}{3} \pi r_0^3 \right) \left(\frac{\kappa}{k_B} \right)$$

where the first bracket term is the volume of the molecule and the value of r_0 will be determine in the next experimental result from the translational diffusion constant, from which we can recalculate κ (the anisotropic interaction parameter) instead of estimating its volume by computer program or other method. The value of the slope is = 2.6823×10^{-7} sec KP^{-1} , by using the first seven point (corr.=0.9993) instead of (corr.=0.9951) for all eight points.

Table 5.3 : Theoretical values of B & C & τ_p for Z'=Y and

N=2.7

$\tau_p(\text{sec } \times 10^{-11})$	B (G)	C (G)
0.688	0.2168	0.07376
0.888	0.2658	0.09757
2.880	0.6578	0.3502
4.880	1.042	0.6053
6.880	1.435	0.8599
8.880	1.833	1.115

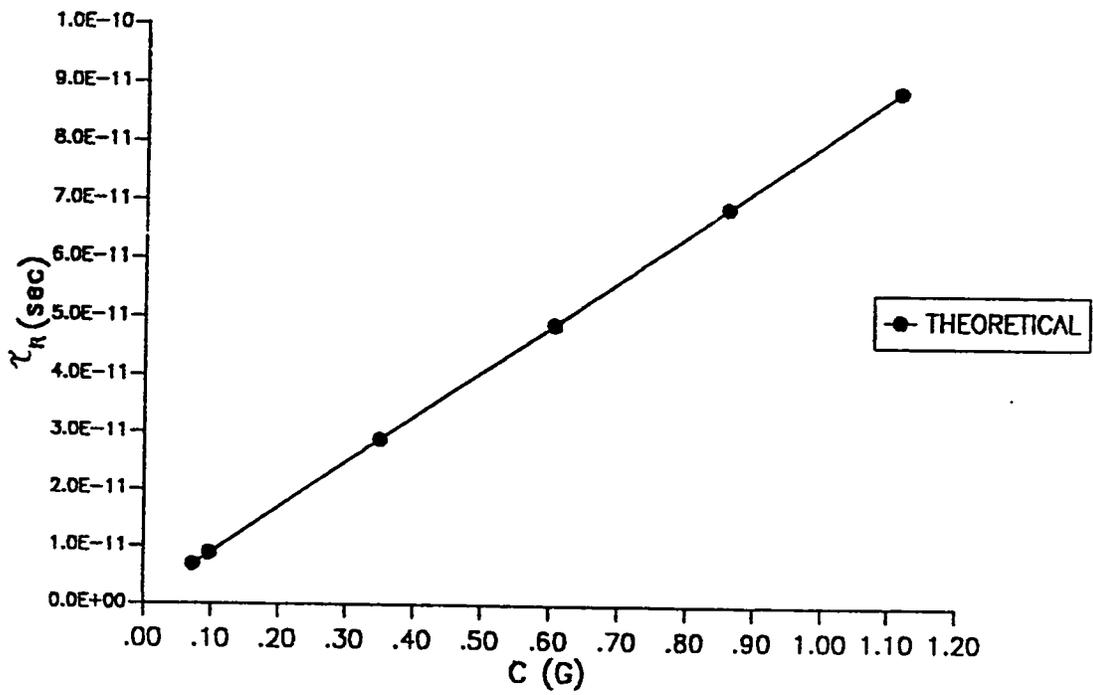


Fig. 5.11 : C Vs. τ_R theoretical values

Table 5.4 : η , η/T and τ_p Values for Cp_2VCl_2 Experiment

t ($^{\circ}C$)	η ($P \times 10^{-2}$)	η/T ($PK^{-1} \times 10^{-4}$)	τ_p (Sec. $\times 10^{-11}$)
21.7	0.57493	0.19500	0.87059
4.6	0.71951	0.25906	1.03664
-12.2	0.92823	0.35573	1.30709
-20.1	1.0622	0.41979	1.50128
-37.7	1.5030	0.63837	2.11845
-51.1	2.0777	0.93574	2.94947
-63.3	2.9862	1.4231	4.13035
-75.2	4.6942	2.3715	5.85921

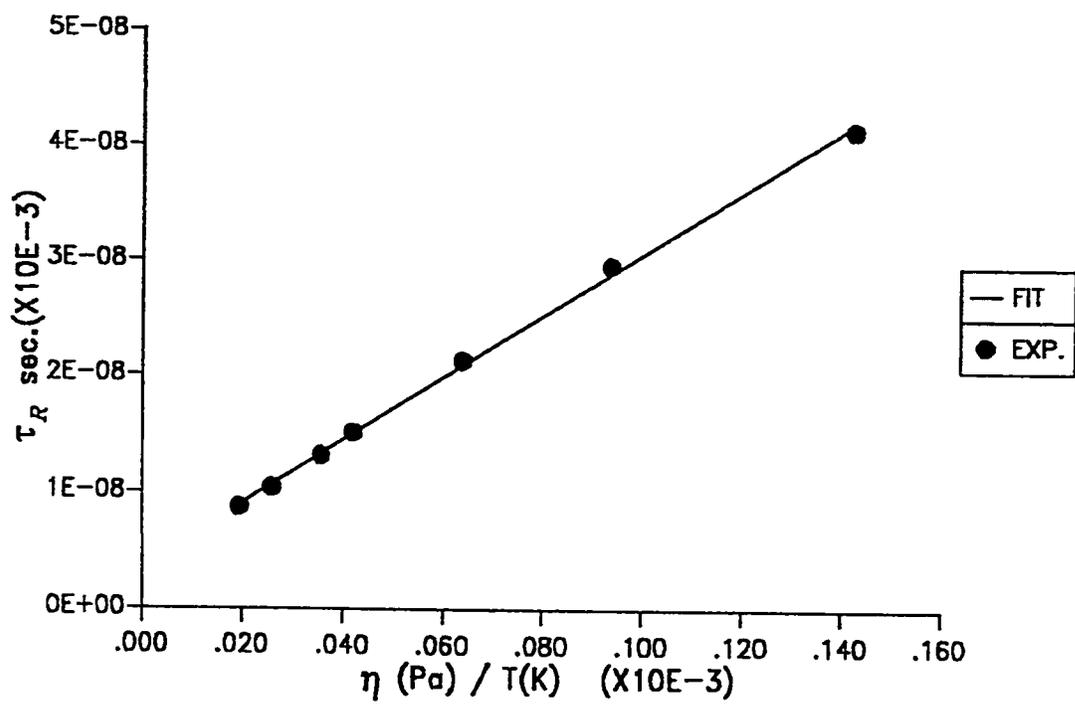


Fig. 5.12 : η/T Vs. τ_R for Cp_2VCl_2 in toluene.

5.3 : DETERMINATION OF THE TRANSLATIONAL DIFFUSION CONSTANT

5.3.1 Cp_2VCl_2

The capillary diffusion experiment which described in chapter three is done for Cp_2VCl_2 as mention. The experiment takes abut 22 days , from which 57 spectra are collected. The time is recorded for each run and each spectra is analysis to find the intensities of the eight lines (if the intensities dropped much at short time this give a result that high concentration difference between the solution in the capillary and the solvent above it, so the experiment will continue until a steady state start, then readjust the ESR spectrum parameter to get maximum intensities again Fig.5.13 show this effect). The data are introduce to the program "DIFF.FORTRAN " and the best least square fit is result form which the time and intensity for this spectra is substituted in eq.(1.13) to get the diffusion constant, and the value is recalculated by the fitting coefficients , also the ratio R ($R = \ell_s / \ell_t$) is recalculated to check the experimental data. The line widths is calculated and adjusted by the following ,

$$\Delta\omega_i = \Delta\omega_o (I_o/I_i)^{\frac{1}{2}} \quad (5.16)$$

$\Delta\epsilon_0$ and I_0 is the line width and intensity of a reference line and in Cp_2VCl_2 it the line of $M = -1/2$, $\Delta\epsilon_i$ and I_i is the line width and intensity for a particular line, the average line widths is calculated and given by the program. The relative total intensity for each spectra and the time is in table 5.5 . In (Fig.5.14) some full spectra are shown. From the diffusion constant the program generate the fitting curve points (time Vs. intensity) which is plotted with the experimental values in (Fig. 5.15). In Appendix 1 the computer program, the input data and the output calculation is given for Cp_2VCl_2 . The diffusion constant calculated by the DIFF. computer program for Cp_2VCl_2 is = $0.520 \times 10^{-5} \text{ cm}^2/\text{s}$
 The value of $R(= \ell_s / \ell_t)$ simulated is with in experimental error $R(\text{exp})=0.57143$, $R(\text{cal.})=0.5700$.

5.3.2 VO[5-MeO-Sal SB](Phen)

For this compound the experiment goes for abut 25 days and 62 spectra were taken , some of the spectra are shown in Fig.5.16. The full data are in appendix.1 as computer output also the calculation of total intensity and time is given . A plot of time vs. intensity is in Fig.5.17 with a simulation of the result, the value of diffusion constant obtained for VO[5-MeO-Sal SB](Phen) is = $0.655 \times 10^{-5} \text{ cm}^2/\text{sec}$.
 The value of $R(= \ell_s / \ell_t)$ simulated is with in experimental error $R(\text{exp})=0.57143$, $R(\text{cal.})=0.6800$.

Table 5.5 : Time and Intensity of Cp_7VCl_2 diffusion

p.#	Time(Sec)	Intensity	p.#	Time(Sec)	Intensity
* 1	0.240×10^3	1.000	30	0.860×10^6	0.451
2	0.480×10^5	0.966	31	0.901×10^6	0.439
3	0.744×10^5	0.961	32	0.927×10^6	0.416
4	0.876×10^5	0.951	33	0.940×10^6	0.407
5	0.138×10^6	0.912	*34	0.986×10^6	0.394
6	0.153×10^6	0.909	35	0.101×10^7	0.384
* 7	0.173×10^6	0.893	36	0.104×10^7	0.370
8	0.212×10^6	0.862	37	0.107×10^7	0.359
9	0.241×10^6	0.826	38	0.107×10^7	0.337
10	0.259×10^6	0.831	39	0.112×10^7	0.335
11	0.297×10^6	0.806	*40	0.117×10^7	0.312
12	0.318×10^6	0.791	41	0.119×10^7	0.298
13	0.346×10^6	0.756	42	0.121×10^7	0.307
*14	0.381×10^6	0.751	43	0.125×10^7	0.291
15	0.409×10^6	0.718	44	0.127×10^7	0.284
16	0.431×10^6	0.706	45	0.130×10^7	0.273
17	0.469×10^6	0.687	46	0.133×10^7	0.266
18	0.498×10^6	0.657	47	0.136×10^7	0.256
19	0.512×10^6	0.651	*48	0.138×10^7	0.244
20	0.555×10^6	0.634	49	0.142×10^7	0.238
*21	0.579×10^6	0.615	50	0.144×10^7	0.222
22	0.604×10^6	0.590	51	0.147×10^7	0.226
23	0.667×10^6	0.556	52	0.151×10^7	0.226
24	0.690×10^6	0.543	53	0.155×10^7	0.208
25	0.727×10^6	0.522	54	0.159×10^7	0.192
26	0.758×10^6	0.510	55	0.169×10^7	0.172
*27	0.778×10^6	0.497	*56	0.180×10^7	0.151
28	0.815×10^6	0.472	57	0.188×10^7	0.145
29	0.843×10^6	0.466			

*The spectra shown in Fig. 5.14 .

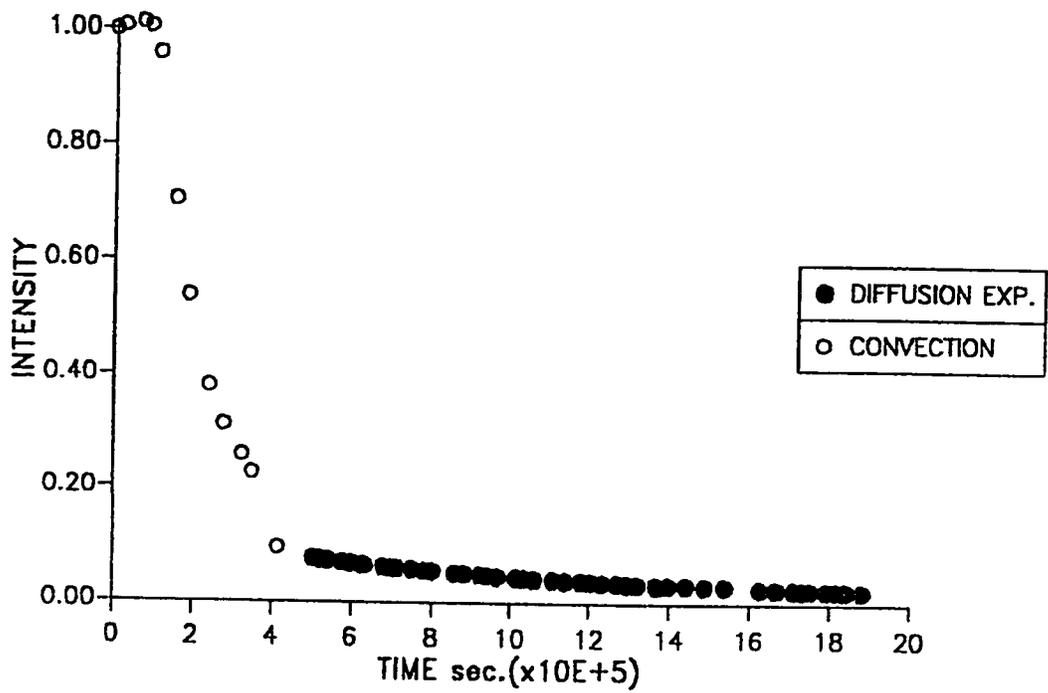


Fig. 5.13 : The effect of high concentration on diffusion.

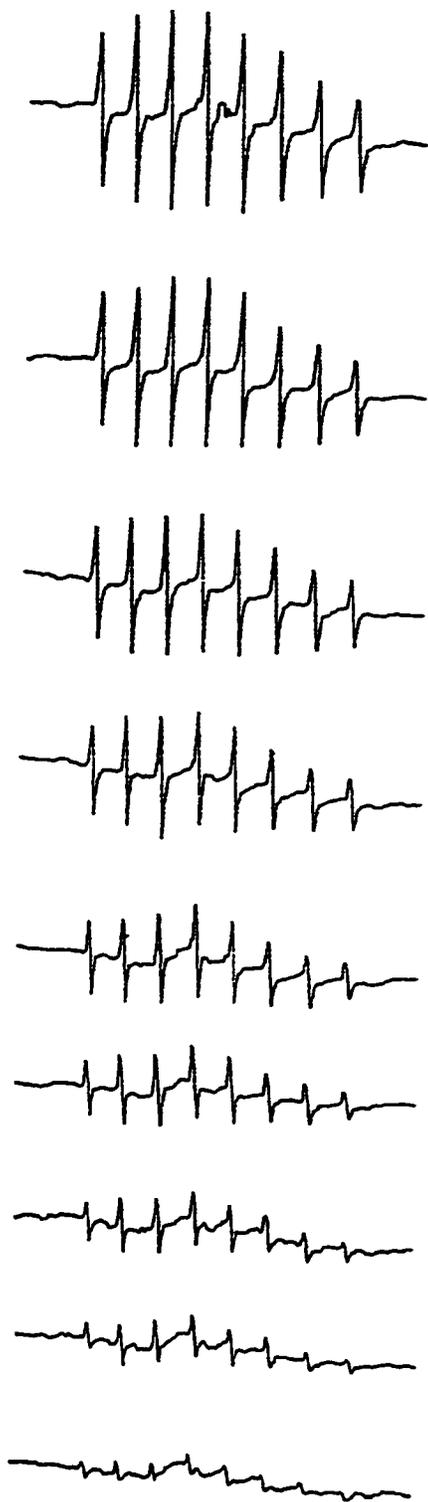


Fig. 5.14 : First derivative ESR spectra of Cp_2VCl_2 at different time, the time and total intensity is marked (*) in table 5.5 .

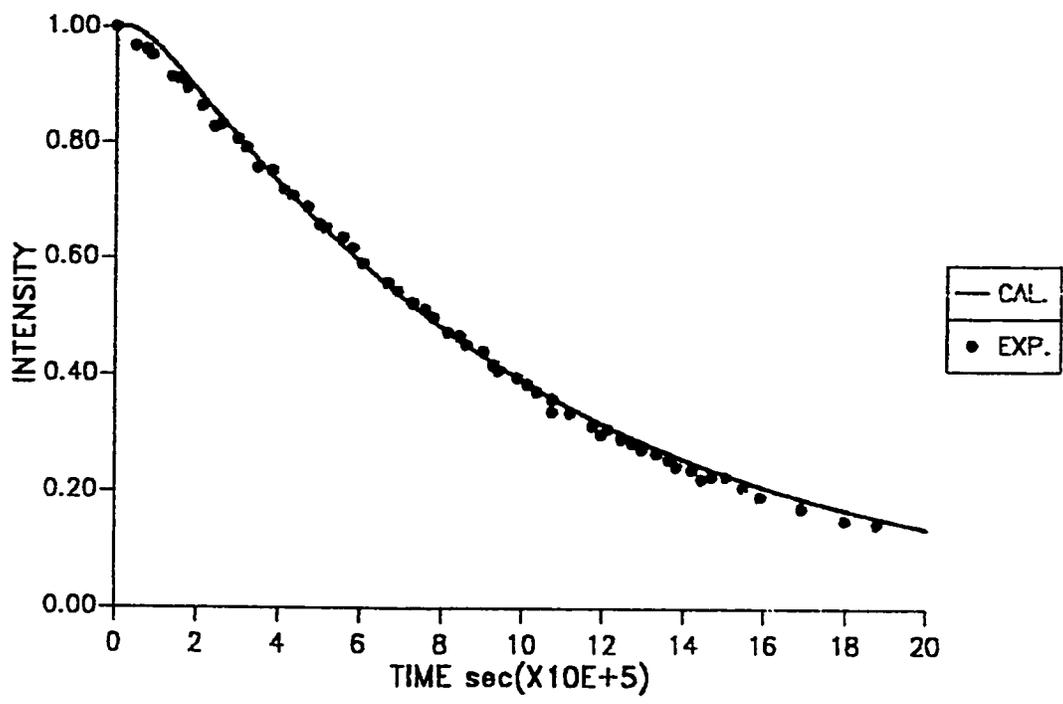


Fig. 5.15 : Total ESR intensities as a function of time for Cp_2VCl_2

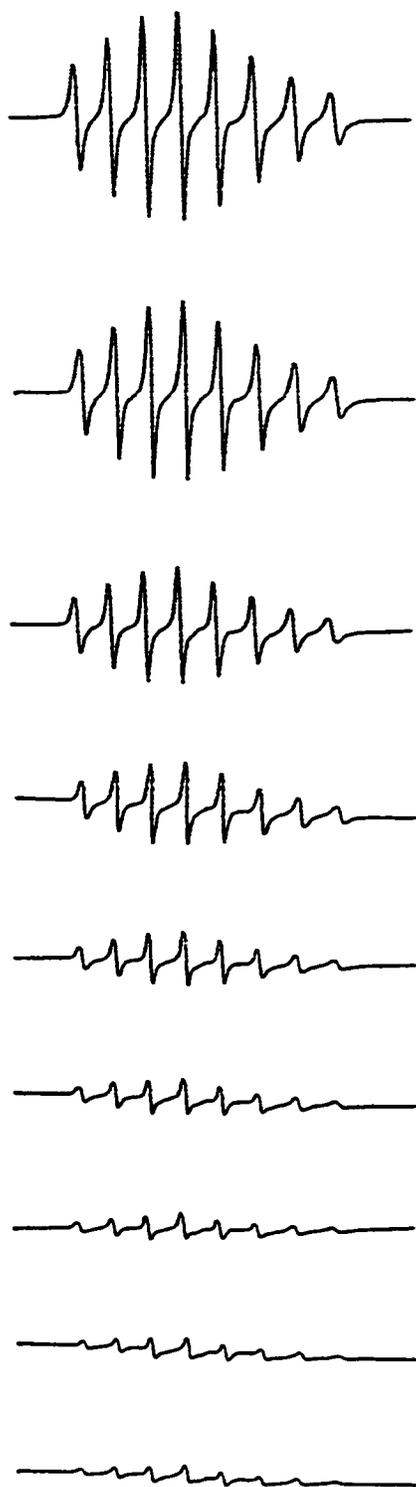


Fig. 5.16 : First derivative ESR spectra of VO[5-MeO-Sal SB](Phen) at different time.

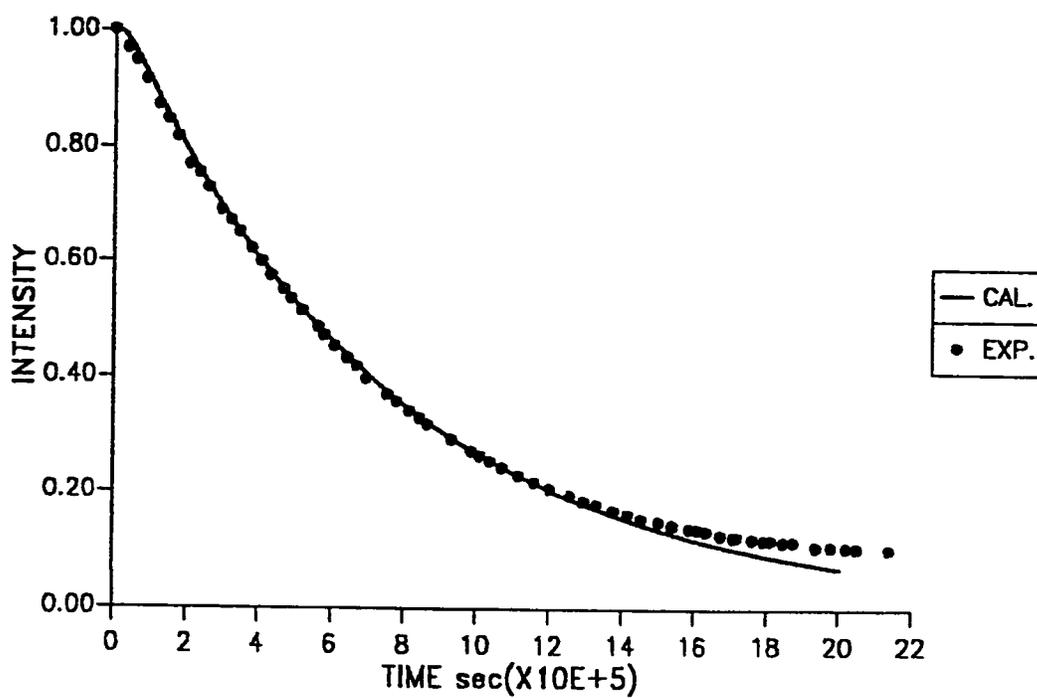


Fig. 5.17 : Total ESR intensities as a function of time for
VO[5-MeO-Sal SB](Phen)

5.3.3 The Hydrodynamic radius

To calculate the hydrodynamic radius of these compound one can use Stokes-Einstein equation(1.1), by substituting the values obtained for the diffusion constant and the other parameter the unit for each are, D in cm^2/sec , k in JK^{-1} , T room temperature in K, η in $\text{Pa}\cdot\text{sec}$, by multiplying with a factor of (10^{14}) the hydrodynamic radius results in (\AA).

This hydrodynamic radius actually is the diameter of the molecules since it is rotating as sphere, so the radius of it will be half of the value obtained from the equation, to get the volume of the molecule by using the formula of the sphere volume $(\frac{4}{3} \pi r_0^3)$. The value results are in the next table,

Compound	Hyd.radius(\AA)	Radius(\AA)	Volume(\AA^3)
Compound I	7.820	3.910	250.0
Compound II	8.073	4.036	275.0

Compound I : Cp_2VCl_2

Compound II : $\text{VO}[5\text{-MeO-Sal SB}](\text{Phen})$

5.3.4 Simulating The Molecular Volume from Disk Top

Molecular Model [DTMM]

A computer program available for drawing the molecular structure and minimizing the energy of the molecules by iteration to minimize the torsional angle, Van der waals forces, and other effect. Then it produce a structure that can be compare with the real molecule, looking at it from different angel and position and rotation on different axis carefully we can measure at least three lengths to determine the volume of this molecule, the structures are shown in Fig.5.18 for the two compound under study and a module compound vanadyl bisacetylacetonate or [VOAA], Fig.5.19 is the crystal structure of VOAA. The lengths obtained for the compounds are

Compound	r_1 (Å)	r_2 (Å)	r_3 (Å)
Compound I	5.78	7.73	4.27
Compound II	8.6	13.5	6.4

Compound I : Cp_2VCl_2 , Compound II : $VO[5-MeO Sal SB](Phen)$

where r_1 , r_2 , r_3 are the length between two atoms in the same plane. Looking at the two tables above we can correlate between the hydrodynamic radius of the first compound (I) and the value of r_2 , similarly for compound (II) the hydrodynamic radius is comparable with r_1 , this give some indication that these compounds are rotating along these axis.

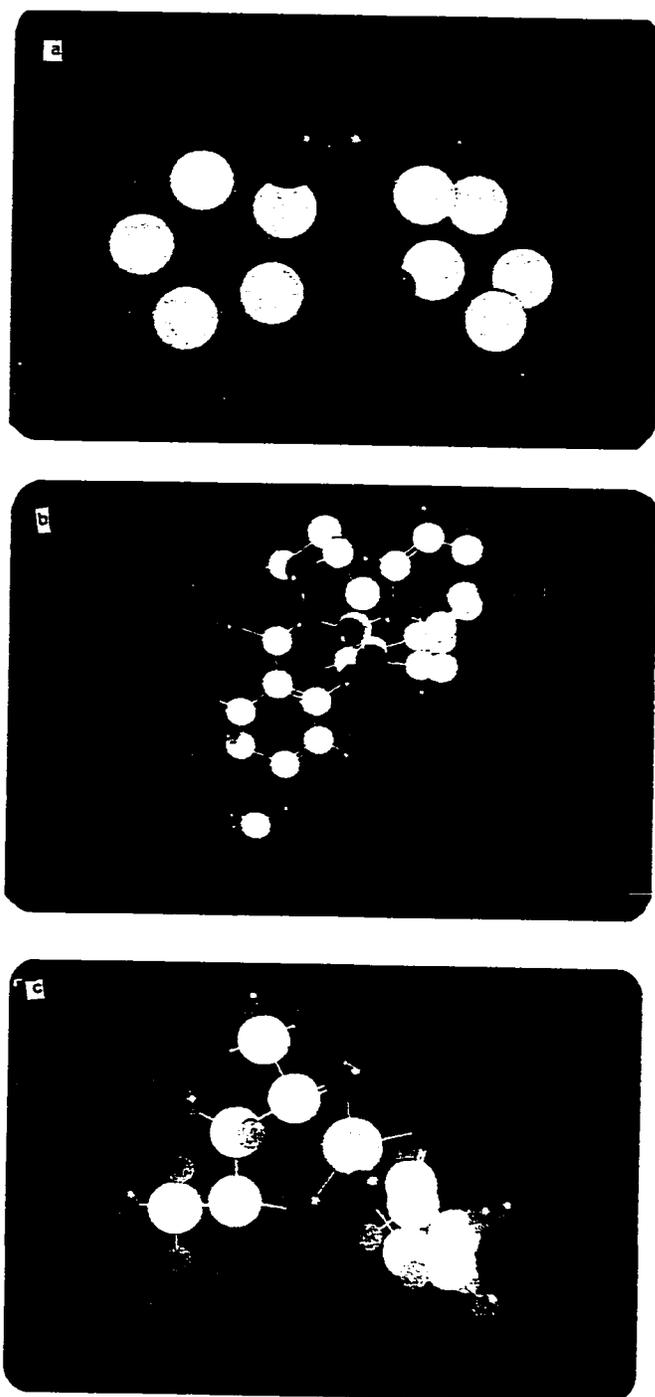


Fig. 5.18 : Structures by DTMM program, (a) Cp_2VCl_2 ,
 (b) $\text{VO}[5\text{-MeO Sal SB}](\text{Phen})$, (c) VOAA

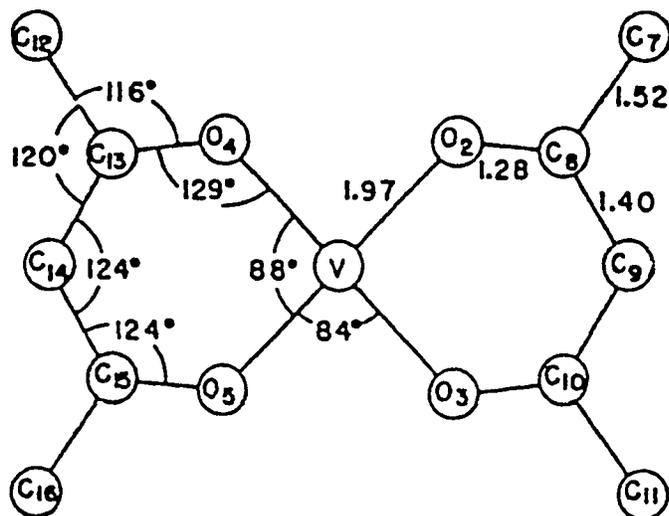
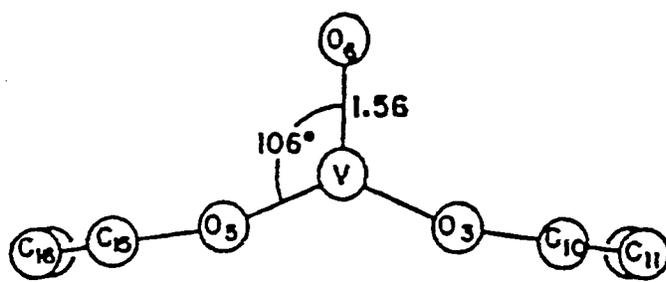
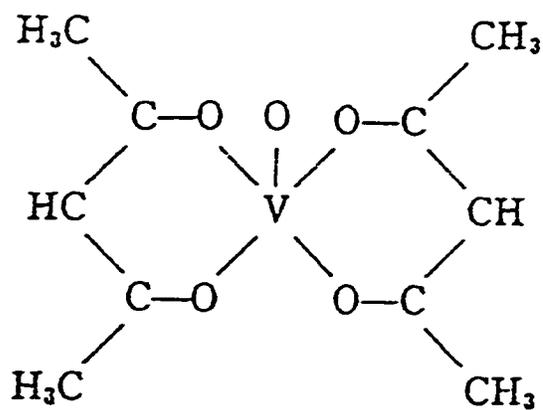


Fig. 5.19 : Crystal Structure of VOAA

5.4 ANISOTROPIC INTERACTION PARAMETER (κ)

In previous discussion in section 5.2 for Cp_2VCl_2 we end with the slope value results from plotting (η/T) Vs. τ_p which are related by eq.(1.2), the slope is equal to

$$\text{The Slope} = \left(\frac{4}{3} \pi r_o^3\right) \left(\frac{\kappa}{k_B}\right)$$

to get κ we substitute by the volume in (\AA^3), the slope in (sec.Pa.K^{-1}), and k_B in (JK^{-1}), then a factor of (10^{31}) is multiplied to get κ as dimensionless unit,

$$\begin{aligned}\kappa &= (\text{slope}) (k_B) (10^{31}) / \text{volume} \\ &= (2.6823 \times 10^{-7}) (1.381 \times 10^{-23}) (10^{31}) / (250.0) \\ \kappa &= 0.150\end{aligned}$$

5.5 STICKINESS FACTOR (S)

We can analyze the data for Cp_2VCl_2 which we get from DTMM program in terms of hydrodynamic free space model for molecular relaxation in liquids [63], a theory based on the existence of free space in the hydrodynamic continuum in which the molecule can rotate. We assign the major axis $r_1 = 7.73$ $r_2 = 5.78$, the symmetry parameter α^1 is defined as (r_2/r_1) . Under stick boundary conditions, where the fluid at the surface of the rotating body sticks to it [68], the dimensionless rotational friction coefficient can be calculated from Perrin's formula [64],

$$f_{stick}^1 = \frac{(2/3) [1 - (\alpha^1)^4]}{\frac{[2 - (\alpha^1)^2] (\alpha^1)^2}{[1 - (\alpha^1)^2]^{\frac{1}{2}}} \ln \left| \frac{1 + [1 - (\alpha^1)^2]^{\frac{1}{2}}}{\alpha^1} \right| - (\alpha^1)^2} \quad (5.17)$$

$\alpha^1 = 0.747$. This gives $f_{stick}^1 = 1.11$. The corresponding value of the dimensionless rotational friction coefficient under slip boundary condition, f_{slip}^1 , is calculated from

$$f_{slip}^+ = f_{stick}^+ [1 - (f_{stick}^+)^{\frac{-2}{3}}] \quad (5.18)$$

the value of $f_{slip}^+ = 0.0761$

from the relation,

$$f_{slip}^+ = f_{stick}^+ C_{slip}^{hyd} \quad (5.19)$$

we can calculate

$$C_{slip}^{hyd} = 0.0761/1.111 = 0.0685$$

The stickiness factor, S is independent of molecular geometry and is zero in the slip limit and 1 in the stick limit. For truly hydrodynamic situation $0 \leq S \leq 1$, from eq.(1.3) we can calculate

$$S = (0.150 - 0.0685)/(1 - 0.0685) = 0.0875 .$$

5.6 STOKES-EINSTEIN MODEL

The anisotropy for rotation (N), defined as $N = R_{\parallel} / R_{\perp}$ is related to the geometric structure factor by means of the Stokes-Einstein equation[11,59,60],

$$R_i = \frac{k_B T}{(8\pi\eta r_i^3 \sigma_i)}, \quad i = \parallel \text{ or } \perp \quad (5.20)$$

where,

$$N = \frac{R_{\parallel}}{R_{\perp}} = \frac{\sigma_{\perp}}{\sigma_{\parallel}} \quad (5.21)$$

$$\sigma_{\parallel} = \frac{2}{3} \lambda^2 (1+\lambda^2) [(1+\lambda^2) \lambda^{-1} \tan^{-1} \lambda - 1]^{-1} \quad (5.22)$$

$$\sigma_{\perp} = \frac{2}{3} \lambda^2 (2+\lambda^2) [1 - (1-\lambda^2) \lambda^{-1} \tan^{-1} \lambda]^{-1} \quad (5.23)$$

and λ is a dimensionless structure factor,

$$\lambda = \frac{(r_f^2 - r_i^2)^{\frac{1}{2}}}{r_i} \quad (5.24)$$

From the DTMM program data,

$$r_f = 7.73$$

$$r_i = 4.28,$$

$$\lambda = 1.504$$

$$\sigma_f = 0.0405$$

$$\sigma_i = 0.133$$

$$N = 3.28$$

which is within experimental error, with our observation that
is $N = 2.7 \pm 0.5$.

5.7 THE ALLOWED - VALUES EQUATION (AVE)

An alternative approach to the determination of N is through the use of the allowed-values equation,

$$A \rho_x = B \rho_y + C \quad (5.25)$$

derived by Kowert [45,60] from line width coefficients B and c when nonsecular terms are unimportant, which shows that only one value of $\rho_x = R_x/R_z$ is consistent with a C/B ratio obtained from experiment for a given $\rho_y = R_y/R_z$. If we let $C = \gamma$ and $B = \beta$, then

$$A = (\Delta a - 3\delta a)[(\gamma/\beta)(\Delta g^- - 3\delta g^-) - (5/16)(\Delta a - 3\delta a)] \quad (a)$$

$$B = (\Delta a + 3\delta a)[(5/16)(\Delta a + 3\delta a) - (\gamma/\beta)(\Delta g^- + 3\delta g^-)] \quad (b)$$

$$C = 4\Delta a[(5/16)\Delta a - (\gamma/\beta)\Delta g^-] \quad (c)$$

with

$$\Delta a = a_z - (\frac{1}{2})(a_x + a_y) \quad (d)$$

$$\delta a = (\frac{1}{2})(a_x - a_y) \quad (e)$$

$$\Delta g^- = [g_z - (\frac{1}{2})(g_x + g_y)](\beta H_0/\hbar) \quad (f)$$

$$\delta g^- = (\frac{1}{2})(g_x - g_y)(\beta H_0/\hbar) \quad (g)$$

and

$$\rho_x = R_x/R_z \quad (h)$$

$$\rho_y = R_y/R_z \quad (i)$$

The values of the various expressions in Eqs. (a-g) using the magnetic parameters g's and a's of set (d) in table 5.1 are listed as follows (in Hz),

$$Aa = 1.426 \times 10^9$$

$$\delta a = 3.850 \times 10^8$$

$$Ag^- = 7.956 \times 10^8$$

$$\delta g^- = 1.609 \times 10^8$$

$$A = 2.834 \times 10^{16}$$

$$B = 8.572 \times 10^{16}$$

$$C = -2.040 \times 10^{17}$$

Using C/B value of 0.605 determined experimentally in table 5.2, the allowed-values equation is

$$2.834 \times 10^{16} \rho_x = 8.572 \times 10^{16} \rho_y - 2.040 \times 10^{17} \quad (j)$$

Dividing the equation by 2.834×10^{16} gives

$$\rho_x = 3.02 \rho_y - 7.20 \quad (k)$$

Equation (k) is not consistent with isotropic reorientation and only consistent with axial symmetry if $\rho_x = 1$, $\rho_y = 2.71$. The allowed-values equation is very sensitive to variation in the γ/B value. The N obtained by using the allowed-values equation is $N = 2.71 \pm 0.5$. This result is consistent with the value determined from motional - narrowing results and the Stokes-Einstein model.

5.8 CONCLUSION

The translational diffusion constant of Cp_2VCl_2 in chloroform has been measured by the capillary diffusion method at room temperature. The translational diffusion processes were monitored by electron spin resonance spectral intensities as a function of time. A computer program based on Fick's second law of diffusion for the capillary diffusion cell was written to get a single value of the diffusion constant. From the best least squares fit of the total intensities as a function of time a single value of the translational diffusion coefficient of $0.520 \times 10^{-5} \text{ cm}^2/\text{s}$ was found for Cp_2VCl_2 in chloroform measured at room temperature ($25 \pm 1 \text{ }^\circ\text{C}$). Using Stokes-Einstein equation for translational diffusion, the hydrodynamic volume of Cp_2VCl_2 in chloroform was found to be 250 \AA^3 . And for $VO[5\text{-MeO-Sal SB}](\text{Phen})$ the diffusion constant is $0.655 \times 10^{-5} \text{ cm}^2/\text{s}$, the volume is 275 \AA^3 .

From an analysis of the temperature dependent ESR line width of Cp_2VCl_2 in toluene in the motional narrowing region it was found that (1) of the four sets of the magnetic parameters for Cp_2VCl_2 available in the literature, only one set of magnetic parameters by Petersen and Dahl could be used to explain the line width data, and (2) Cp_2VCl_2 in toluene was undergoing anisotropic reorientation with $N = 2.7 \pm 0.5$ at an axis $z' = Y$, where N is the ratio of R_{\parallel}/R_{\perp} and R_{\parallel} is the rotational diffusion constant along the molecular Z axis and

R_1 is the rotational diffusion constant perpendicular to the molecular Z axis. This conclusion is also supported by the analysis of anisotropic and nonsecular contributions to the ESR line widths using the allowed-values equation. The AVE analysis for Cp_2VCl_2 in toluene gives $\rho_x = 3.02 \rho_y - 7.20$ which is not consistent with isotropic reorientation ($\rho_x = \rho_y = 1$) and only consistent with $\rho_x = 1, \rho_y = 2.71$, agrees with our result. This value is also supported by the Stokes-Einstein model. The stickiness factor (S) calculated to be (0.0875) which is close to the slip limit (S = 0).

The anisotropic interaction parameter κ which is a measure of the anisotropy of intermolecular interactions, was found to be 0.15 for Cp_2VCl_2 in toluene. κ for VOAA in toluene is 0.53. The small κ value for Cp_2VCl_2 in toluene could be due to the absence of oxygen atom at the apex of a five coordination square pyramidal structure present in VOAA, resulting in a more slippery boundary condition. Therefore, it seems that κ value decreases with decreasing solute corrugation. So we expect that κ for VO[5-MeO-Sal SB](Phen) to be higher value.

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APPENDIX 1

```

C
C THIS PROGRAM IS A SIMULATION FOR CALCULATION OF THE TRANSLATION DIF00010
C DIFFUSION CONSTANTS THE EQUATIONS ARE TAKEN FROM THE FOLLOWING DIF00020
C REFERENCE" JOURNAL OF MAGNETIC RESONANCE 22,289-293(1976) " DIF00030
C BY THE TITEL " DIFFUSION COEFFICIENT OF PARAMAGNETIC SPECIES DIF00040
C IN SOLUTION " REPORTED BY " NYONG-KU AMM" . DIF00050
C THE CAPILLARY DIFFUSION METHOD ARE USED AND MONITORING THE DIF00060
C INTENSITIES BY ESR MACHIN AS A FUNCTION OF TIME . DIF00070
C NOTE :- DIF00080
C DIF00090
C DIF00100
C 1)THE INPUT ARE :(FIRST READ STATEMENT) DIF00110
C READ THE NAME OF THE COMPOUND DIF00120
C 2)THE INPUT ARE :(SECOND READ STATEMENT) DIF00130
C SPIN=SPIN Q.# FOR THE ATOM IN THE SAMPLE DIF00140
C IDAY=THE NUMBER OF PLOTS TAKEN DIF00150
C REFIN= TOTAL INTINSETY OF THE LINES IN THE REFERENCE SPECTROM DIF00160
C NFIT=THE WANTED FIT ORDER DIF00170
C LT= TOTAL LENGTH OF THE CAPILLARY DIF00180
C LS= LENGTH OF THE CAPILLARY IN THE CAVITY DIF00190
C THE INPUT ARE :(THAERED READ STATEMENT) DIF00200
C EACH LINE FOR ONE SPECTROM DATA,THE ORDER OF INPUT DATA ARE DIF00210
C IRO=INTINSETY OF A REFERENCE LINE IN THE SPECTROM DIF00220
C WLO=THE LINE WIEDTH OF THE REFERENCE LINE IN THE SPECTROM DIF00230
C T=TIME(SEC) OF THE PLOT TAKEN WITH RESPECT TO THE START OF EXP. DIF00240
C E1(K)=THE INTINSETY OF ALL LINES IN THE SPECTROM DIF00250
C DIF00260
C 3) THE OUTPUT RESULTS ARE ON SEPARET TWO FIELDS DIF00270
C 4) THE PROGRAM FINED THE LEAST SQUIRE VALUE OF THE ALL PLOT AND DIF00280
C FROM WHICH IT CALCULATE THE EXP. VALUE OF THE DIF00290
C DIFFUSION COEFFICIENT. DIF00300
C 5) THE CALCULATED DIFFUSION COEFFICIENT VALUE CALCULATED BY CHANGING DIF00310
C THE R-VALUE ( R < 1 ) IN SUBROUTINE DIFSR UNTIL YOU GET BEST FIT DIF00320
C 6) ON PLOT DATA FILE THE PROGRAM GENARETE THE CAL.INTINSETY AND DIF00330
C THE TIME WHICH CAN BE USED AS DATA FOR PLOTING BY ICUSYS,THE DIF00340
C TIME IS SCALED BY DIVIDING IT BY 10E3 FOR PLOTING ONLY DIF00350
C
C ----- DIF00360
C WRITEN BY -----> HASAN,B.BALJHOYOR DIF00370
C DIF00380
C ----- DIF00390
C DIF00400
C DIF00410
C DIF00420
C DOUBLE PRECISION T, ISUM,X,Y, TM, SMOIF, MINSOR, VAR, STNDEV, DIF00430
C 1 ISM,DIF,YCAL,BB,R DIF00440
C CHARACTER NAME*30 DIF00450
C REAL LT,LS,IRO,IO DIF00460
C DIMENSION X(100),Y(100),E1(100),W(100),AVGW(100,20),BB(20), DIF00470
C 2 DIFSR(100) DIF00480
C PI=3.1415926 DIF00490
C ACCUR=0.00001 DIF00500
C SCAL=100000 DIF00510
C ----- DIF00520
C FIRST READ----- DIF00530
C READ(5,16) NAME DIF00540
C 16 FORMAT(A30) DIF00550
C ----- DIF00560
C SECOND READ----- DIF00570
C READ(5,*) SPIN, IDAY, REFIN, NFIT, LT, LS DIF00580
C R=LS/LT DIF00590
C NP=IDAY DIF00600
C M=(2*SPIN)+1 DIF00610
C WRITE(6,3) DIF00620
C 3 FORMAT('1') DIF00630
C WRITE(7,123) DIF00640
C 123 FORMAT('1',' ','001',' ','2') DIF00650
C WRITE(7,321) IDAY DIF00660
C 321 FORMAT(14) DIF00670
C MM=1 DIF00680
C DO 66 L=1, IDAY DIF00690
C Z=-SPIN DIF00700
C ISUM=0 DIF00710
C ----- READ STATEMENT THREE ----- DIF00720
C READ(5,*) IRO,WLO,T,(E1(K),K=1,M) DIF00730
C IO=IRO/REFIN DIF00740
C DO 77 K=1,M DIF00750
C E1(K)=E1(K)/REFIN DIF00760
C ISUM=ISUM+E1(K). DIF00770

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77      CONTINUE
          IF(IISUM.GT.1) GO TO 108
          JJ=08E
          DO 118 N=1,N
              W(N)=WLO*((10/EI(N))**.5)
              AVGM(JJ,N)=W(N)
118      CONTINUE
          IF(IISUM.GT.1) GO TO 108
          NN=JJ+1
108     CONTINUE
          X(L)=T
          Y(L)=ISUM
          IPL=L-1
          TSCAL=T/SCAL
          WRITE(7,100) TSCAL,ISUM
100     FORMAT(2(2X,F16.5))
66      CONTINUE
          N=-SPIN
          WRITE(6,15) NAME
15      FORMAT(/,'THE NAME OF THE SAMPLE IS ',A30)
          WRITE(6,5) N
5        FORMAT(/,'THE NUMBER OF LINES = ',I4)
          WRITE(6,102)
102     FORMAT(/,'          ','THE SPIN Q.# ',' ')
          'THE AVERAGE LINE WIDTH ')
          WRITE(6,700)
700     FORMAT(' ')
          DO 88 K=1,N
              AVRM=0
              DO 99 N=1,JJ
                  AVRM=AVRM+AVGM(N,K)
99      CONTINUE
                  AVRM=AVRM/JJ
          WRITE(6,92) N,AVRM
92      FORMAT(12X,F6.1,7X,F16.5)
          N=N+1
88      CONTINUE
C
C TO CAL. THE VAR = STN.DIV & LOWEST
C
          SHDIF=0
          -----CALLING SIMULATION TO FIND THE BEST FIT AND LEAST SQUARE
          -----VALUE
          CALL SIM(NP,NFIT,X,Y,88)
          WRITE(6,607)
607     FORMAT(/,'PL.#',4X,'TIME.SEC',7X,'EXP.INT1',7X
1        , 'CAL.INT1',7X,'SOUR.DIFF')
          WRITE(6,608)
608     FORMAT(' ')
          DO 60 K=1,NP
              ISM=0
              DO 70 N=1,NFIT
                  YCAL=88(N)*X(K)**(N-1)
                  ISM=ISM+YCAL
70      CONTINUE
              DIF=(Y(K)-ISM)
              DIFSOR(K)=(Y(K)-ISM)**2
              SHDIF=SHDIF+DIFSOR(K)
              WRITE(6,80) K,X(K),Y(K),ISM,DIFSOR(K)
80      FORMAT(13,4D15.3)
60      CONTINUE
              MINSQR=DIFSOR(1)
              J=1
              DO 111 I=2,NP
                  IF(DIFSOR(I).GT.MINSQR) GO TO 12
                  MINSQR=DIFSOR(I)
                  J=I
12      CONTINUE
111     CONTINUE
          WRITE(6,81) J
81      FORMAT(/,'THE LEAST SQUARES FIT IS FOR THE PLT.# ',I3)
          VAR=(SHDIF/(N-1))
          STNDEV=VAR**.5)
          WRITE(6,90) VAR
90      FORMAT(/,'T H E   V A R I A N C E = ',D15.5)
          WRITE(6,91) STNDEV
91      FORMAT(/,'THE STANDARD DEVIATION = ',D15.5)
          ISUM=Y(J)
          T=X(J)

```

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DIF00760
DIF00770
DIF00780
DIF00790
DIF00800
DIF00810
DIF00820
DIF00830
DIF00840
DIF00850
DIF00860
DIF00870
DIF00880
DIF00890
DIF00900
DIF00910
DIF00920
DIF00930
DIF00940
DIF00950
DIF00960
DIF00970
DIF00980
DIF00990
DIF01000
DIF01010
DIF01020
DIF01030
DIF01040
DIF01050
DIF01060
DIF01070
DIF01080
DIF01090
DIF01100
DIF01110
DIF01120
DIF01130
DIF01140
DIF01150
DIF01160
DIF01170
DIF01180
DIF01190
DIF01200
DIF01210
DIF01220
DIF01230
DIF01240
DIF01250
DIF01260
DIF01270
DIF01280
DIF01290
DIF01300
DIF01310
DIF01320
DIF01330
DIF01340
DIF01350
DIF01360
DIF01370
DIF01380
DIF01390
DIF01400
DIF01410
DIF01420
DIF01430
DIF01440
DIF01450
DIF01460
DIF01470
DIF01480
DIF01490
DIF01500
DIF01510
DIF01520
DIF01530
DIF01540

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```

C-----CALLING TRSDIFF TO CAL THE EXP.DIFFUSION COEFF. -----
      CALL TRSDIFF (PI,T,LT,R,ACCUR,ISUM,SCAL,DIFF,RINT)
      WRITE(6,23)
23  FORMAT(//,'          TIME . SEC   INTENSITY   R-VALUE
1    DIFF. COEF   ')
      WRITE(6,24) T,ISUM,R,DIFF
24  FORMAT(//,' EXP. VAL   ',E10.4,F13.6,D20.6,D18.6)
C-----CALLING DIFSIM TO GENERATE THE CALCULATED INTINSEY AND TIME
      CALL DIFSIM(R,LT,PI,ACCUR,SCAL,DIFF,BB)
      CALL TRSDIFF (PI,T,LT,R,ACCUR,ISUM,SCAL,DIFF,RINT)
      WRITE(6,996) T,RINT,R,DIFF
996  FORMAT(//,' CAL. VAL   ',E10.4,F13.6,D20.6,D18.6)
      WRITE(6,2)
2    FORMAT('1')
      STOP
      END
C .....
SUBROUTINE TRSDIFF (PI,T,LT,R,ACCUR,ISUM,SCAL,DIFF,RINT)
      DOUBLE PRECISION T,ISUM,R
      REAL LT
      DIFF=0.0
      K=10000000
      A=8/(PI*PI*R)
      DO 11 I=1,K
          N=0
          SUM=0
          DO 22 J=1,100
              EEE=((((2*N)+1)**2)*(PI**2)*DIFF*T)/(4*(LT**2))
              B=((EXP(-EEE))/(((2*N)+1)**2))
              IF (B.LE.ACCUR) GO TO 33
              FFF=(((2*N)+1)*PI*R)/2
              C=SIN(FFF)
              D=((-1)**N)*B*C
              SUM=SUM+D
              N=N+1
          22  CONTINUE
      33  CONTINUE
          RINT=A*SUM
          IF (RINT.LE.ISUM) GO TO 44
          DIFF=DIFF+0.1E-9
      11  CONTINUE
      44  CONTINUE
      RETURN
      END
C .....
SUBROUTINE DIFSIM(R,LT,PI,ACCUR,SCAL,DIFF,BB)
      DOUBLE PRECISION BB,R
      DIMENSION BB(20)
      REAL LT
      WRITE(6,613)
C 613  FORMAT(//,'THIS IS SIMULATION DATA TIME(E+5) VS INTENSITIES ')
      WRITE(6,999)
999  FORMAT(' ')
      WRITE(7,300)
300  FORMAT(' ','R01')
C
C CHANGE R-VALUE TO GET BEST FIT
C
      R=0.5700+00
C
C
      T=0
      A=8/(PI*PI*R)
      DO 11 I=1,R01
          N=0
          SUM=0
          DO 22 J=1,50
              EEE=((((2*N)+1)**2)*(PI**2)*DIFF*T)/(4*(LT**2))
              B=((EXP(-EEE))/(((2*N)+1)**2))
              IF (B.LE.ACCUR) GO TO 33
              FFF=(((2*N)+1)*PI*R)/2
              C=SIN(FFF)
              D=((-1)**N)*B*C
              SUM=SUM+D
              N=N+1
          22  CONTINUE
      33  CONTINUE
          RINT=A*SUM
          TSCAL=T/SCAL

```

```

DIF01550
DIF01560
DIF01570
DIF01580
DIF01590
DIF01600
DIF01610
DIF01620
DIF01630
DIF01640
DIF01650
DIF01660
DIF01670
DIF01680
DIF01690
DIF01700
DIF01710
DIF01720
DIF01730
DIF01740
DIF01750
DIF01760
DIF01770
DIF01780
DIF01790
DIF01800
DIF01810
DIF01820
DIF01830
DIF01840
DIF01850
DIF01860
DIF01870
DIF01880
DIF01890
DIF01900
DIF01910
DIF01920
DIF01930
DIF01940
DIF01950
DIF01960
DIF01970
DIF01980
DIF01990
DIF02000
DIF02010
DIF02020
DIF02030
DIF02040
DIF02050
DIF02060
DIF02070
DIF02080
DIF02090
DIF02100
DIF02110
DIF02120
DIF02130
DIF02140
DIF02150
DIF02160
DIF02170
DIF02180
DIF02190
DIF02200
DIF02210
DIF02220
DIF02230
DIF02240
DIF02250
DIF02260
DIF02270
DIF02280
DIF02290
DIF02300
DIF02310
DIF02320
DIF02330
DIF02340
DIF02350

```


FUNCTION POL(I,X)	DIF03140
DOUBLE PRECISION X,POL	DIF03150
POL=X**0(I-1)	DIF03160
RETURN	DIF03170
END	DIF03180
C	DIF03190
C	DIF03200
C	DIF03210
C	DIF03220
.....	DIF03220
C	DIF03230
C	DIF03240
C	DIF03250
C	DIF03260
C	DIF03270
C	DIF03280
C	DIF03290
C	DIF03300
C	DIF03310
C	DIF03320
C	DIF03330
C	DIF03340
C	DIF03350
C	DIF03360
C	DIF03370
C	DIF03380
C	DIF03390
C	DIF03400
C	DIF03410
C	DIF03420
C	DIF03430
C	DIF03440
C	DIF03450
C	DIF03460
C	DIF03470
C	DIF03480
C	DIF03490
C	DIF03500
C	DIF03510
C	DIF03520
C	DIF03530
C	DIF03540
C	DIF03550
C	DIF03560
C	DIF03570
C	DIF03580
C	DIF03590
C	DIF03600
C	DIF03610
C	DIF03620
C	DIF03630
C	DIF03640
C	DIF03650
C	DIF03660
C	DIF03670
C	DIF03680
C	DIF03690
C	DIF03700
C	DIF03710
C	DIF03720
C	DIF03730
C	DIF03740
C	DIF03750
C	DIF03760
C	DIF03770
C	DIF03780
C	DIF03790
C	DIF03800
C	DIF03810
C	DIF03820
C	DIF03830
C	DIF03840
C	DIF03850
.....	DIF03850

SUBROUTINE DGELS

PURPOSE

TO SOLVE A SYSTEM OF SIMULTANEOUS LINEAR EQUATIONS WITH SYMMETRIC COEFFICIENT MATRIX UPPER TRIANGULAR PART OF WHICH IS ASSUMED TO BE STORED COLUMNWISE.

USAGE

CALL DGELS(R,A,M,N,EPS,IER,AUX)

DESCRIPTION OF PARAMETERS

R - DOUBLE PRECISION M BY N RIGHT HAND SIDE MATRIX (DESTROYED). ON RETURN R CONTAINS THE SOLUTION OF THE EQUATIONS.
A - UPPER TRIANGULAR PART OF THE SYMMETRIC DOUBLE PRECISION M BY M COEFFICIENT MATRIX. (DESTROYED)
M - THE NUMBER OF EQUATIONS IN THE SYSTEM.
N - THE NUMBER OF RIGHT HAND SIDE VECTORS.
EPS - SINGLE PRECISION INPUT CONSTANT WHICH IS USED AS RELATIVE TOLERANCE FOR TEST ON LOSS OF SIGNIFICANCE.
IER - RESULTING ERROR PARAMETER CODED AS FOLLOWS
IER=0 - NO ERROR,
IER=-1 - NO RESULT BECAUSE OF M LESS THAN 1 OR PIVOT ELEMENT AT ANY ELIMINATION STEP EQUAL TO 0,
IER=k - WARNING DUE TO POSSIBLE LOSS OF SIGNIFICANCE INDICATED AT ELIMINATION STEP K+1, WHERE PIVOT ELEMENT WAS LESS THAN OR EQUAL TO THE INTERNAL TOLERANCE EPS TIMES ABSOLUTELY GREATEST MAIN DIAGONAL ELEMENT OF MATRIX A.
AUX - DOUBLE PRECISION AUXILIARY STORAGE ARRAY WITH DIMENSION M-1.

REMARKS

UPPER TRIANGULAR PART OF MATRIX A IS ASSUMED TO BE STORED COLUMNWISE IN $M*(M+1)/2$ SUCCESSIVE STORAGE LOCATIONS, RIGHT HAND SIDE MATRIX R COLUMNWISE IN $M*N$ SUCCESSIVE STORAGE LOCATIONS. ON RETURN SOLUTION MATRIX R IS STORED COLUMNWISE TOO.
THE PROCEDURE GIVES RESULTS IF THE NUMBER OF EQUATIONS M IS GREATER THAN 0 AND PIVOT ELEMENTS AT ALL ELIMINATION STEPS ARE DIFFERENT FROM 0. HOWEVER WARNING IER=k - IF GIVEN - INDICATES POSSIBLE LOSS OF SIGNIFICANCE. IN CASE OF A WELL SCALED MATRIX A AND APPROPRIATE TOLERANCE EPS, IER=k MAY BE INTERPRETED THAT MATRIX A HAS THE RANK K. NO WARNING IS GIVEN IN CASE M=1.
ERROR PARAMETER IER=-1 DOES NOT NECESSARILY MEAN THAT MATRIX A IS SINGULAR, AS ONLY MAIN DIAGONAL ELEMENTS ARE USED AS PIVOT ELEMENTS. POSSIBLY SUBROUTINE DGELS (WHICH WORKS WITH TOTAL PIVOTING) WOULD BE ABLE TO FIND A SOLUTION.

SUBROUTINES AND FUNCTION SUBPROGRAMS REQUIRED

NONE

METHOD

SOLUTION IS DONE BY MEANS OF GAUSS-ELIMINATION WITH PIVOTING IN MAIN DIAGONAL, IN ORDER TO PRESERVE SYMMETRY IN REMAINING COEFFICIENT MATRICES.

C		DIF03860
C	SUBROUTINE DGELS(R,A,M,N,EPS,IER,AUX)	DIF03870
C		DIF03880
C	DIMENSION A(1),R(1),AUX(1)	DIF03890
	DOUBLE PRECISION R,A,AUX,PIV,TB,TOL,PIV1	DIF03900
	IF(M)24,24,1	DIF03910
		DIF03920
C	SEARCH FOR GREATEST MAIN DIAGONAL ELEMENT	DIF03930
1	IER=0	DIF03940
	PIV=0.00	DIF03950
	L=0	DIF03960
	DO 3 K=1,M	DIF03970
	L=L+K	DIF03980
	TB=DABS(A(L))	DIF03990
	IF(TB-PIV)3,3,2	DIF04000
2	PIV=TB	DIF04010
	I=L	DIF04020
	J=K	DIF04030
3	CONTINUE	DIF04040
	TOL=EPS*PIV	DIF04050
C	MAIN DIAGONAL ELEMENT A(I)=A(J,J) IS FIRST PIVOT ELEMENT.	DIF04060
C	PIV CONTAINS THE ABSOLUTE VALUE OF A(I).	DIF04070
C		DIF04080
C		DIF04090
C	START ELIMINATION LOOP	DIF04100
	LST=0	DIF04110
	NN=M*N	DIF04120
	LEND=M-1	DIF04130
	DO 10 K=1,M	DIF04140
C		DIF04150
C	TEST ON USEFULNESS OF SYMMETRIC ALGORITHM	DIF04160
	IF(PIV)24,24,4	DIF04170
4	IF(IER)7,5,7	DIF04180
5	IF(PIV-TOL)6,6,7	DIF04190
6	IER=K-1	DIF04200
7	LT=J-K	DIF04210
	LST=LST+K	DIF04220
C		DIF04230
C	PIVOT ROW REDUCTION AND ROW INTERCHANGE IN RIGHT HAND SIDE R	DIF04240
	PIV1=1.00/A(I)	DIF04250
	DO 8 L=K,NN,M	DIF04260
	LL=L+LT	DIF04270
	TB=PIV1*R(LL)	DIF04280
	R(LL)=R(L)	DIF04290
8	R(L)=TB	DIF04300
C		DIF04310
C	IS ELIMINATION TERMINATED	DIF04320
	IF(K=M)9,19,19	DIF04330
C		DIF04340
C	ROW AND COLUMN INTERCHANGE AND PIVOT ROW REDUCTION IN MATRIX A.	DIF04350
C	ELEMENTS OF PIVOT COLUMN ARE SAVED IN AUXILIARY VECTOR AUX.	DIF04360
9	LR=LST+(LT*(K+J-1))/2	DIF04370
	LL=LR	DIF04380
	L=LST	DIF04390
	DO 14 I=K,LEND	DIF04400
	L=L+11	DIF04410
	LL=LL+1	DIF04420
	IF(L-LR)12,10,11	DIF04430
10	A(LL)=A(LST)	DIF04440
	TB=A(L)	DIF04450
	GO TO 13	DIF04460
11	LL=L+LT	DIF04470
12	TB=A(LL)	DIF04480
	A(LL)=A(L)	DIF04490
13	AUX(1)=TB	DIF04500
14	A(L)=PIV1*TB	DIF04510
C		DIF04520
C	SAVE COLUMN INTERCHANGE INFORMATION	DIF04530
	A(LST)=I	DIF04540
C		DIF04550
		DIF04560

C	ELEMENT REDUCTION AND SEARCH FOR NEXT PIVOT	
	PIV=0.D0	D1F04570
	LLST=LST	D1F04580
	LT=0	D1F04590
	DO 18 I1=K,LEND	D1F04600
	PIV1=-ABS(I1)	D1F04610
	LL=LLST	D1F04620
	LT=LT+1	D1F04630
	DO 15 LLD=I1,LEND	D1F04640
	LL=LL+LLD	D1F04650
	L=LL+LT	D1F04660
15	A(L)=A(L)+PIV1*A(LL)	D1F04670
	LLST=LLST+11	D1F04680
	LR=LLST+LT	D1F04690
	TB=DABS(A(LR))	D1F04700
	IF(TB-PIV)17,17,16	D1F04710
16	PIV=TB	D1F04720
	I=LR	D1F04730
	J=I+1	D1F04740
17	DO 18 LR=J,MM,M	D1F04750
	LL=LR+LT	D1F04760
18	R(LL)=R(LL)+PIV1*R(LR)	D1F04770
	END OF ELIMINATION LOOP	D1F04780
C		D1F04790
C		D1F04800
C		D1F04810
C	BACK SUBSTITUTION AND BACK INTERCHANGE	D1F04820
19	IF(LEND)24,23,20	D1F04830
20	I1=M	D1F04840
	DO 22 I=2,M	D1F04850
	LST=LST-I1	D1F04860
	I1=I1-1	D1F04870
	L=A(LST)+.500	D1F04880
	DO 22 J=I1,MM,M	D1F04890
	TB=R(J)	D1F04900
	LL=J	D1F04910
	K=LST	D1F04920
	DO 21 LT=I1,LEND	D1F04930
	LL=LL+1	D1F04940
	K=K+LT	D1F04950
21	TB=TB-A(K)*R(LL)	D1F04960
	K=J+L	D1F04970
	R(J)=R(K)	D1F04980
22	R(K)=TB	D1F04990
23	RETURN	D1F05000
C		D1F05010
C		D1F05020
C	ERROR RETURN	D1F05030
24	IER=-1	D1F05040
	RETURN	D1F05050
	END	D1F05060

VO(5-MEO-SAL SB)(PHEN)

3.5 62 41684 5 3.500 2.000

7781 12.3940 480 3982 5965 7512 7781 6547 4781 3154 2002
 7542 12.3940 36180 3841 5772 7288 7542 6381 4561 3069 1950
 7828 12.3940 59580 3784 5637 7161 7828 6200 4466 3001 1892
 7151 12.3940 87480 3647 5444 6874 7151 6010 4312 2881 1843
 6854 12.3940 121980 3441 5191 6568 6854 5738 4104 2733 1733
 6649 12.3940 147480 3345 5014 6376 6649 5600 3995 2648 1669
 6410 12.3940 173580 3227 4872 6189 6410 5395 3859 2548 1604
 6036 12.3940 208080 3051 4593 5799 6036 5091 3600 2409 1495
 5914 12.3940 233880 2960 4511 5676 5914 4983 3552 2369 1505
 5735 12.3940 259380 2871 4343 5515 5735 4793 3421 2279 1441
 5423 12.3940 294780 2705 4098 5214 5423 4572 3254 2146 1343
 5272 12.3940 321780 2632 4000 5087 5272 4423 3166 2084 1322
 5078 12.3940 345780 2584 3896 4878 5078 4304 3052 2064 1248
 4883 12.3940 379080 2435 3719 4710 4883 4095 2919 1938 1208
 4719 12.3940 405780 2342 3569 4559 4719 3951 2810 1862 1159
 4515 12.3940 432180 2269 3434 4355 4515 3771 2682 1812 1119
 4339 12.3940 466680 2180 3265 4169 4339 3609 2584 1705 1071
 4211 12.3940 487080 2101 3175 4032 4211 3531 2486 1659 1047
 4034 12.3940 517980 2007 3066 3885 4034 3381 2406 1608 1018
 3858 12.3940 562080 1879 2879 3678 3858 3227 2255 1516 915
 3716 12.3940 577080 1843 2807 3578 3716 3100 2204 1472 919
 3547 12.3940 604980 1767 2696 3434 3547 2976 2130 1436 877
 3417 12.3940 640080 1673 2572 3283 3417 2869 2033 1318 841
 3334 11.7417 645580 1604 2492 3169 3334 2800 1948 1267 817
 3141 11.7417 692280 1528 2361 3020 3141 2629 1840 1221 773
 2929 12.3940 751980 1422 2206 2831 2929 2468 1711 1154 701
 2831 11.7417 777780 1373 2126 2733 2831 2377 1664 1110 689
 2709 11.7417 813180 1332 2023 2618 2709 2270 1597 1050 630
 2605 12.3940 841380 1270 1958 2518 2605 2194 1551 1015 595
 2561 11.7417 864180 1259 1903 2413 2561 2104 1482 957 599
 2324 11.7417 930480 1150 1743 2225 2324 1924 1388 878 543
 2155 11.7417 985080 1076 1621 2067 2155 1778 1253 841 514
 2103 11.7417 1007880 1019 1567 1991 2103 1746 1241 809 486
 2032 11.7417 1036380 990 1517 1928 2032 1660 1154 779 513
 1939 11.7417 1070580 922 1454 1863 1939 1622 1107 755 453
 1828 11.7417 1115880 844 1355 1737 1828 1529 1095 695 425
 1785 11.7417 1159380 844 1287 1662 1785 1404 979 697 438
 1660 11.7417 1201980 807 1216 1562 1660 1366 968 636 380
 1575 11.7417 1257780 756 1158 1486 1575 1279 868 614 347
 1451 12.3940 1296780 695 1121 1398 1451 1226 847 582 356
 1426 11.7417 1330380 690 1057 1341 1426 1177 822 564 338
 1166 11.7417 1376280 650 1032 1262 1166 1119 761 500 331
 1111 11.7417 1416780 604 961 1208 1111 1087 728 493 312
 1266 11.7417 1454280 590 925 1132 1266 998 701 484 300
 1233 11.7417 1503480 556 855 1136 1233 964 682 475 302
 1140 11.7417 1541280 564 846 1092 1140 926 658 454 268
 1122 11.7417 1589580 524 798 1053 1122 898 642 427 260
 1067 11.7417 1611480 526 797 1044 1067 900 641 403 274
 1075 12.3940 1635480 485 781 991 1075 896 637 422 260
 987 11.7417 1676280 490 744 950 987 828 615 391 255
 969 11.0894 1711980 466 746 932 969 804 582 377 229
 979 12.3940 1720680 491 739 930 979 809 531 409 225
 945 11.7417 1762380 448 717 907 945 769 589 359 219
 917 11.7417 1794780 463 686 893 917 786 545 373 224
 944 11.7417 1814280 449 697 883 944 768 541 375 230
 921 11.7417 1849080 450 699 847 921 749 520 361 231
 924 11.7417 1875480 456 681 885 924 740 513 356 225
 845 11.0894 1938480 387 640 797 845 729 476 354 208
 827 11.7417 1980780 423 638 804 827 721 482 337 205
 870 11.7417 2021580 419 605 806 870 665 489 313 209
 820 12.3940 2050080 386 623 784 820 692 494 310 224
 813 11.7417 2140080 388 605 761 813 692 471 330 179

FILE: DIFFVL OUTPUT A1 KING FAHD UNIVERSITY OF PETROLEUM AND MINERALS, DHAHRAM
 THE NAME OF THE SAMPLE IS CP2VCL2
 THE NUMBER OF LINES = 8
 THE SPIN α/β THE AVERAGE LINE WIDTH

-1.5	6.52724
-2.5	5.95845
-1.5	5.82001
-0.5	5.85751
0.5	6.14222
1.5	6.72765
2.5	7.52331
3.5	8.55251

THE β OF DATA POINTS IS 57 THE DESIRED POLYNOMIAL FIT IS 3

THE COEFFICIENTS OF POLYNOMIAL ARE
 $A(0) = 1.0170+00$
 $A(1) = -7.8520-07$
 $A(2) = 1.3850-13$
 $A(3) = 1.8270-20$

PL.#	TIME.SEC	EXP.INTI	CAL.INTI	SQR.OIFF
1	0.2400+03	0.1000+01	0.1020+01	0.2690-03
2	0.4800+05	0.9640+00	0.9790+00	0.1710-03
3	0.7440+05	0.9610+00	0.9590+00	0.4780-05
4	0.8760+05	0.9510+00	0.9490+00	0.3320-05
5	0.1380+06	0.9120+00	0.9110+00	0.1020-05
6	0.1530+06	0.9090+00	0.9000+00	0.8920-04
7	0.1730+06	0.8930+00	0.8930+00	0.7310-04
8	0.2120+06	0.8620+00	0.8570+00	0.3070-04
9	0.2410+06	0.8260+00	0.8360+00	0.9630-04
10	0.2590+06	0.8310+00	0.8230+00	0.6990-04
11	0.2970+06	0.8060+00	0.7960+00	0.9470-04
12	0.3180+06	0.7910+00	0.7810+00	0.8810-04
13	0.3460+06	0.7560+00	0.7630+00	0.4870-04
14	0.3810+06	0.7510+00	0.7360+00	0.1630-03
15	0.4090+06	0.7180+00	0.7200+00	0.4280-05
16	0.4110+06	0.7060+00	0.7060+00	0.6290-06
17	0.4490+06	0.6870+00	0.6810+00	0.4140-04
18	0.4980+06	0.6370+00	0.6620+00	0.2930-05
19	0.5120+06	0.6510+00	0.6510+00	0.4870-05
20	0.5550+06	0.6340+00	0.6270+00	0.5830-04
21	0.5790+06	0.6150+00	0.6120+00	0.1180-04
22	0.6040+06	0.5900+00	0.5970+00	0.4800-04
23	0.6670+06	0.5560+00	0.5600+00	0.1220-04
24	0.6900+06	0.5430+00	0.5470+00	0.1570-04

THE LEAST SQUARES FIT IS FOR THE $\alpha/\beta = 34$

THE VARIANCE = 0.687690-03

THE STANDARD DEVIATION = 0.262220-01

PL.#	VAL	INTENSITY	R-VALUE	DIFF.COEF
25	0.2270+06	0.3220+00	0.3260+00	0.1490-04
26	0.7580+06	0.5100+00	0.5090+00	0.1420-05
27	0.7780+06	0.4970+00	0.4980+00	0.1280-05
28	0.8130+06	0.4720+00	0.4780+00	0.14610-04
29	0.8430+06	0.4640+00	0.4640+00	0.4270-05
30	0.8600+06	0.4510+00	0.4530+00	0.2000-04
31	0.9010+06	0.4390+00	0.4330+00	0.1910-04
32	0.9270+06	0.4160+00	0.4220+00	0.3860-04
33	0.9400+06	0.4070+00	0.4160+00	0.8640-04
34	0.9660+06	0.3940+00	0.3940+00	0.5220-08
35	0.1010+07	0.3840+00	0.3830+00	0.6860-06
36	0.1040+07	0.3700+00	0.3700+00	0.4160-05
37	0.1070+07	0.3590+00	0.3590+00	0.5640-05
38	0.1070+07	0.3370+00	0.3340+00	0.1640-03
39	0.1120+07	0.3350+00	0.3380+00	0.1030-04
40	0.1170+07	0.3120+00	0.3120+00	0.8160-05
41	0.1190+07	0.2980+00	0.3070+00	0.8470-04
42	0.1210+07	0.3070+00	0.3010+00	0.3960-04
43	0.1230+07	0.2910+00	0.2880+00	0.6870-05
44	0.1270+07	0.2840+00	0.2790+00	0.1930-04
45	0.1300+07	0.2730+00	0.2710+00	0.1160-05
46	0.1330+07	0.2660+00	0.2590+00	0.3830-04
47	0.1360+07	0.2560+00	0.2500+00	0.3870-04
48	0.1380+07	0.2440+00	0.2440+00	0.1810-06
49	0.1420+07	0.2380+00	0.2330+00	0.2620-04
50	0.1440+07	0.2250+00	0.2270+00	0.2280-04
51	0.1470+07	0.2260+00	0.2200+00	0.3230-04
52	0.1510+07	0.2260+00	0.2110+00	0.2280-03
53	0.1530+07	0.2080+00	0.2010+00	0.5110-04
54	0.1590+07	0.1920+00	0.1910+00	0.5910-06
55	0.1690+07	0.1720+00	0.1730+00	0.3740-06
56	0.1800+07	0.1510+00	0.1590+00	0.6100-04
57	0.1880+07	0.1450+00	0.1510+00	0.4370-04

1 THE NAME OF THE SAMPLE IS VO(S-MEO-SAL SB)(PHEN)
 THE NUMBER OF LINES = 6
 THE SPIN Q, β THE AVERAGE LINE WIDTH

-1.5	17.19795
-2.5	13.69622
-1.5	12.32136
-0.5	12.02574
0.5	13.19674
1.5	15.71477
2.5	19.20445
3.5	24.37782

THE β OF DATA POINTS IS 62 THE DESIRED POLYNOMIAL FIT IS 4

THE COEFFICIENTS OF POLYNOMIAL ARE

A(0)= 1.0150+00
 A(1)= -1.2700-06
 A(2)= 6.3180-13
 A(3)= -1.4120-19
 A(4)= 9.6000-27

PL.#	TIME, SEC	EXP. INTI	CAL. INTI	SQR. DIFF
1	0.4600+03	0.1000+01	0.1010+01	0.2100+03
2	0.3620+03	0.9690+00	0.9700+00	0.5010+06
3	0.5960+03	0.9480+00	0.9420+00	0.4340+04
4	0.8750+03	0.9160+00	0.9080+00	0.4390+04
5	0.1220+06	0.8720+00	0.8700+00	0.7360+05
6	0.4470+06	0.8470+00	0.8420+00	0.2750+04
7	0.1740+06	0.8170+00	0.8140+00	0.1340+04
8	0.2080+06	0.7690+00	0.7780+00	0.6910+04
9	0.2340+06	0.7550+00	0.7520+00	0.9360+05
10	0.2590+06	0.7290+00	0.7270+00	0.4750+05
11	0.2950+06	0.6990+00	0.6940+00	0.1550+04
12	0.3220+06	0.6710+00	0.6690+00	0.4500+05
13	0.3460+06	0.6500+00	0.6480+00	0.4430+05
14	0.3790+06	0.6220+00	0.6200+00	0.3050+05
15	0.4060+06	0.5990+00	0.5980+00	0.1510+05
16	0.4320+06	0.5750+00	0.5770+00	0.4380+05
17	0.4670+06	0.5500+00	0.5500+00	0.2300+06
18	0.4870+06	0.5340+00	0.5350+00	0.2830+05
19	0.5180+06	0.5140+00	0.5130+00	0.1620+06
20	0.5620+06	0.4850+00	0.4830+01	0.3300+05
21	0.5770+06	0.4710+00	0.4730+00	0.1760+05
22	0.6050+06	0.4530+00	0.4550+00	0.7310+05
23	0.6490+06	0.4320+00	0.4340+00	0.2980+05

24	0.6660+06	0.4180+00	0.4190+00	0.2460+06
25	0.6920+06	0.3960+00	0.4030+00	0.5390+04
26	0.7520+06	0.3700+00	0.3720+00	0.2420+03
27	0.7780+06	0.3580+00	0.3590+00	0.1020+03
28	0.8130+06	0.3410+00	0.3410+00	0.1350+07
29	0.8410+06	0.3290+00	0.3280+00	0.9460+07
30	0.8640+06	0.3190+00	0.3180+00	0.1820+07
31	0.9300+06	0.2920+00	0.2910+00	0.1240+03
32	0.9850+06	0.2710+00	0.2700+00	0.6460+06
33	0.1010+07	0.2630+00	0.2620+00	0.4120+06
34	0.1040+07	0.2540+00	0.2530+00	0.9610+06
35	0.1070+07	0.2430+00	0.2420+00	0.1020+03
36	0.1120+07	0.2290+00	0.2280+00	0.1100+03
37	0.1160+07	0.2170+00	0.2160+00	0.1900+03
38	0.1200+07	0.2060+00	0.2050+00	0.1760+03
39	0.1260+07	0.1940+00	0.1920+00	0.4960+03
40	0.1300+07	0.1840+00	0.1830+00	0.8090+06
41	0.1330+07	0.1780+00	0.1760+00	0.1990+03
42	0.1380+07	0.1680+00	0.1680+00	0.2740+06
43	0.1420+07	0.1610+00	0.1610+00	0.1230+10
44	0.1450+07	0.1530+00	0.1530+00	0.2580+03
45	0.1500+07	0.1490+00	0.1480+00	0.7670+06
46	0.1540+07	0.1430+00	0.1430+00	0.4430+08
47	0.1590+07	0.1370+00	0.1370+00	0.3030+07
48	0.1610+07	0.1360+00	0.1350+00	0.1660+03
49	0.1640+07	0.1330+00	0.1320+00	0.8250+06
50	0.1680+07	0.1260+00	0.1260+00	0.3810+03
51	0.1710+07	0.1220+00	0.1220+00	0.5920+03
52	0.1720+07	0.1230+00	0.1230+00	0.1020+03
53	0.1760+07	0.1190+00	0.1210+00	0.8070+03
54	0.1790+07	0.1170+00	0.1180+00	0.1810+03
55	0.1810+07	0.1170+00	0.1170+00	0.1130+06
56	0.1850+07	0.1150+00	0.1150+00	0.5090+08
57	0.1880+07	0.1150+00	0.1130+00	0.3220+03
58	0.1940+07	0.1060+00	0.1090+00	0.7750+03
59	0.1980+07	0.1040+00	0.1070+00	0.2590+03
60	0.2020+07	0.1050+00	0.1050+00	0.4380+06
61	0.2050+07	0.1040+00	0.1040+00	0.1970+06
62	0.2140+07	0.1020+00	0.9950+01	0.4670+03

THE LEAST SQUARES FIT IS FOR THE PLT.# = 43

THE VARIANCE = 0.116950+03

THE STANDARD DEVIATION = 0.109040+01

EXP. VAL	TIME, SEC	INTENSITY	R-VALUE	DIFF. COEF
0.1617E+07	0.160973	0.5714290+00	0.6763190+03	
0.1617E+07	0.160970	0.680000+00	0.6551420+03	

APPENDIX 2

FILE: CPT2A OUTPUT AT KING FAND UNIVERSITY OF PETROLEUM AND MINERALS, DAMMAM.

THE CAL. B & C VALUES FOR 2-K USING STEWART & PORTE MAG. PAR FOR
DIFF. VALUES IN M, THE FREQ. 9.511 GHZ

COMP. TIME	B	C
6. RAYLINE-12	.1641	.054076
7. RAYLINE-12	.704407	.076332
8. RAYLINE-11	.49759	.77955
9. RAYLINE-11	.76488	.48481
10. RAYLINE-11	1.04135	.69165
11. RAYLINE-11	1.323	.89902
SAME ORDER OF		
COMP. TIME AS	.18137	.04914
ABOVE	.177534	.049182
	.43968	.2366
	.6642	.41256
	.8954	.54926
	1.1319	.765434
	1.3929	.943005
	1.5765	.954191
	1.7964	.801461
	1.94456	1.469116
	1.79616	.5161
	.976682	.67187
	1.179	.92072
	1.0296	.69662
	1.685	1.01915
	1.558	.32581
	.7227	.661
	.96464	.66464

FILE: CPT2A OUTPUT AT KING FAND UNIVERSITY OF PETROLEUM AND MINERALS, DAMMAM.

THE CAL. B & C VALUES FOR 2-Z USING STEWART & PORTE MAG. PAR FOR
DIFF. VALUES IN M, THE FREQ. 9.511 GHZ

COMP. TIME	B	C
6. RAYLINE-12	.1641	.05406
7. RAYLINE-12	.7045	.07632
8. RAYLINE-11	.4976	.77955
9. RAYLINE-11	.76488	.48481
10. RAYLINE-11	1.04135	.69165
11. RAYLINE-11	1.323	.8990
SAME ORDER OF		
COMP. TIME AS	.2032	.078033
ABOVE	.24982	.10315
	.60685	.37201
	.94185	.69267
	1.35433	.91676
	1.6921	1.18032
	.22996	.094192
	.28013	.12683
	.69651	.483711
	1.11695	.7677237
	1.54647	1.06926
	1.9603	1.51183
	.25045	.100362
	.204077	.14372
	.17282	.501559
	1.25132	.07196
	1.221758	1.236812
	1.2181	1.6937

FILE: CPT2A OUTPUT AT KING FAND UNIVERSITY OF PETROLEUM AND MINERALS, DAMMAM.

THE CAL. B & C VALUES FOR 2-Y USING STEWART & PORTE MAG. PAR FOR
DIFF. VALUES IN M, THE FREQ. 9.511 GHZ

COMP. TIME	B	C
6. RAYLINE-12	.19946	.0274212
7. RAYLINE-12	.24416	.09571
8. RAYLINE-11	.19547	.384568
9. RAYLINE-11	.93819	.400999
10. RAYLINE-11	1.29076	.854265
11. RAYLINE-11	1.44822	1.10789
SAME ORDER OF		
COMP. TIME AS	.22173	.084227
ABOVE	.21224	.11173
	.67316	.39933
	1.0161	.60997
	1.4876	.919926
	1.9635	1.27626
	.24004	.0950756
	.29187	.12594
	.741164	.44344
	1.1961	.767255
	1.65832	1.0694
	.26904	.11394
	.3223	.15066
	.69279	.32025
	1.4013	.90736
	1.93282	1.27991
	.292322	.12034
	.32403	.122593
	.94185	.501645
	1.54872	1.01794
	.388101	.172024
	.42942	.26375
	1.2763	.78435
	.361876	.183978
	.4449	.241769
	1.2896	.81331
	.19293	.20296
	.46293	.2701
	1.425	.9041
	.65412	.2531
	.56418	.35043
	1.7133	1.0999

THE CAL. B & C VALUES FOR Z' = X USING BAKALIK & HAYES MAG. PAR FOR
DIFF. VALUES OF N, THE FREQ.=9.511 GHZ

CORR. TIME	B	C	
6.88710E-12	.142395	.069719	
8.88710E-12	.17748	.091617	
2.44710E-11	.431849	.33435	N = 1
4.88710E-11	.66343	.58424	
6.88710E-11	.903776	.832302	
8.88710E-11	1.1484	1.0795	
SAME ORDER OF	.13581	.079244	
CORR. TIME AS	.16832	.104483	
ABOVE	.43074	.37152	N = 4
	.6671	.64568	
	.90598	.919643	
	1.14861	1.19354	
	.13576	.094162	
	.16755	.12409	
	.456016	.42868	N = 8
	.7201	.740543	
	.98047	1.05478	
	1.2426	1.37079	

THE CAL. B & C VALUES FOR Z' = Y USING BAKALIK & HAYES MAG. PAR FOR
DIFF. VALUES OF N, THE FREQ.=9.511 GHZ

CORR. TIME	B	C	
6.88710E-12	.140063	.0602662	
8.88710E-12	.171574	.0796071	
2.44710E-11	.4170408	.2880488	N = 2
4.88710E-11	.6629788	.4991402	
6.88710E-11	.9159114	.709363	
8.88710E-11	1.171911	.9197114	
SAME ORDER OF	.181371	.0808904	
CORR. TIME AS	.2189958	.1072044	
ABOVE	.5358927	.3789403	N = 4
	.9071857	.6522152	
	1.264468	.925948	
	1.624466	1.201343	
	.2507477	.1261714	
	.3030084	.1645478	N = 10
	.8414513	.5683533	
	1.402059	.9735658	
	.3792556	.221347	
	.470338	.289466	N = 30
	1.425019	.9633199	
	.6403737	.4096137	
	.8137168	.5314634	N = 100
	2.579395	1.770348	

THE CAL. B & C VALUES FOR Z' = Z USING BAKALIK & HAYES MAG. PAR FOR
DIFF. VALUES OF N, THE FREQ.=9.511 GHZ

CORR. TIME	B	C	
6.88710E-12	.14239	.069719	
8.88710E-12	.17748	.091617	
2.44710E-11	.43185	.33435	N = 1
4.88710E-11	.66383	.58424	
6.88710E-11	.903776	.832302	
8.88710E-11	1.1484	1.07947	
SAME ORDER OF	.13586	.064047	
CORR. TIME AS	.168377	.08414	
ABOVE	.43088	.30049	N = 4
	.66732	.52604	
	.906342	.75144	
	1.14908	.976399	
	.13584	.066919	
	.16765	.087896	
	.45627	.30514	N = 8
	.72055	.53168	
	.98111	.76034	
	1.2434	.98986	

FILE: CP2XC OUTPUT AT KING FAHD UNIVERSITY OF PETROLEUM AND MINERALS, DHAHRAN

THE CAL. B & C VALUES FOR Z'=X USING CASEY & RAYNOR MAG. PAR. FOR
DIFF. VALUES OF N, THE FREQ.=9.511 GHZ

CORR.TIME	B	C	
6.88X10E-12	.20938	.051146	
8.88X10E-12	.26096	.067211	
2.88X10E-11	.63499	.24529	N = 1
4.88X10E-11	.97608	.42869	
6.88X10E-11	1.3289	.61084	
8.88X10E-11	1.6886	.7925	
SAME ORDER OF CORR.TIME AS ABOVE	.1457807	.03365571	
	.1845777	.04388695	
	.4753854	.1581491	N = 4
	.705093	.280999	
	.9357591	.4038153	
	1.172129	.5260099	
	.1118138	.02519779	
	.1425162	.03273453	
	.3889812	.1155009	N = 8
	.5730618	.2068584	
	.7467171	.2996174	
	.9226121	.3921521	

FILE: CP22C OUTPUT AT KING FAHD UNIVERSITY OF PETROLEUM AND MINERALS, DHAHRAN

THE CAL. B & C VALUES AT Z'=Z USING CASEY & RAYNOR MAG. PAR. FOR
DIFF. VALUES OF N, THE FREQ.=9.511 GHZ

CORR.TIME	B	C	
6.88X10E-12	.20938	.051146	
8.88X10E-12	.260964	.067211	
2.88X10E-11	.63499	.245297	N = 1
4.88X10E-11	.9761	.428694	
6.88X10E-11	1.3289	.61084	
8.88X10E-11	1.68859	.7924662	
SAME ORDER OF CORR.TIME AS ABOVE	.3298447	.09292165	
	.3994029	.1232105	N = 4
	1.014741	.4352556	
	1.64664	.748131	
	.4116183	.1314404	
	.4971032	.1738935	N = 8
	1.350281	.5976089	

FILE: CP2VC OUTPUT AT KING FAHD UNIVERSITY OF PETROLEUM AND MINERALS, DHAHRAN

THE CAL. B & C VALUES FOR Z'=Y USING CASEY & RAYNOR MAG. PAR FOR
DIFF. VALUES OF N, THE FREQ.=9.511 GHZ

CORR.TIME	B	C	
6.88X10E-12	.20938	.051146	
8.88X10E-12	.260964	.067211	
2.88X10E-11	.63499	.2453	N = 1
4.88X10E-11	.9761	.4287	
6.88X10E-11	1.3289	.61084	
SAME ORDER OF CORR.TIME AS ABOVE	.2947512	.0866697	
	.358445	.1188427	N = 4
	.911908	.4060238	
	1.467125	.698853	
	.3544574	.1202329	
	.4294979	.1590027	N = 8
	1.1670	.5467515	
	1.915462	.9349394	

FILE: CP2XD OUTPUT AT KING FAHD UNIVERSITY OF PETROLEUM AND MINERALS, DHAHRAN

THE CAL. B & C VALUES FOR Z'²X USING PETERSEN & DAHL MAG. PAR AT
DIFF. VALUES OF N THE FREQ.=9.511 GHZ

CORR. TIME	B	C	
6.88X10E-12	.17439	.052638	
8.88X10E-12	.21736	.069172	
2.88X10E-11	.52889	.252457	N = 1.0
4.88X10E-11	.81300	.441216	
6.88X10E-11	1.10687	.6287	
8.88X10E-11	1.40645	.81567	
SAME ORDER OF	.123124	.0349899	
CORR. TIME AS	.155723	.0456389	
ABOVE	.400941	.164411	N = 4.0
	.595986	.291979	
	.791918	.419518	
	.992593	.546438	
	.0959023	.0265643	
	.121981	.0345285	
	.332873	.121737	N = 8.0
	.492641	.217747	
	.6438824	.315219	
	.796998	.412501	

FILE: CP2YD OUTPUT AT KING FAHD UNIVERSITY OF PETROLEUM AND MINERALS, DHAHRAN

THE CAL. B & C VALUES FOR Z'²Y USING PETERSEN & DAHL MAG. PAR AT
DIFF. VALUES OF N THE FREQ.= 9.511 GHZ

CORR. TIME	B	C	
6.88X10E-12	.174392	.052638	
8.88X10E-12	.21736	.069172	N = 1
2.88X10E-11	.52889	.252457	
4.88X10E-11	.812997	.441216	
6.88X10E-11	1.10687	.5287	
8.88X10E-11	1.40646	.81567	
SAME ORDER OF	.21706	.073773	
CORR. TIME AS	.26607	.097589	
ABOVE	.65864	.35019	N = 2.7
	1.04358	.605337	
	1.437498	.859878	
	.282285	.119276	
	.342421	.157710	N = 8
	.930494	.542482	
	1.52378	.930171	

FILE: CP2ZD OUTPUT AT KING FAHD UNIVERSITY OF PETROLEUM AND MINERALS, DHAHRAN

THE CAL. B & C VALUES FOR Z'²Z USING PETERSEN & DAHL MAG. PAR FOR
DIFF. VALUES OF N , THE FREQ.=9.511 GHZ

CORR. TIME	B	C	
6.88X10E-12	.17439	.052638	
8.88X10E-12	.21736	.0691718	
2.88X10E-11	.528891	.252457	N = 1
4.88X10E-11	.812997	.441216	
6.88X10E-11	1.10686	.6287	
8.88X10E-11	1.406455	.8156679	
SAME ORDER OF	.280983	.097772	
CORR. TIME AS	.3399639	.1296689	N = 4
ABOVE	.86351	.497974	
	1.403485	.78691	
	.353024	.139111	
	.4261	.18406	N = 8
	1.157812	.632495	
	1.91514	1.082584	

FILE: CP2Y27AL OUTPUT AT KING FAHD UNIVERSITY OF PETROLEUM AND MINERALS, DHAHRAN

THE CAL B & C VALUES FOR Z'²=Y AT N=2.7 USING DIFFERANTE MAG. PAR

CORR. TIME	B	C	
6.88710E-12	.21542	.081018	
8.88710E-12	.26338	.107161	
2.88710E-11	.6.111	.384558	STEWART & FORT
4.88710E-11	1.03699	.6648	
6.88710E-11	1.431913	.944341	
SAME ORDER	.15701	.068077	
CORR. TIME	.19095	.090098	
AS ABOVE	.470725	.323174	BAKALIK & HAYES
.	.75775	.558183	
.	1.05155	.792746	
8.88710E-11	1.34814	1.02792	
	.267064	.073245	
	.32698	.09692	
	.80891	.34763	CASEY & RAYBOR
	1.28475	.60063	
	1.7717	.85295	
	.21706	.073773	
	.26607	.097589	
	.65864	.35019	PETERSEN & DAHL
	1.04358	.605337	
	1.437498	.859878	

FILE: CP2YDFRE OUTPUT AT KING FAHD UNIVERSITY OF PETROLEUM AND MINERALS, DHAHRAN

THE CAL. B & C VALUES FOR Z'²=Y AND N=2.7 AT DIFF. FREQ USING MAG. PAR FOR PETERSEN & DAHL

CORR. TIME	B	C	
6.88710E-12	.2131	.0736264	
8.88710E-12	.261461	.097368	
2.88710E-11	.64653	.34979	AT 9.3 GHZ
4.88710E-11	1.0225	.60501	
6.88710E-11	1.4072	.85962	
SAME ORDER	.21498	.073696	
OF CORR. TIME	.26365	.097473	
AS ABOVE FOR	.65227	.349982	AT 9.4 GHZ
ALL SETS	1.03247	.605168	
	1.4216	.85975	
	.21706	.073773	
	.26607	.097589	
	.65864	.35019	AT 9.511 GHZ
	1.04358	.605337	
	1.437498	.859878	
	.21872	.0738345	
	.268003	.097642	
	.663759	.350357	AT 9.6 GHZ
	1.05249	.605468	
	1.45028	.85998	
	.22058	.0739035	
	.270165	.0977856	
	.66951	.350538	AT 9.7 GHZ
	1.06251	.605612	
	1.4646	.860094	